

Chapter 5

Inverse Spectral Theory for Symmetric and Self-adjoint Operators

We prove that an operator measure in general is non-orthogonal and unbounded. We prove that two orthogonal spectral measures are unitarily equivalent. In accordance with the stieltjes inversion formula the spectral measure admits an analytic continuation [90]. We discuss and prove a sharp estimate that a strictly monotone function on each component interval of the inverse function is analytic and also Strictly monotone. We prove that a non-orthogonal spectral measure in a gab of any self-adjoint extensions can be calculated, if exist a boundary triple and have various examples [93].

Section (5-1): Inverse Spectral Problem for Direct Sum of Symmetric Operators

Let S be a densely defined symmetric operator in Hilbert space H with deficiency indices $n_+(S) = n_-(S) \leq \infty$. We recall that abounded open interval $J = (\alpha, \beta)$ is called a gap for S if

$$\|2S - (\alpha + \beta)\| \geq (\alpha - \beta) \|f\|, f \in \text{dom } S, \quad (1)$$

if $\beta = \infty$, then (1) turns into $(Sf, f) \geq \beta \|f\|^2$ for all $f \in \text{dom } S$, meaning that

$(-\infty, \beta)$, is a gap for S if S is semi bounded below with the lower bound β .

Theorem (5-1-1) [131]:

Let $\{S_k\}_{k=1}^\infty$ be a family of closed symmetric operators S_k , defined in the separable Hilbert space H such that the operators S_k are unitarily equivalent to a closed symmetric operator S in H with equal positive deficiency indices. If there exists a boundary triple $\Gamma_b = \{H_b, \Gamma_b^0, \Gamma_b^1\}$ for S such

that the corresponding Weyl function $M(\cdot)$ is monotone with respect to open set $J \subseteq \rho(A_0)$, $A_0 = A^*|_{\ker(\Gamma_0^0)}$, then for any auxiliary self-adjoint operator \mathcal{A} in some separable Hilbert space \mathcal{H} the closed symmetric operator \mathcal{S} admits a self-adjoint extension \mathcal{S}^* such that the spectral parts \mathcal{E}_J and \mathcal{E}_K are unitarily equivalent i.e. $\mathcal{S}_J^0 \cong_{R_J}$ [95.109,110].

The following result is known as a generalized Nuimark dilation theorem.

Proposition (5-1-2) [131]:

If $\Sigma(\cdot): B(R) \rightarrow \mathcal{K}(H)$ is a bounded operator measure, then there exist a Hilbert space K bounded operator $k \in [H, K]$ and an orthogonal measure $E(\cdot): B(R) \rightarrow \mathcal{K}(K)$ (an orthogonal dilation) such that

$$\Sigma(\mathcal{S}) = k^* E(\mathcal{S}) k, \quad \mathcal{S} \in B(R) \quad (2)$$

If the orthogonal dilation is minima i.e.,

$$\text{span}\{E(\mathcal{S}) \text{ran}(k) : \mathcal{S} \in B(R)\} = K, \quad (3)$$

then it is uniquely determined up to unitary equivalence that is if one has two bounded operator $k \in [H, K]$ and $K' \in [H, K]$ as well as two minimal orthogonal dilation $E(\cdot): B(R) \rightarrow \mathcal{K}(K)$ and $E'(\cdot): B(R) \rightarrow \mathcal{K}(K')$ obeying $\Sigma(\mathcal{S}) = k^* E(\mathcal{S}) k = K'^* E'(\mathcal{S}) K', \mathcal{S} \in B(R)$, then there exists an isometry $v: K' \rightarrow K$ such that $E'(\mathcal{S}) = v^* E(\mathcal{S}) v, \mathcal{S} \in B(R)$.

Definition (5-1-3) [131]:

We call $E(\cdot)$ satisfying (2) and (3) the minimal orthogonal measure associated to $\Sigma(\cdot)$, or the minimal orthogonal dilation of $\Sigma(\cdot)$.

Every operator measure $\Sigma(\cdot)$ admits the Lebesgue Jordan decomposition

$$\Sigma = \Sigma^{ac} + \Sigma^s, \quad \Sigma^s = \Sigma^{sc} + \Sigma^{pp} \quad \text{where } \Sigma^{ac}, \Sigma^s, \Sigma^{sc} \text{ and } \Sigma^{pp} \text{ are the absolutely}$$

continuous, singular, singular continuous and pure point components

(measure) of $\Sigma(\cdot)$, respectively. Non-topological supports of mutually disjoint, therefore if an operator measure \mathbb{J} is orthogonal, $\Sigma(\cdot) = E_T(\cdot)$, then the ortho-projections $P^\tau = E_T^\tau(R)$ ($\tau \in \{ac, sc, pp\}$) are pair wise orthogonal. Every subspace $h_t^\tau = P^\tau h$ reduces the operator $T = T^*$ and the Lebesgue-Jordan decomposition yields

$$\begin{aligned} h &= h_t^{ac} \oplus h_t^{sc} \oplus h_t^{pp} \\ T &= T^{ac} \oplus T^{sc} \oplus T^{pp} \end{aligned} \quad (4)$$

Where $T^\tau = P^\tau T \upharpoonright_{h_t^\tau}$, $T \in \{ac, sc, pp\}$. Now we show Nevanlinna functions:

Let \mathcal{H} be a separable Hilbert space, we recall that an operator-valued function $F : C_+ \rightarrow \mathcal{H}$ is said to be a Nevanlinna (or Herglotz or \mathcal{H}) one if it is holomorphic and takes values in the set of dissipative operators on \mathcal{H} i.e.,

$$\operatorname{Im}(F(z)) = \frac{F(z) - F(z)^*}{2i} \geq 0, z \in C_+$$

Usually one considers a continuation of F in \mathcal{H} by setting $F(z) =$

$F(\bar{z}), z \in C_-$. Bounded operator $k \in [\mathcal{H}, K]$ obeying $\ker(K) = \ker \sum_F^0(R)$ and

$\sum_F^0(\delta) = k^* E_F(\delta) k, \delta \in B(R)$. By

$$\sum_F(\delta) = \int_{\delta} (1+t^2) d \sum_F^0(t), \delta \in B_b(R) \quad (5)$$

One defines an operator measure which in general is non-orthogonal and unbounded. It is called the unbounded spectral measure of $F(\cdot)$. Using

\mathbb{J}_F the representation [118],

$$F(z) = C_0 + C_1 z + \int_{-\infty}^{+\infty} \left(\frac{1}{t-z} - \frac{1}{1+t^2} \right) d \sum_F(t), z \in C_+ \cup C_- \quad (6)$$

To show this we have [5]:

From this representation $F(z) = C_0 + C_1 z + \int_{-\infty}^{+\infty} \frac{1+tz}{t-z} d\sum_F^0(t), z \in C_+ \cup C_-$. To prove representation (6) use equation (5)

$$\sum_F(\delta) = \int_{\delta} (1+t^2) d\sum_F^0(t), \delta \in B_b(R)$$

so $d\sum_F(\delta) = (1+t^2) d\sum_F^0(t)$, which implies that $d\sum_F^0(t) = \frac{1}{1+t^2} d\sum_F(t)$, put this in the representation above we have

$$F(z) = C_0 + C_1 z + \int_{-\infty}^{+\infty} \frac{1+tz}{t-z} \left(\frac{1}{1+t^2} \right) d\sum_F(t) = C_0 + C_1 z + \int_{-\infty}^{+\infty} \frac{1+tz}{(t-z)(1+t^2)} \left(\frac{1}{1+t^2} \right) d\sum_F(t)$$

To analysis this component we use this $\frac{1+tz}{(t-z)(1+t^2)} = \frac{A}{t-z} + \frac{Bt}{1+t^2} + \frac{C}{1+t^2} = 1+tz$

and $A(1+t^2) + Bt(t-z) + C(t-z) = 1+tz$ put $t=z$ we get $A(1+z^2) = 1+z^2$, so

$A=1$ at $t=0, A-Cz = 1$ implies that $C=0$ since $A=1, C=0$. Our equation

become $1+t^2 + Bt(t-z) + 0 = 1+tz$, $Bt(t-z) = 1+tz - 1 - t^2 = -t(t-z)$, $Bt = -t \left(\frac{t-z}{t-z} \right)$

$B=-1$. Substituted A, B , and C the equation

$$\frac{1+tz}{(t-z)(1+t^2)} = \frac{A}{t-z} + \frac{Bt}{1+t^2} + \frac{C}{1+t^2} = 1+tz$$

We get the following

$$F(z) = C_0 + C_1 z + \int_{-\infty}^{+\infty} \left(\frac{1}{t-z} - \frac{t}{1+t^2} \right) d\sum_F^0(t) \quad (7)$$

$$z \in C_+ \cup C_-$$

Which complete the proof. From representation

$$F(z) = C_0 + C_1 z = \int_{-\infty}^{\infty} \frac{1-tz}{t-z} d\sum_F^0(t), z \in C_+ \cup C_-$$

F determines uniquely the unbounded spectral measure $\sum_F^0(\cdot)$ by means of the Stieltjes inversion formula, which is given by

$$\sum_F((a, b)) = s - \lim_{\delta \rightarrow +0} s - \lim_{\varepsilon \rightarrow +0} \frac{1}{\pi} \int_{a+\delta}^{b-\delta} \text{Sm}(F(x+i\varepsilon)) dx \quad (8)$$

By $\text{supp}(F)$ we denote the topological (minimal closed) support of the spectral measure Σ_F . Since $\text{supp}(F)$ is closed the set $O_F = \mathbb{R} \setminus \text{supp}(F)$ is open. The Nevanlinna function $F(\cdot)$ admits an analytic continuation to \mathbb{C} given by

$$F(\lambda) = C_0 + C_1 \lambda + \int_{-\infty}^{+\infty} \left(\frac{1}{t-\lambda} - \frac{t}{1+t^2} \right) d\Sigma_F(t), \lambda \in O_F$$

Using this representation we immediately find that $F(\cdot)$ is monotone on each component interval Δ of O_F i.e., $F(\lambda) \leq F(\mu)$, $\lambda < \mu$, $\lambda, \mu \in \Delta$. In general, this relation is not satisfied if λ and μ belong to different component interval.

Definition (5-1-4) [131]:

Let $F(\cdot)$ be a Nevanlinna function, the Nevanlinna function is monotone with respect to the open set $J \subseteq O_F$ if for any two component intervals I_1 and I_2 of J one has $F(\lambda) \leq F(\mu)$ for all $\lambda \in I_1$ and $\mu \in I_2$ or $F(\lambda) \geq F(\mu)$ for all $\lambda \in I_1$ and $\mu \in I_2$.

Let $L \in \mathbb{N} \cup \infty$ be the number of component interval of J . obviously if $F(\cdot)$ is monotone with respect to J and $L < \infty$, then there exists an enumeration

$\{J_k\}_{k=1}^L$ of the components of J such that

$$F(\lambda) \leq F(\mu) \leq \dots \leq F(\nu)$$

Holds for $\{\lambda, \mu, \dots, \nu\} \in J_1 \times J_2 \times \dots \times J_L$. If $L = \infty$, then it can happen that such an enumeration does not exist. If $F(\cdot)$ is a scalar Nevanlinna function, then $F(\cdot)$ is monotone with respect to J if and only if the condition

$$F(J_1) \cap F(J_2) = \emptyset$$

is satisfied for any two component intervals I_1 and I_2 of J .

Definition (5-1-5) [131]:

A triple $\Pi = \{H, \Gamma_0, \Gamma_1\}$ consisting of an auxiliary Hilbert space H and linear mappings $\Gamma_i : \text{dom}(A^*) \rightarrow H, i = 0, 1$. Called a boundary triple for the adjoint operator A^* of A if the following two conditions are satisfied:

(i) The second Green's formula takes place

$$(A^*f, g) - (f, A^*g) = (\Gamma_1 f, \Gamma_0 g) - (\Gamma_0 f, \Gamma_1 g), f, g \in \text{dom}(A^*)$$

(ii) The mapping $\Gamma = \{\Gamma_0, \Gamma_1\} : \text{dom}(A^*) \rightarrow H \oplus H$, $\Gamma f = \{\Gamma_0 f, \Gamma_1 f\}$ is surjective. The above definition allows one to describe the set Ext_A in the following way.

Proposition (5-1-6) [131]:

Let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for A then the mapping $\Gamma : \text{Ext}_A \rightarrow \mathcal{L}(H, H)$ established abjective correspondence $\mathcal{L}(H, H) \rightarrow \Gamma(\text{dom}(A^*))$ between the set Ext_A of self-adjoint linear relations in H . By proposition (5-1-6) the following definition is natural [144,145].

Definition (5-1-7) [131]:

Let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for A . We put $A_\theta = A^*|_{D_\theta}$, if

$$\theta = \Gamma(\text{dom}(A^*)) \text{ that is } A_\theta = A^*|_{D_\theta},$$

$$\text{dom}(A_\theta) = D_\theta = \{f \in \text{dom}(A^*) : \{\Gamma_0 f, \Gamma_1 f\} \in \theta\} \quad (9)$$

If $\theta = G(B)$ is the graph of an operator $B = B^* \in C(H)$, then $\text{dom}(A_\theta)$ is determined by the equation $\text{dom}(A_B) = D_B = \ker(\Gamma_1 - B\Gamma_0)$. We set $A_B = A_\theta$.

Let us recall the basic facts on Weyl functions:

Definition (5-1-8) [131]:

Let A be a densely defined closed symmetric operator and

$\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for A . The unique mapping

$$M(\cdot) = \mathcal{P}(A_0) \rightarrow \mathcal{L}(H) \text{ defined by } \Gamma_1 f_z = M(z) \Gamma_0 f_z, f_z \in N_z = \ker(A^* - z), z \in C_+$$

Is called the Weyl function corresponding to the boundary triple Π .

Proposition (5-1-9) [131]:

Let T be a simple closed symmetric operator and let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for T^* with Weyl function $M(\lambda)$. Suppose that M is self-adjoint linear relation in \mathcal{H} and $\lambda \in \rho(A_0)$ then

- (i) $\delta(A_0) = \text{supp } M$
- (ii) $\lambda \in \rho(A_0)$ if and only if $\Theta \in \rho(\Theta - M(\lambda))$
- (iii) $\lambda \in \delta_r(A_0)$ if and only if $0 \in \delta_r(\Theta - M(\lambda)), T \in [p, c]$

We need the following simple proposition.

Proposition (5-1-10) [131]:

Let T be a closed symmetric operator and let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for T^*

- (i) If T is simple and $\Pi_1 = \{H_1, \Gamma_0^1, \Gamma_1^1\}$ is another boundary triple for T^* such that $\ker(\Gamma_0) = \ker(\Gamma_1^1)$, then the Weyl functions $M(\cdot)$ and $M_1(\cdot)$

of Π and Π_1 , respectively are related by $M_1(z) = k^* M(z) k + D$,

$z \in C_+ \cup C_-$. Where $D = D^* \in [H]$ and $k \in [H_1, H]$ is boundedly

invertible.

- (ii) If $\Theta = G(B)$, $B = B^* \in H$, then the Weyl function $M_B(\cdot)$ corresponding

to the boundary triple $\Pi_B = \{H, \Gamma_0^B, \Gamma_1^B\} = \{H, B\Gamma_0 - \Gamma_1, \Gamma_0\}$ is given by

$$M_B(z) = (B - M(z))^{-1}, z \in \mathbb{C}_+ \cup \mathbb{C}_-$$

Definition (5-1-11) [131]:

Let T be a densely defined closed symmetric operator and let

$\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for T^* . The mapping $\rho(A_0) \ni z \rightarrow \gamma(z) \in [H, N_z]$

$$\gamma(z) = (\Gamma_0|_{N_z})^{-1} : H \rightarrow N_z, z \in \rho(A_0)$$

is called the γ -filed of the boundary triple Π . One can easily have

$$\gamma(z) = (A_0 - z_0)(A_0 - z_0)^{-1} \gamma(z_0), z, z_0 \in \rho(A_0) \quad (10)$$

The \mathcal{H} -field and the Weyl function $M(\cdot)$ are related by

$$M(z) - M(z_0)^* = (z - \bar{z}_0) \gamma(z_0)^* \gamma(z)$$

Lemma (5-1-12) [131]:

Let T be a simple densely defined closed symmetric operator on a separable Hilbert space \mathcal{H} with equal deficiency indices. Further let

$\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for T with Weyl function $M(\cdot)$. If $E_{A_0}(\cdot)$

is the orthogonal spectral measure of T define on \mathcal{H} and $E_M(\cdot)$ the

associated minimal orthogonal spectral dilation of $\Sigma_M^0(\cdot)$ defined on such that

$$E_{A_0}(\delta) = W^* E_M(\delta) W \quad \text{for any Borel set } \delta \in B(\mathbb{R}).$$

Proof: By (10) one obtains

$$S(M(x+iy)h, h) = \gamma(\gamma(x+iy)h, (x+iy)h)h \in \mathcal{H} \quad (11)$$

To show this we have [5]:

$$Sm(M(z)h, h) = \frac{(M(z)h, h) - (M(z)h, h)^*}{2i}$$

Where

$$\begin{aligned} z = x + iy &= |h| \left[(M(z), 1) - (M(z), 1)^* \right] / 2i \\ &= |h| \left[(z - \bar{z}_0) \gamma(z_0)^* \gamma(z) + M(z_0) - (z - \bar{z}_0)^* \gamma(z)^* - M(z_0)^* \right] / 2i \end{aligned}$$

Multiply and divided by $(z - \bar{z}_0) \gamma(z_0)^*$

$$\begin{aligned} &= \frac{|h|}{2i} \left[\frac{(z - \bar{z}_0) \gamma(z_0)^* \gamma(z)}{(z - \bar{z}_0) \gamma(z_0)^*} - \frac{(z - \bar{z}_0) \gamma(z_0) \gamma(z_0)^*}{(z - \bar{z}_0) \gamma(z_0)^*} \right] \\ &= \frac{|h|}{2i} [\gamma(z) - \gamma^*(z)] = \frac{|h|}{2i} [\gamma(z) - \gamma^*(z)] \\ &= \frac{|h|}{2i} [\gamma(z) - (\bar{z}_0) \gamma(\bar{z}_0)^* \gamma(z)] \end{aligned}$$

$$\begin{aligned}
&= |h| \left[\frac{\gamma(z) - \gamma^*(z)}{2} (z - \bar{z}_0) \gamma^*(z_0) \right] \\
&= \frac{|h|}{2i} \left[(z - \bar{z}_0) \gamma^*(\bar{z}_0) \gamma(z) - \gamma^*(z) \gamma^*(z_0) (z - \bar{z}_0) \right] \\
&= \frac{|h|}{2i} \gamma^*(z_0) \left[(z - \bar{z}_0) \gamma(z) - (z - \bar{z}_0) \gamma^*(z) \right]
\end{aligned}$$

Where $\gamma^*(z_0)/i \cdot 2 = y \implies |h|y [\gamma(z), \gamma^*(z)] = y(\gamma(z)h, \gamma^*(z)h)$

Since $z = x + iy$, we get

$$Sm(M(x + iy)h, h) = y(\gamma(x + iy)h, \gamma^*(x + iy)h)$$

Which is the prove of (11). Further, it follows from (10) that

$$\gamma(x + iy) = \left[I + (x + i(y - 1))(A_0 - x - iy)^{-1} \right] \gamma(i) \quad (12)$$

To prove (12) we use (10) [5]:

$$\begin{aligned}
\gamma(z) &= (A_0 - z)(A_0 - z)^{-1} \gamma(z_0) \\
\gamma(z) &= A_0(A_0 - z)^{-1} \gamma(z_0) - z_0(A_0 - z)^{-1} \gamma(z_0) \\
&= A_0 \frac{1}{A_0} \left(I - \frac{z}{A_0} \right)^{-1} \gamma(z_0) - z_0(A_0 - z)^{-1} \gamma(z_0) \\
&= \left[(I - ZA_0^{-1})^{-1} - Z_0(A_0 - z)^{-1} \right] \gamma(z_0) \\
&= \left[\left(I + \sum_{n=1}^{\infty} Z^n \|A_0^{-1}\|^n - Z_0(A_0 - Z)^{-1} \right) \right] \gamma(Z_0) \\
&= \left[I + \sum_{n=1}^{\infty} Z^{n+1} \|A_0^{n+1}\| - Z_0(A_0 - Z)^{-1} \right] \gamma(Z_0)
\end{aligned}$$

Since $A_0 = A^*$ is self adjoint spectrum and $\|A_0^{n+1}\| = 1$, so

$$\begin{aligned}
\gamma(z) &= \left[I + \sum_{n=0}^{\infty} Z^{n+1} - Z_0(A_0 - Z)^{-1} \right] \gamma(z_0) \\
&= \left[I + \sum_{n=0}^{\infty} Z^n \cdot Z - Z_0(A_0 - Z)^{-1} \right] \gamma(z_0)
\end{aligned}$$

But $\sum_{n=0}^{\infty} z^n = (A_0 - z)^{-1}$

Hence $\mathcal{H}(z) = [I + z(A_0 - z)^{-1} - Z_0(A_0(A_0 - z)^{-1})]\mathcal{H}(Z_0)$

$$= [I + (z - Z_0)(A_0 - Z)^{-1}]\mathcal{H}(Z_0)$$

Let $x = 0, y = 1 \implies z_0 = 0 + i$

Therefore $\mathcal{H}(z) = [I + (z - i)(A_0 - z)^{-1}]\mathcal{H}(i)$

Since $z = x + iy$

$$\mathcal{H}(x + iy) = [I + (x + iy - i)(A_0 - (x + iy))^{-1}]\mathcal{H}(i)$$

$$= [I + (x + i(y - 1))(A_0 - x - iy)^{-1}]\mathcal{H}(i)$$

Which is the proof of (12). Inserting (12) into (11) one gets

$$Sm(M(x + iy)h, h) = y \int_{-\infty}^{+\infty} \frac{1+t^2}{(t-x)^2 + y^2} d(E_{A_0}(t) \mathcal{H}(i)h, \mathcal{H}(i)h), h \in H$$

On the other hand we obtain that $d(\sum_M(t)h, h) = (1+t^2)d(E_{A_0}(t) \mathcal{H}(i)h, \mathcal{H}(i)h)$,

inserting in the above representation we get

$$Sm(M(x + iy)h, h) = \int_{-\infty}^{+\infty} \frac{d(\sum_M(t)h, h)}{(t-x)^2 + y^2}, \quad h \in H$$

Applying the stieltjes inversion formula (8) we find

$$\left(\sum_M((a, b))h, h \right) = \int_{(a, b)} (1+t^2) d(E_{A_0}(t) \mathcal{H}(i)h, h), h \in H$$

Which yields

$$\sum_M^0((a, b)) = \mathcal{H}(i)^* E_{A_0}((a, b)) \mathcal{H}(i) \quad (13)$$

for any bounded open interval $(a, b) \subseteq \mathbb{R}$. Since \mathcal{H} is simple it follows from (12) that

$$\{(A_0 - \lambda)^{-1} \tan(\mathcal{H}(i)) : \lambda \in C_+ \cup C_-\} = H \quad (14)$$

By (13) and (14), $E_{A_0}(\cdot)$ is a minimal orthogonal dilation of $\Sigma_M^0(\cdot)$. By proposition (5-1-2) we find that the spectral measure $E_{A_0}(\cdot)$ and $E_M(\cdot)$ are unitarily equivalent.

Definition (5-1-13) [131]:

Let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for \mathcal{H} with corresponding Weyl function $M(\cdot)$. We will call $\Sigma_M^0(\cdot)$ the bounded non-orthogonal spectral measure of the extension $A_0 = (A^* | \ker(\Gamma_0))$.

Corollary (5-1-14) [131]:

Let \mathcal{H} be a simple densely defined closed symmetric operator in a separable Hilbert space \mathcal{H} with equal deficiency indices. Further, let

$\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for \mathcal{H} and $M(\cdot)$ the corresponding Weyl function, then

$$\mathcal{S}(A_0) = \text{supp}(M) = \text{supp}\left(\sum_M^r\right), \quad \mathcal{S}_r(A_0) = \text{supp}\left(\sum_M^r\right). \quad \text{Where } \tau \in \{ac, s, sc, pp\}$$

Remark (5-1-15) [131]:

$M_B(\cdot)$ of the form $M_B(z) = (B - M(z))^{-1} = (B - m(z) \cdot I_H)^{-1}$ is the Weyl function of the generalized boundary triple Π_B . Being a Weyl function. $M_B(\cdot)$ admits the representation

$$M_B(z) = C_0 + \int_{-\infty}^{+\infty} \left(\frac{1}{t-z} - \frac{t}{1+t^2} \right) d \sum_B(t), \quad z \in C_+ \cup C_- \quad (15)$$

Where $\sum_B(\cdot) = \sum_{MB}(\cdot)$ is the (unbounded) non-orthogonal spectral measure of

$M_B(\cdot)$. In accordance with the Stieltjes inversion formula (8) the spectral measure can be re-obtained by

$$\sum_B(a, b) = s\text{-}\lim_{\delta \rightarrow 0} s\text{-}\lim_{\epsilon \rightarrow 0} \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} (M_B(x+i\epsilon) - M_B(x-i\epsilon)) dx \quad (16)$$

With $M(z) = \overline{M(\bar{z})}^*$. We get

$$M_B(x+i\varepsilon) - M_B(x-i\varepsilon) = \int_{-\infty}^{+\infty} (\lambda - m(x+i\varepsilon))^{-1} - \int_{-\infty}^{+\infty} (\lambda - m(x-i\varepsilon))^{-1} dE_B(\lambda) \quad (17)$$

Where $z = x + i\varepsilon$ and $\bar{z} = x - i\varepsilon$. The representation admits this

$$M_B(x+i\varepsilon) - M_B(x-i\varepsilon) = \int_{-\infty}^{+\infty} ((\lambda - m(x+i\varepsilon))^{-1} - (\lambda - m(x-i\varepsilon))^{-1}) dE_B(\lambda)$$

By taking the integration both sides of equation (16) which leads to the expression [5]:

$$\begin{aligned} & \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} (M_B(x+i\varepsilon) - M_B(x-i\varepsilon)) dx \\ &= \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} \int_{-\infty}^{+\infty} ((\lambda - m(x+i\varepsilon))^{-1} - (\lambda - m(x-i\varepsilon))^{-1}) d\lambda dx \\ &= \int_{-\infty}^{+\infty} \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} ((\lambda - m(x+i\varepsilon))^{-1} - (\lambda - m(x-i\varepsilon))^{-1}) dE_B(\lambda) \end{aligned}$$

Put $= \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} ((\lambda - m(x+i\varepsilon))^{-1} - (\lambda - m(x-i\varepsilon))^{-1}) dx = k_\Delta(\lambda, \delta, t)$

We get the following

$$= \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} ((M_B(x+i\varepsilon))^{-1} - (M_B(x-i\varepsilon))^{-1}) dx = \int_{-\infty}^{+\infty} k_\Delta(\lambda, \delta, t) dE_B(\lambda), \varepsilon > 0 \quad (18)$$

and

$$k_\Delta(\lambda, \delta, \varepsilon) = \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} ((\lambda - m(x+i\varepsilon))^{-1} - (\lambda - m(x-i\varepsilon))^{-1}) dE_B(\lambda) \quad (19)$$

$\lambda \in \mathbb{R}, \Delta = (a, b) \subseteq \mathbb{R}$ and $\varepsilon > 0$ with $m(z) = \overline{m(\bar{z})}$, $z \in \mathbb{C}_-$ we denote by the family of the component intervals $\Delta_\varepsilon = (a_\varepsilon, b_\varepsilon)$ of $O_\varepsilon = \mathbb{R} \setminus \text{Supp}(m)$.

Further the function $M(\cdot)$ admits an analytic continuation to \mathbb{C}_+ such that

$$m(x) = C_0 + \int_{-\infty}^{+\infty} \left(\frac{1}{t-x} - \frac{t}{1+t^2} \right) d\mu(t), x \in O_m$$

Hence the function $m(\cdot)$ restricted to \mathbb{D} is analytic. Moreover one easily verifies that for every component interval I of \mathbb{D}

$$m(x) < m(y), x < y, x, y \in \Delta$$

Therefore for every component interval I of \mathbb{D} the set $\Delta' = m(\Delta)$ is again an open interval. Thus $O'_m = m(O_m)$ is also open and the union of the sets $O'_m = m(\Delta)$ where the union is taken over all component intervals I of \mathbb{D} .

Lemma (5-1-16) [5]:

Let $m(\cdot)$ be a scalar Nevalinna function. If $\Delta = (a, b)$ is contained in a component interval I of \mathbb{D} then $C_\Delta(\mathcal{S}) = \sup_{\lambda \in \mathbb{R}, \varepsilon \in (0,1]} |k_\Delta(\lambda, \mathcal{S}, \varepsilon)| < \infty$, for each

$$\delta \in \left(0, \frac{b-a}{2}\right) \quad (20)$$

Proof: we have

$$m(x + i\varepsilon) = m(x) - \varepsilon^2 T_0(\varepsilon, x) + \bar{z} \mathcal{E}_1(\varepsilon, x), x \in O_m \quad (21)$$

Where

$$T_0(\varepsilon, x) = \int_{-\infty}^{+\infty} \frac{1}{y-x} \cdot \frac{1}{(y-x)^2 + \varepsilon^2} d_\mu(y) \quad (22)$$

and

$$T_1(\varepsilon, x) = \int_{-\infty}^{+\infty} \frac{1}{(y-x)^2 + \varepsilon^2} d_\mu(y) \quad (23)$$

using (21) and (22) we find constant $x_0(\delta), k_1(\delta)$ and $w_1(\delta)$ such that

$$|T_0(\varepsilon, x)| \leq x_0(\delta) \quad \text{and} \quad 0 < w_1(\delta) \leq T_1(t, x) \leq x_1(\delta), \quad x \in (a + \delta, b - \delta) \quad (24)$$

For $\varepsilon \in [0, 1]$ further we get from (20)

$$\begin{aligned} P(\lambda, x, \varepsilon) &= \frac{1}{\lambda - m(x + i\varepsilon)} - \frac{1}{\lambda - m(x) - i\varepsilon T_1(\varepsilon, x)} \\ &= \frac{\lambda - m(x) - i\varepsilon T_1(\varepsilon, x) - \lambda + m(x + i\varepsilon)}{(\lambda - m(x + i\varepsilon))(\lambda - m(x) - i\varepsilon T_1(\varepsilon, x))} \end{aligned} \quad (25)$$

From (20) we get

$$P(\lambda, x, \varepsilon) = \frac{\varepsilon^2 T_0(\varepsilon, x)}{(\lambda - m(x + i\varepsilon))(\lambda - m(x) - i\varepsilon T_1(\varepsilon, x))}, \quad \lambda \in R, x \in O_m, \varepsilon > 0. \quad \text{Since both } m(x) \text{ and}$$

$T_0(\varepsilon, x)$ are real for $x \in O_m$ we have from (20) that $|\lambda - m(x + i\varepsilon)| \geq \varepsilon T_1(\varepsilon, x)$ and

$|\lambda - m(x) - i\varepsilon T_1(\varepsilon, x)| \geq \varepsilon T_1(\varepsilon, x), \lambda \in R$. In view of (36) these inequalities yield

$$|P(\lambda, x, \varepsilon)| \leq \frac{|T_0(t, x)|}{|T_1(t, x)|^2}, \lambda \in R, x \in O_m, \varepsilon > 0 \quad (26)$$

Combining (23) with (25) we obtain the estimate [5]:

$$|P(\lambda, x, \varepsilon)| \leq \frac{x_0(\delta)}{w_1(\delta)^2}, \lambda \in R, x \in (a + \delta, b - \delta), \varepsilon \in (0, 1] \quad (27)$$

We set [5]:

$$r_{\Delta}(\lambda, \delta, \varepsilon) = \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} \left(\frac{1}{\lambda - m(x) - i\varepsilon T_1(\varepsilon, x)} - \frac{1}{\lambda - m(x) + i\varepsilon T_1(\varepsilon, x)} \right) dx$$

for $\lambda \in R$ and $\varepsilon > 0$. By the representation

$$\begin{aligned} r_{\Delta}(\lambda, \delta, \varepsilon) &= \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} \frac{\lambda - m(x) + i\varepsilon T_1(\varepsilon, x) - \lambda - m(x) + i\varepsilon T_1(\varepsilon, x)}{(\lambda - m(x) - i\varepsilon T_1(\varepsilon, x))(\lambda - m(x) + i\varepsilon T_1(\varepsilon, x))} dx \\ &= \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} \left(\frac{2i\varepsilon T_1(\varepsilon, x)}{(\lambda - m(x))^2 + \varepsilon^2 T_1(\varepsilon, x)^2} \right) dx \\ &= \frac{1}{\pi} \int_{a+\delta}^{b-\delta} \left(\frac{\varepsilon T_1(\varepsilon, x)}{(\lambda - m(x))^2 + \varepsilon^2 T_1(\varepsilon, x)^2} \right) dx \end{aligned}$$

and the estimate (23) we obtain that $T_1(\varepsilon, x) = x_1(\delta)$ and $T_1(\varepsilon, x)^2 = w_1^2(\delta)$ put this in the above equation we get

$$r_{\Delta}(\lambda, \delta, \varepsilon) = \frac{1}{\pi} \int_{a+\delta}^{b-\delta} \frac{\varepsilon x_1(\delta)}{(\lambda - m(x))^2 + \varepsilon^2 w_1^2(\delta)} dx, \lambda \in R, \varepsilon \in (0, 1] \quad (28)$$

Form this equation

$$m(x) = C_0 + \int_{-\infty}^{+\infty} \left(\frac{1}{t-x} - \frac{t}{1+t^2} \right) d\mu(t), x \in O_m$$

The derivation $m'(x), x \in O_m$, admits the representation

$$m'(x) = \int_{-\infty}^{+\infty} \frac{1}{(1-x)^2} d\mu(t), x \in O_m \quad (29)$$

Obviously, there exist constants $w_2(\delta)$ and $x_2(\delta)$ such that

$$0 < w_2(\delta) \leq m'(x) \leq x_2(\delta), x \in (a + \delta b - \delta) \quad (30)$$

By combining the equation (27) and equation (29) where $0 < w_2(\delta) \leq m'(x)$,

$x \in (a + \delta b - \delta)$ we have the following

$$r_{\Delta}(\lambda, \delta, \varepsilon) \leq \frac{x_1(\delta)}{\pi w_2(\delta)} \int_{a+\delta}^{b-\delta} \frac{\varepsilon m'(x)}{(\lambda - m(x))^2 + \varepsilon^2 w_1^2(\delta)} dx, \lambda \in R, \varepsilon \in (0, 1].$$

Using the substitution $y = m(x)$ we derive that $\frac{dy}{dx} = m'(x)$ so $dx = \frac{dy}{m'(x)}$ in

the equation we get

$$\begin{aligned} r_{\Delta}(\lambda, \delta, \varepsilon) &\leq \frac{x_1(\delta)}{\pi w_2(\delta)} \int_{a+\delta}^{b-\delta} \frac{\varepsilon m'(x)}{(\lambda - m(x))^2 + \varepsilon^2 w_1^2(\delta)} \frac{dy}{m'(x)} \\ &\leq \frac{x_1(\delta)}{\pi w_2(\delta)} \int_{m(a+\delta)}^{m(b-\delta)} \frac{\varepsilon}{(\lambda - y)^2 + \varepsilon^2 w_1^2(\delta)} dy, \lambda \in R, \varepsilon \in (0, 1] \end{aligned}$$

Finally, we get

$$r_{\Delta}(\lambda, \delta, \varepsilon) \leq \frac{x_1}{w_1 w_2}, \lambda \in R, \varepsilon \in (0, 1] \quad (31)$$

Obviously we have

$$k_{\Delta}(\lambda, \delta, \varepsilon) = \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} (\rho(\lambda, \delta, \varepsilon) - \overline{\rho(\lambda, \delta, \varepsilon)}) dx + r_{\Delta}(\lambda, \delta, \varepsilon), \lambda \in R, \varepsilon > 0$$

Hence we find the estimate

$$|k_{\Delta}(\lambda, \delta, \varepsilon)| \leq \frac{1}{\pi} \int_{a+\delta}^{b-\delta} |\rho(\lambda, \delta, \varepsilon)| dx + r_{\Delta}(\lambda, \delta, \varepsilon), \lambda \in R, \varepsilon > 0$$

Taking into account equation $|\rho(\lambda, \delta, \varepsilon)| \leq \frac{x_0(\delta)}{w_1(\delta)^2}$ and the equation

$$r_{\Delta}(\lambda, \delta, \varepsilon) \leq \frac{x_1}{w_1 w_2} \quad \text{we arrive at the estimate} \quad |k_{\Delta}(\lambda, \delta, \varepsilon)| \leq \frac{x_0}{\pi w_1(\delta)} (b-a)$$

$$+ \frac{x_1(\delta)}{w_1(\delta) w_2(\delta)}, \lambda \in R, \varepsilon \in (0, 1] \quad . \text{ Which proves (19).}$$

Since the function φ_i is strictly monotone on each component interval I_i of I the inverse function $\varphi_i^{-1}(\cdot)$ exists there. The function $\varphi_i(\cdot)$ is analytic and also strictly monotone, its first derivative $\varphi_i'(\cdot)$ exists, it is analytic and non-negative.

Lemma (5-1-17) [131]:

Suppose that $m(\cdot)$ is a scalar Nevanlinna function, let $\Delta=(a,b)$ be contained in some component interval I of $O_m = \mathbb{R} \setminus \text{supp}(m)$, then (with λ defined as in (18)).

$$\lim_{\varepsilon \rightarrow +0} k_{\Delta}(\lambda, \delta, \varepsilon) = \theta_L(\lambda, \delta) = \begin{cases} 0 & \lambda \in \mathbb{R} \setminus [m(a+\delta), m(b-\delta)] \\ \frac{1}{2} \phi'_L(\lambda) & \lambda \in \{m(a+\delta), m(b-\delta)\} \\ \phi'_L(\lambda) & \lambda \in (m(a+\delta), m(b-\delta)) \end{cases} \quad (32)$$

For $\delta \in (0, (b-a)/2)$ and

$$\lim_{\varepsilon \rightarrow +0} \lim_{\delta \rightarrow +0} k_{\Delta}(\lambda, \delta, \varepsilon) = \theta_L(\lambda, \delta) = \begin{cases} 0 & \lambda \in \mathbb{R} \setminus (m(a), m(b)) \\ \phi'_L(\lambda) & \lambda \in (m(a), m(b)) \end{cases} \quad (33)$$

Proof [5]:

At first let us show that

$$\lim_{\varepsilon \rightarrow 0} \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} \rho(\lambda, x, \varepsilon) dx = 0, \quad \lambda \in \mathbb{R} \quad (34)$$

by (24) one immediately gets that

$$\begin{aligned} \lim_{\varepsilon \rightarrow 0} \rho(\lambda, x, \varepsilon) &= \lim_{\varepsilon \rightarrow 0} \left(\frac{1}{\lambda - m(x+i\varepsilon)} - \frac{1}{\lambda - m(x) - i\varepsilon T_1(\varepsilon, x)} \right)^{\frac{1}{i}} \\ &= \lim_{\varepsilon \rightarrow 0} \rho \left(\frac{\varepsilon^2 T_0(\varepsilon, x)}{(\lambda - m(x+i\varepsilon))(\lambda - m(x) - i\varepsilon T_1(\varepsilon, x))} \right)^{\frac{1}{i}} = 0, \quad \lambda \in \mathbb{R}, x \in O_m, \varepsilon > 0 \end{aligned}$$

Which implies that $\lim_{\varepsilon \rightarrow 0} \rho(\lambda, x, \varepsilon) = 0$ by lemma (5-1-16). Now (33) is implied by (26) and the Lebesgue dominated convergence theorem. Next we set Lebesgue

$$T_3(t, x) = \int_{-\infty}^{+\infty} \frac{1}{(y-x)^2 + \varepsilon^2} \cdot \frac{1}{(y-x)^2} d\mu(y), x \in O_m, t \geq 0 \quad (35)$$

Obviously there is a constant $x_3(\delta) > 0$ such that

$$0 \leq T_3(\varepsilon, x) \leq x_3(\delta), x \in (a+\delta, b-\delta), \varepsilon \in [0, 1] \quad (36)$$

Let

$$\rho_0(\lambda, x, t) = \frac{1}{\lambda - m(x) - i\varepsilon\tau_1(\varepsilon, x)} - \frac{1}{\lambda - m(x) - i\varepsilon T_1(0, x)}, \lambda \in R, x \in O_m \quad (37)$$

For $\varepsilon > 0$, it follows from (20), (35) and (37)

That

$$\rho_0(\lambda, x, \varepsilon) = \frac{-i\varepsilon^3\tau_3(\varepsilon, x)}{(\lambda - m(x) - i\varepsilon T_1(\varepsilon, x))(\lambda - m(x) - i\varepsilon T_1(0, x))} \quad (38)$$

for $\varepsilon > 0$, since $\lambda \in R$ and $m(x)$ is real for $x \in O_m$ we get from (38)

$$|\rho_0(\lambda, x, \varepsilon)| \leq \frac{\tau_3(\varepsilon, x)}{\tau_1(\varepsilon, x) T_1(0, x)}, \lambda \in R, x \in O_m, \varepsilon > 0 \quad \text{where}$$

$$\begin{aligned} \tau_1(\varepsilon, x) &= \lambda - m(x) - i\varepsilon\tau_1(\varepsilon, x), \\ \tau_1(0, x) &= \lambda - m(x) - i\varepsilon\tau_1(0, x), \end{aligned}$$

by using (23) and (36) we obtain the estimate [5]:

$$|\rho_0(\lambda, x, \varepsilon)| \leq \frac{\varepsilon\tau_3(\delta)}{w_1(\delta)^2}, \lambda \in R, x \in (a + \delta, b - \delta), \varepsilon \in (0, 1]$$

Which immediately yields

$$\lim_{\varepsilon \rightarrow 0} \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} \rho_0(\lambda, x, \varepsilon) dx = 0, \lambda \in R, \delta > 0 \quad (39)$$

Finally, let us introduce

$$q_\Delta(\lambda, \delta, \varepsilon) = \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} \left(\frac{1}{\lambda - m(x) - i\varepsilon\tau_1(0, x)} - \frac{1}{\lambda - m(x) + i\varepsilon\tau_1(0, x)} \right) dx \quad (40)$$

For $\lambda \in R$ and $\varepsilon > 0$. Using the representation

$$\begin{aligned} q_\Delta(\lambda, \delta, \varepsilon) &= \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} \left(\frac{(\lambda - m(x) + i\varepsilon\tau_1(0, x)) - (\lambda - m(x) - i\varepsilon\tau_1(0, x))}{(\lambda - m(x) - i\varepsilon\tau_1(0, x))(\lambda - m(x) + i\varepsilon\tau_1(0, x))} \right) dx \\ &= \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} \left(\frac{2i\varepsilon\tau_1(0, x)}{(\lambda - m(x))^2 + \varepsilon^2\tau_1(0, x)^2} \right) dx \end{aligned}$$

Form the equation (20) $\tau_1(0, x) = \int_{-\infty}^{+\infty} \frac{1}{(y-x)^2} d\mu$ and the equation

$m'(x) = \int_{-\infty}^{+\infty} \frac{1}{(t-x)} d\mu \geq (y), x \in O_m$. We get this relation $m'(x) = \tau_1(0, x), x \in O_m$ from

the equation (20) and equation (28) we get after change of variable $y = m(x)$ that

$$\begin{aligned} q_{\Delta}(\lambda, \delta, \varepsilon) &= \frac{1}{\pi} \int_{m(a+\delta)}^{(b-\delta)} \frac{\varepsilon m'(x)}{(\lambda - m(x))^2 + \varepsilon^2 \tau_1(0, x)^2} dx \\ &= \frac{1}{\pi} \int_{m(a+\delta)}^{(b-\delta)} \frac{\varepsilon m'(x)}{(\lambda - y)^2 + \varepsilon^2 \tau_1(0, \varphi_L(y))^2} \frac{dx}{m'(x)}, \lambda \in R, \varepsilon > 0 \\ &= \frac{1}{\pi} \int_{m(a+\delta)}^{m(b-\delta)} \frac{\varepsilon}{(\lambda - y)^2 + \varepsilon^2 \tau_1(0, \varphi_L(y))^2} dy \end{aligned}$$

where $x = \varphi(y)$

By $\tau_1(0, \varphi(y)) = m'(\varphi(y)) = 1/\varphi'(y), y \in \Delta$, we finally obtain that

$$q_{\Delta}(\lambda, \delta, \varepsilon) = \frac{1}{\pi} \int_{m(a+\delta)}^{m(b-\delta)} \frac{t \varphi'(y)^2}{\varphi'(y)^2 (\lambda - y)^2 + \varepsilon^2} dy, y \in R, \varepsilon > 0 \quad (41)$$

Next we prove the relation

$$\lim_{\varepsilon \rightarrow 0} q_{\Delta}(\lambda, \delta, \varepsilon) = q_L(\lambda, \delta), \delta \in (0, (b-a)/2), \lambda \in R \quad (42)$$

We consider only the case when $\lambda \in (m(a+\delta), m(b-\delta))$. The other cases can be treated in a similar way.

Noting that $\varphi'(\lambda) > 0$ choose an arbitrary $C \in (0, \varphi'(\lambda))$. Since φ is continuous

we can choose $\eta > 0$ such that $m(a+\delta) < \lambda - \eta < \lambda + \eta < m(b-a)$ and

$$0 < \varphi(\lambda) - C \leq \varphi(y) \leq \varphi(\lambda) + C, \lambda - \eta < y \leq \lambda + \eta \quad (43)$$

Let $a, b > 0$. The change of variables $x = (y - \lambda)/\varepsilon$ yields [5]:

$$\int_{\lambda-\eta}^{\lambda+\eta} \frac{a^2 \varepsilon}{b^2 (\lambda - y)^2 + \varepsilon^2} dy = \frac{a^2}{\varepsilon} \int_{-\frac{b\eta}{\varepsilon}}^{\frac{b\eta}{\varepsilon}} \frac{1}{1+x^2} \cdot \frac{\varepsilon}{b} dx \rightarrow \frac{\pi a^2}{b} \text{ as } \varepsilon \rightarrow 0 \quad (44)$$

Setting $a = \varphi'(\lambda) - C$ and $b = \varphi' - C$ in (43) and using (44) we obtain

$$\pi \frac{(\varphi'(\lambda) - C)^2}{\varphi'(\lambda) + C} \leq \liminf_{\varepsilon \rightarrow 0} \int_{\lambda-\eta}^{\lambda+\eta} \frac{\varepsilon \varphi'_i(y)^2}{\varphi'_i(y)^2 (\lambda - y)^2 + \varepsilon^2} dy \quad (45)$$

$$\liminf_{\varepsilon \rightarrow 0} \int_{\lambda-\eta}^{\lambda+\eta} \frac{\varepsilon \varphi'_i(y)^2}{\varphi'_i(y)^2 (\lambda - y)^2 + \varepsilon^2} dy \leq \pi \frac{(\varphi'(\lambda) - C)^2}{\varphi'(\lambda) + C}$$

Setting $G = (m(a + \delta), m(b - a)) \setminus (\lambda - \eta, \lambda + \eta)$ and applying the Lebesgue dominated convergence theorem we get

$$\lim_{\varepsilon \rightarrow 0} \int_G \frac{\varepsilon \varphi'_i(y)^2}{\varphi'_i(y)^2 (\lambda - y)^2 + \varepsilon^2} dy = 0 \quad (45)$$

By (44) and (45)

$$\begin{aligned} \pi \frac{(\varphi'_i(\lambda) - C)^2}{\varphi'_i(\lambda) + C} &\leq \liminf_{\varepsilon \rightarrow 0} \int_{m(a+\delta)}^{m(b-\delta)} \frac{\varepsilon \varphi'_i(y)^2}{\varphi'_i(y)^2 (\lambda - y)^2 + \varepsilon^2} dy \\ &\leq \liminf_{\varepsilon \rightarrow 0} \int_{m(a+\delta)}^{m(b-\delta)} \frac{\varepsilon \varphi'_i(y)^2}{\varphi'_i(y)^2 (\lambda - y)^2 + \varepsilon^2} dy \leq \frac{(\varphi'_i(\lambda) + c)^2}{\varphi'_i(\lambda) - c} \end{aligned} \quad (46)$$

Since (46) holds for every $C \in (0, \varphi'(\lambda))$, (46) in combination with (40) imply (41) combining (18), (26), (36) and (39) we derive the representation

$$\begin{aligned} k_{\Delta}(\lambda, \delta, \varepsilon) &= \frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} \left(\rho(\lambda, x, \varepsilon) - \overline{\rho(\lambda, x, \varepsilon)} \right) + \\ &\frac{1}{2\pi i} \int_{a+\delta}^{b-\delta} \left(\rho_0(\lambda, x, \varepsilon) - \overline{\rho_0(\lambda, x, \varepsilon)} \right) + q_{\Delta}(\lambda, x, \varepsilon) \end{aligned} \quad (47)$$

Where $\lambda \in \mathbb{R}$ and $\varepsilon > 0$. Now combining the relation (33), (38) and (41) with (37) we arrive at (41). The relation (32) immediately follows from (31). Now

we are ready to calculate a non-orthogonal spectral measure $\Sigma_{\mathbb{B}}^0$ in a gap of any self-adjoint extension $A_{\mathbb{B}} = A_{\mathbb{B}}^* \in \mathcal{E}_{\mathcal{H}_1, \mathcal{H}_2}$ if only \mathcal{H}_1 admits a boundary triple of a scalar-type Weyl function.

Section (5-2): Gaps and Examples

The symmetric operator S admits a boundary triple $\Pi = \{H, \Gamma_0, \Gamma_1\}$ and is of scalar-type. On the spectrum $\sigma(A_B)$ of the operator A_B , outside the gaps $O_m^c = \mathbb{R} \setminus O_m = \text{Supp}(m)$. We obtain results on the absolutely continuous spectrum [116,146].

Theorem (5-2-1) [131]:

Let $m(\cdot)$ be a scalar Nevanlinna function in \mathcal{N} with the integral representation $F(z) = C_0 + C_1 z + \int_{-\infty}^{+\infty} \left(\frac{1}{t-z} - \frac{1}{1+t^2} \right) d \sum_F(t), z \in C_+ \cup C_-$ and the imaginary part $\nu(z) = S(m(z))$ with admits the representation

$$\nu(x, y) = C_1 y + \int_{\mathbb{R}} \frac{y d\mu(t)}{(t-x)^2 + y^2}, \int_{\mathbb{R}} \frac{d\mu(t)}{1+t^2} < \infty$$

Where $\nu(x, y) = \nu(x + iy), z = x + iy \in C_+$. Then

(i) For any $x \in \mathbb{R}$ the $\lim_{t \rightarrow 0} \nu(x + it) = \lim_{t \rightarrow 0} \nu(x + iy)$ exists and is finite if and

$$\text{only if the symmetric derivation } D_\mu(x) = \lim_{\varepsilon \rightarrow 0} \frac{\mu(x + \varepsilon) - \mu(x - \varepsilon)}{2\varepsilon}$$

exists and is finite. In case one has $\nu(x + i0) = \pi D_\mu(x)$.

(ii) If the symmetric derivative $D_\mu(x)$ exists and is infinite, then

$$\nu(z) \rightarrow +\infty \text{ as } z \rightarrow x + i0.$$

(iii) For each $x \in \mathbb{R}$ one has $S m(z - x) \nu(z) \rightarrow \nu(\{x\})$ as $z \rightarrow x$.

(iv) $\nu(z)$ converges to a finite constant as $z \rightarrow x$ if and only if the

$$\text{derivative } \mu'(t) = \frac{d\mu(t)}{dt} \text{ exists at } t = x \text{ and is finite.}$$

Moreover, one has $\nu(x_0 + i0) = \pi \mathcal{L}(x)$. The symbol \rightarrow means that the limit

$$\lim_{r \rightarrow 0} \nu(x + r e^{i\theta}), x \in \mathbb{R} \text{ exist uniformly in } \theta \in [\varepsilon, \pi - \varepsilon] \text{ for each } \varepsilon \in (0, \pi/2).$$

Theorem (5-2-2) [131]:

Let T be a simple symmetric operator in \mathcal{H} with infinite deficiency indices. Further, let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for T with scalar-type Weyl function $M(\cdot)$ i.e., $M(z) = m(z)I_H$ and let $B = B^* \in C(H)$.

- (i) Then $\mathfrak{S}_{ac}(A_B) \supseteq \mathfrak{S}_{ac}(A_0)$, $A_0 = A^* / \ker(\Gamma_0)$.
- (ii) If the operator T is purely absolutely continuous, then the self-adjoint extension A_B is purely absolutely continuous to T .

Proof:

By corollary (5-1-14) we get that $\mathfrak{S}_{ac}(A_0) = \text{Supp}_{ac}(\mu)$ where μ is the random measure of the representation

$$\Omega_{ac}(m) = \left[x \in R : \exists m(x+io) = \lim_{y \rightarrow 0} m(x+iy) \text{ and } 0 < \nu(x, 0) = \int m(x+io) < \infty \right]$$

Notice that the limit $m(x+io) = \lim_{y \rightarrow 0} m(x+iy)$ exists for almost all $x \in R$.

Further, let us introduce the set $cl_{ac}(x) = \{x \in R : \text{mes}(x - \varepsilon, x + \varepsilon) > 0 \text{ for all } \varepsilon > 0\}$. We get $cl_{ac}(\Omega_{ac}(m)) = \text{Supp}_{ac}(\mu)$. By remark (5-1-15) the

Weyl function $M_B(\cdot)$ of the extension A_B is given by $M_B(z) = (B - m(z))^{-1}$

$= (B - m(z)I_H)^{-1}$, $z \in C_+$. Let us introduce the scalar-function

$$\begin{aligned} M_{B,h}(z) &= (M_B(z)h, h) = ((B - m(z)I_H)^{-1}h, h) \\ &= \int_R \frac{d(E_B(t)h, h)}{t - m(z)}, z \in C_+ \end{aligned} \quad (48)$$

For $h \in H$. If $z = x + iy$ and $m(z) = u(x, y) + i v(x, y)$, then we get from (48)

$$F_{B,n}(z) = S(m_{B,h}(z)) = \int_R \frac{v(x, y) d(E_B(t)h, h)}{(t - u(x, y))^2 + v(x, y)^2} \quad (49)$$

Let $x \in \Omega_{ac}(m)$. Notice the limits $v(x, 0) = \lim_{y \rightarrow 0} v(x, y) > 0$ and $u(x, 0) = \lim_{y \rightarrow 0} u(x, y)$

exists if $x \in \Omega_{ac}(m)$. If $y_0 > 0$ is small enough, then

$$\frac{v(x, y)}{(t - u(x, y))^2 + v^2(x, y)} \leq \frac{1}{v(x, y)} \leq \frac{2}{v(x, 0)}, y \in [0, y_0], x \in \Omega_{ac}(m) \quad (50)$$

Taking in to account (50) and applying the Lebesgue dominated convergence theorem we obtain from (48) that

$$\begin{aligned} F_{B,h}(x + i0) &= \lim_{y \rightarrow 0} F_{B,h}(x, y) = \\ v(x, 0) \int_R \frac{d(E_B(t)h, h)}{(t - u(x, 0))^2 + v(x, 0)^2}, x &\in \Omega_{ac}(m) \end{aligned} \quad (51)$$

Since $v(x, 0) > 0$ for $x \in \Omega_{ac}(m)$ we find

$$0 < F_{B,h}(x, i0) < \infty, x \in \Omega_{ac}(m)$$

Furthermore we have

$$G_{B,h}(z) = \operatorname{Re}(M_{B,h}(z)) = \int_R \frac{(t - u(x, y)) d(E_B(t)h, h)}{(t - u(x, y))^2 + v(x, y)^2}$$

$$\text{Since } \frac{|t - u(x, y)|}{(t - u(x, y))^2 + v(x, y)^2} \leq \frac{1}{\sqrt{(t - u(x, y))^2 + v(x, y)^2}} \leq \frac{\sqrt{2}}{v(x, 0)}$$

For $x \in \Omega_{ac}(m)$ and $y \in (0, y_0)$. A gain by the Lebesgue dominated convergence theorem we find

$$G_{B,h}(x + i0) = \lim_{y \rightarrow 0} G_{B,h}(x + iy) = \int_R \frac{(t - u(x, 0)) d(E_B(t)h, h)}{\sqrt{(t - u(x, y))^2 + v(x, 0)^2}} \leq \frac{\sqrt{2}}{v(x, 0)}$$

Which shows that $x \in \Omega_{ac}(m)$ implies $x \in \Omega_{ac}(M_{B,h})$ for every $h \in H$ where

$$\Omega_{ac}(M_{B,h}) = \left\{ x \in R : \exists M_{B,h}(x + i0) = \lim_{y \rightarrow 0} M_{B,h}(x + iy) \text{ and } 0 < \operatorname{Im}(M_{B,h}(x + i0)) < \infty \right\}$$

Since $\Omega_{ac}(m) \subseteq \Omega_{ac}(M_{B,h})$ one gets $\mathcal{E}_{ac}(A_0) \subseteq \operatorname{Supp}_{ac}(\mathcal{A}) \subseteq \mathcal{I}_{ac}(\Omega_{ac}(m))$

$$\subseteq \mathcal{I}_{ac}(\Omega_{ac}(M_{B,h})) \text{ for each } h \in H.$$

If $B = B^\infty$ then the measure $\rho_h(\cdot) = (E_B(\cdot)h, h)$ is absolutely continuous for any $h \in H$, that is $d\rho_h(t) = \rho'_h(t)dt$, where $\rho'_h(\cdot) \in L^1(R)$ for any $h \in H$. One rewrite (48) as

$$F_{B,h}(z) = \int_R \frac{v(x, y) \rho'_h(t) dt}{(t - u(x, y))^2 + v(x + y)} \quad (52)$$

and the subset $H_\infty = \{h \in H : \rho'_h \in L^\infty(R) \cap L^1(R)\}$ is dense in $H = H^\infty(B)$. For

$h \in H_\infty$ we obtain from (48) that

$$C_\infty(h) = \sup_{0 < y < 1} \sup_{x \in R} \Im F_{B,h}(x + iy) \leq \|\rho'_h\|_{L^\infty} \sup_{v > 0} \sup_{u \in R} \int_R \frac{v ds}{(s - u)^2 + v^2} \quad (53)$$

Corollary (5-2-3) [131]:

Let \mathcal{H} be a simple symmetric operator with infinite deficiency indices. Further, let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ a boundary triple for \mathcal{H} with scalar-type Weyl function $M(\cdot)$. If $A \notin E_{\pi_A} A$, then $\delta_\infty(A_0) \subseteq \delta_\infty(A)$.

Corollary (5-2-4) [131]:

Shows that under the assumption of a scalar-type Weyl function the absolutely continuous spectrum of any extension always contains $\delta_\infty(A_0)$. The above result implies the following corollary.

Corollary (5-2-5) [131]:

Let \mathcal{H} be a simple symmetric operator with infinite deficiency indices on the separable Hilbert space \mathcal{H} . Further, let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for \mathcal{H} with scalar-type Weyl function $M(\cdot) = m(\cdot)I_H$ which is monotone with respect to the open set $J \subseteq O_m \subset \rho(A_0)$. Then for any operator \mathcal{A} on some separable Hilbert space there is a self-adjoint extension \mathcal{A}_J such that

$\mathcal{A}_J \cong R_J^\infty$ and \mathcal{A}_J is absolutely continuous [111,130].

Theorem (5-2-6) [131]:

Let T be a simple symmetric operator in \mathcal{H} with infinite deficiency indices. Further, let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for T with scalar-type Weyl function $M(\cdot)$. i.e., $M(z) = m(z)I_H$ and let $B = B^* \in C(H)$

- (i) If T is singular i.e., $B^* = B$, then the absolutely continuous parts $\mathcal{E}_B^{\text{ac}}$ and $\mathcal{E}_0^{\text{ac}}$ is unitarily equivalent, in particular $\mathcal{E}_{ac}(A_B) = \mathcal{E}_{ac}(A_0)$.
- (ii) If T and T_0 are singular, then T is singular.
- (iii) If T is pure point and the spectrum of T consist of isolated eigenvalues, then T is pure point.

Proposition (5-2-7) [131]:

Let T be a simple symmetric operator in \mathcal{H} with infinite deficiency indices. Further, let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple for T with scalar-type Weyl function $M(\cdot)$, i.e., $M(z) = m(z)I_H$, and $\text{Supp}^+(\mu) = \{x \in \text{supp } \mu : \exists$

$D_\mu(x) \text{ and } D_\mu(x) > 0\}$ where μ is the radon measure of representation

$$m(z) = C_0 + C_1 z + \int_{-\infty}^{\infty} \left(\frac{1}{t-z} - \frac{t}{1+t^2} \right) d\mu(t).$$

If $B \in C(H)$, then

$$E_{A_B}^T(\text{Supp}^+(\mu)) = 0, T \in \{s, pp, sc\} \quad (54)$$

In particular, it holds

- (i) $\mathcal{E}_B^{\text{ac}}(A_B) \upharpoonright_{\text{supp } \mu(\mu)} \leq \text{Supp}(\mu) \setminus \text{supp}^+(\mu)$ and
- (ii) $E_{A_B}^{\text{ac}}(\text{supp } \mu(\mu)) = 0$ provided $\text{supp}(\mu) \setminus \text{supp}^+(\mu)$ is either finite or countable.

Proof:

We set $\text{Supp}_\infty^+(\mu) = \{x \in \text{supp}^+(\mu) : D_\mu(x) = \infty\}$. By theorem (5-2-1) we

derive that the limit $\lim_{y \rightarrow 0} (v(x, y))$ exists and is finite for

$x \in \text{supp } \mu^+(\mu) \setminus \text{supp } \mu_\infty^+(\mu)$ and

$$v(x, 0) = \lim_{y \rightarrow 0} v(x, y) = \mathcal{D}_\mu(x) > 0, \quad x \in \sup p^+(\mu) \setminus \sup p_\infty^+(\mu) \quad (55)$$

by proposition (5-1-5) there exists an operator $B = B^* \in C(H)$ such that

$\mathcal{A} \ominus A_B = A^* \upharpoonright_{\ker(\Gamma_1 - B \Gamma_0)}$. We consider the generalized Weyl function

$M_B(z) = (B - M(z))^{-1}$ and define $F_{B,h}$ by (48). Following the line of reasoning of theorem (5-2-12) we obtain

$$0 < F_{B,h}(x + i0) < \infty, \quad x \in \sup p^+(\mu) \setminus \sup p_\infty^+(\mu), \quad h \in H \quad (56)$$

Further, let $x \in \sup p_\infty^+(\mu)$. By theorem (5-2-1) (ii) and (iii) we find

$$v(x, 0) = \lim_{y \rightarrow 0} v(x, y) = \infty \quad \text{and} \quad \lim_{y \rightarrow 0} y v(x, y) = \mu(\{x\}).$$

Therefore for every $y_0 > 0$ there exists $N = N(y_0)$ such that $v(x, y) \geq N$ for

$y \in (0, y_0)$. Hence

$$\frac{v(x, y)}{(t - u(x_0, y))^2 + v^2(x, y)} \leq \frac{1}{N}$$

By Lebesgue dominated theorem we obtain from (48) that

$$\lim_{y \rightarrow 0} F_{B,h}(x, iy) = 0, \quad x \in \sup p_\infty^+(\mu), \quad h \in H \quad (57)$$

Let $\sum_B(\cdot)$ be the unbounded non-orthogonal spectral measure of the Weyl

function $M_B(z) = (B - M(z))^{-1}, z \in C_+$ and $\sum_{B,h}(\cdot) = \left(\sum_B(\cdot) h, h \right), h \in H$. If

$$\delta_s \left(\sum_{B,h} \right) = \{x \in R : F_{B,h}(z) \rightarrow \infty \text{ as } z \rightarrow x\}, \quad h \in H$$

then we find from (55) and (56) that $\delta_s \left(\sum_{B,h} \right) \perp \sup p^+(\mu) = 0$. Let $\tau = \{h_k\}_{k=1}^\infty$ be

a total set in H . Setting $\delta_s \left(\sum_B, \tau \right) = \bigcup_{k=1}^\infty \delta_s \left(\sum_{B, h_k} \right)$. One gets $\delta_s \left(\sum_B, \tau \right) \perp \sup p(\mu^+) = 0$,

and we get

$$E_{A_B}^{sc}(\sup p^+(\mu)) = E_{A_B}(\sup p^+(\mu) \upharpoonright \mathcal{S}_s(\sum_B, \tau)) = 0$$

Which proves (53) for $\tau = s$. Similarly, setting

$$\mathcal{S}_{pp}(\sum_{B,h}) = \{x \in \mathbb{R} : \lim_{z \rightarrow \infty} (z - x) F_{B,h}(z) > 0\}, h \in H$$

and

$$\mathcal{S}_{pp}(\sum_B, \tau) = \bigcup_{k=1}^{\infty} \mathcal{S}_{pp}(\sum_{B,h_k})$$

We verify $\mathcal{S}_{pp}(\sum_B, \tau) \subseteq \mathcal{S}_s(\sum_B, \tau)$, one proves (53) for $\tau = pp$. Finally, setting

$$\mathcal{S}_{sc}(\sum_{B,h}) = \{x \in \mathbb{R} : F_{B,h}(z) \rightarrow \infty \text{ and } (z - x) F_{B,h}(z) \rightarrow 0 \text{ as } z \rightarrow \infty\}, h \in H, \text{ and}$$

$$\mathcal{S}_{sc}(\sum_B, \tau) = \bigcup_{k=1}^{\infty} \mathcal{S}_{sc}(\sum_{B,h_k}) \setminus \mathcal{S}_{pp}(\sum_B, \tau)$$

We obtain $\mathcal{S}_{sc}(\sum_{B,h}) \subseteq \mathcal{S}_s(\sum_{B,h})$ which yield (53) for $\tau = sc$.

(i) We have $\mathcal{S}_p(A_B) = \mathcal{S}_{pp}(\sum_B, \tau)$ which yields $\mathcal{S}_p(A_B) \upharpoonright \text{supp}(\mu) \subseteq \text{supp}(\mu) \setminus \text{supp}^+(\mu)$.

$$\subseteq \text{supp}(\mu) \setminus \text{supp}^+(\mu).$$

(ii) We have $E_{A_B}^{sc}(\text{supp}(\mu)) = E_{A_B}^{sc}(\text{supp}^+(\mu)) + E_{A_B}^{sc}(\text{supp}(\mu) \setminus \text{supp}^+(\mu))$

$$= E_{A_B}^{sc}(\text{supp}(\mu) \setminus \text{supp}^+(\mu)).$$

Since by assumption $\text{supp}(\mu) \setminus \text{supp}^+(\mu)$ is countable we obtain

$$E_{A_B}^{sc}(\text{supp}(\mu) \setminus \text{supp}^+(\mu)) = 0, \text{ which shows } E_{A_B}^{sc}(\text{supp}(\mu)) = 0.$$

We consider several examples in order to illustrate the previous results [148].

Example (5-2-8) [131]:

Let $h = L^2((0,1))$. By [131] we denote the closed symmetric operator

$$(Af)(x) =$$

$-i \frac{d}{dx} f(x), x \in (0,1), f \in \text{dom}(A) = \{f \in W_2^1((0,1)) : f(0) = f(1) = 0\}$. Which is simple

and has deficiency indices $(1,1)$. We note that A^* is given by $(A^*f)(x) =$

$-i \frac{d}{dx} f(x), f \in \text{dom}(A^*) = W_2^1((0,1))$. A straight forward computation shows that

$\Gamma = \{\Gamma_0, \Gamma_1\}$ where $H = \mathbb{C}$

$$\Gamma_0 f = \frac{f(0) - f(1)}{\sqrt{2}}, \Gamma_1 f = i \frac{f(0) + f(1)}{\sqrt{2}}, f \in \text{dom}(A^*) = W_2^1((0,1)) \quad (58)$$

forms a boundary triple for A_0 . The operator $A_0 = A^* \upharpoonright_{\ker(\Gamma_0)}$ is given by

$$(A_0 f)(x) = -i \frac{d}{dx} f(x), x \in (0,1), f \in \text{dom}(A_0) \\ = \{W_2^1((0,1)) : f(0) = f(1)\}$$

The spectrum of A_0 is discrete. It consist of isolated eigenvalues we have

$\sigma(A_0) = \{\lambda_l\}_{l \in \mathbb{Z}}$ with $\lambda_l = 2L\pi$. Obviously we have $\rho(A_0)U_{L \in \mathbb{Z}} \Delta_L$ where

$$\Delta_L = (2L\pi, 2(L+1)\pi).$$

Trivially the open intervals Δ_L are gaps of the operator $A_0 = A_0^*$. Hence they are gaps of the symmetric operator A_0 . The extension $A_1 = A^* \upharpoonright_{\ker(\Gamma_1)}$

has the domain $(A_1), \text{dom}(A_1) = \{f \in W^{(1,2)}((0,1)) : f(0) = -f(1)\}$. Its spectrum is

discrete and consists of the eigenvalues $\lambda_l = (2L+1)\pi, L \in \mathbb{Z}$. Any other

extension of A_0 is given by a real constant $\theta \in \mathbb{R}$ and the boundary triple

$\Gamma_\theta = \{C, \Gamma_0^\theta, \Gamma_1^\theta\}$, where $\Gamma_1^\theta = \Gamma_0$ and $\Gamma_0^\theta = \theta \Gamma_0 - \Gamma_1$. The domain $\text{dom}(A_\theta)$ of the

self-adjoint extension $A_\theta = A^* \upharpoonright_{\ker(\Gamma_\theta^\theta)}$ can be alternatively described by

$$\text{dom}(A_\theta) = \{f \in W_2^1((0,1)) : (\theta - i)(\theta + i)^{-1} f(0) = f(1)\}$$

the spectrum of A_θ is also describe and consists of eigenvalues. Setting

$$\theta = -\cot(T/2), \quad T \in (0, 2\pi) \quad \text{one easily verifies that} \quad \lambda_l^{(\theta)} = T + L, L \in \mathbb{Z}.$$

In other words any extension of \mathcal{H} , which is different from \mathcal{H} , has an eigenvalues in the gaps $\Delta_L, L \in \mathbb{N}$. i.e., it does not preserve the gaps Δ_L . It is easily seen that the Weyl function corresponding to the boundary triple

$\Pi = \{H, \Gamma_0, \Gamma_1\}$ other form (57) is

$$m(z) = -\frac{\cos(z/2)}{\sin(z/2)} = -\cot(z/2), z \in C_+ \cup C_-$$

The open set $O_m = \mathbb{R} \setminus \sup p(m)$ coincides with $P(A_0)^\perp$. i.e., $O_m = U_{L \in \mathbb{N}} \Delta_L$.

The Weyl function admits an extension to \mathbb{R} which is given by

$m(\lambda) = -\cot(\lambda/2), \lambda \in O_m$. Obviously the Weyl function $m(\cdot)$ is increasing on

each open interval Δ_L . However, choosing $J = O_m$ one easily verifies that the

Weyl function $m(\cdot)$ is not monotone with respect to \mathbb{R} . The lack of

monotonicity is related to the fact that there does not exist an extension \tilde{m} of

m which has only an eigenvalue in one gap Δ_L as we have seen above.

Let us consider the closed symmetric operator $S = \bigoplus_{k=1}^{\infty} S_k$ on the Hilbert space $\mathcal{H} = \bigoplus_{k=1}^{\infty} \mathcal{H}_k$ where the operators S_k are unitarily equivalent to S defined above. Obviously the operator S is unitarily equivalent to the operator \tilde{S} defined on $\mathcal{H} = L^2((0, \infty))$

$$(Cf)(x) = -i \frac{d}{dx} f(x), f \in \text{dom}(C) = \{W_2^1(\mathbb{R}_+) : f(k) = 0, k \in \{0\} \cup \mathbb{N}\}.$$

We note that $O_m = U_{L \in \mathbb{N}} (2\pi L, 2\pi(L+1))$ and $\varphi(t) = -2 \arccot(t) + 2\pi(L+1), L \in \mathbb{N}$.

By the associated non-orthogonal spectral measure $\Sigma_B^0(\cdot)$ and $\Sigma_B^1(\cdot)$ of the

Weyl function $M_B(z) = (B - m(z), I)^{-1}$ are given by

$$\begin{aligned} \Sigma_B^0(\delta) &= \varphi'_L(B) \left(1 + \varphi_L(B)^2\right)^{-1} E_B(m(\delta)) = \\ &= \left(1 + 2\pi(L+1) - 2 \arccot(B)^2\right)^{-1} E_B(-\cot(\delta/2)) \end{aligned} \quad (59)$$

and

$$\sum_{\mathcal{S} \in B(\Delta_L)} (\mathcal{S}) = \varphi^1(B) E_B(-\cot(\mathcal{S}/2)) = 2(1+B^2)^{-2} E_B(-\cot(\mathcal{S}/2)) \quad (60)$$

It follows from (59) that the measure $\sum_B(\cdot)$ is periodic:

$$\sum_{\mathcal{S} \in B(\Delta_L)} (\mathcal{S} + 2\pi L) = \sum_{\mathcal{S} \in B(\Delta_L)} (\mathcal{S}), \quad L \in \mathbb{C}$$

Having in mind this fact one obtains that for any $L \in \mathbb{C}$ the operator

$$\mathcal{S}_{\mathcal{B}E_{s_B}}(2\pi L, 2\pi(L+1)) \quad \text{is unitarily equivalent to the operator} \quad \mathcal{S}_{\mathcal{B}E_{s_B}}((0, 2\pi)) \quad .$$

Example (5-2-9) [131]:

Let $\mathfrak{h} = L^2(\mathfrak{i}_+)$ and let \mathfrak{s} be a closed symmetric operator in \mathfrak{h} defined by

$$\begin{aligned} (\mathfrak{s}f)(x) &= -\frac{d^2}{dx^2} f(x), f \in \text{dom}(\mathfrak{s}_1) = W_2^2(\mathfrak{i}_+) \\ &= \{f \in W_2^2(\mathfrak{i}_+) : f(0) = f'(0) = 0\} \end{aligned} \quad (61)$$

Obviously $\delta_1 \geq 0$. Setting

$$\begin{aligned} \Gamma_0^*(\theta)f &= f'(0) - \theta f(0), \Gamma_1^*(\theta)f = -f(0), f \in \text{dom}(\mathfrak{s}_1^*) = \\ &= W_2^2(\mathfrak{i}_+) \quad 0 \in \mathfrak{i} \end{aligned}$$

We obtain the boundary triple $\Gamma \mathfrak{F} = \{\mathfrak{F}, \Gamma_0^*(\theta), \Gamma_1^*(\theta)\}$. It is clear that the

extension \mathfrak{s}_θ is non-negative if $\theta \geq 0$. The corresponding Weyl function is

$$m_\theta(\lambda) = (\theta - i\sqrt{\lambda})^{-1}. \quad \text{It is regular in } \mathfrak{E} \setminus \mathfrak{i}_+ \quad \text{if } \theta \geq 0, \quad \text{where the branch of } \sqrt{\cdot} \quad \text{is}$$

fixed by the condition $\sqrt{1} = 1$. The Weyl function $m_\theta(\cdot)$ admits the following integral representation

$$m_\theta(\lambda) = (\theta - i\sqrt{\lambda})^{-1} = \frac{1}{\pi} \int_0^\infty \frac{\sqrt{t}}{(t-z)(t+\theta^2)} dt, \quad \theta \geq 0$$

and the corresponding spectral measure is given by $d\mu_\theta = \pi^{-1} t^{\frac{1}{2}} (t+\theta^2)^{-1} dt$.

Clearly $m_\theta(\cdot)$ is holomorphic within $(-\infty, 0)$ such that $m_\theta(-\infty) = (0, \theta^2)$. The

inverse function $\mathfrak{M}(\cdot) : (0, \theta^2) \rightarrow (-\infty, 0)$ is given by $\mathfrak{M}(\zeta) = -(\zeta^{-1} - \theta^2)^2$,

$\zeta \in (0, \theta^{-1})$. We set $\Delta = (-\infty, 0)$ and $\Delta' = m_\theta(\Delta) = (0, \theta^{-1})$. Notice that

$$\varphi(\zeta) = 2(\zeta^{-1} - \theta)\zeta^{-2}.$$

Let $h = \bigoplus_{k=1}^\infty h_k$, $A = \bigoplus_{k=1}^\infty A_k$ and $\Gamma = \{H, \Gamma_0, \Gamma_1\} = \bigoplus_{k=1}^\infty \Gamma_k$ where $h_k = h$,

function $M(\cdot)$ is of scalar-type, i.e., $M(z) = m_\theta(z)I_H$. Further, let

$B = B^* \in C(H)$. To the self-adjoint extension \mathfrak{A} it corresponding the Weyl

function $M_B(z) = (B - m_\theta(z)I_H)^{-1}$. Let $\Sigma_B(\cdot)$ be the unbounded non-

orthogonal spectral measure of the Weyl function $M_B(\cdot)$. It follows from (60)

$$\sum_B(\mathcal{S}) = 2(B_{\Delta'}^{-1} - \theta)B_{\Delta'}^{-2}E_B(m_\theta(\mathcal{S}), B_{m_\theta(\Delta)} = B_{\Delta'} = BE_B(\Delta), \mathcal{S} \in B(\Delta) \quad (62)$$

Let $\mathcal{S} = (x, 0)$, $x < 0$. Since $m_\theta((x, 0)) = \left((\theta + \sqrt{|x|})^{-1}, \theta^{-1} \right)$ for $x < 0$ we get from (62)

$$\sum_B((x, 0)) = 2(B_{\Delta'}^{-1} - \theta)B_{\Delta'}^{-2}E_B\left(\left(\theta + \sqrt{|x|}\right)^{-1}, \theta^{-1}\right), x < 0 \quad (63)$$

we note that $\sum_B(x) \in [H]$ for every $x < 0$, while $B_{\Delta'}^{-1}$ may be unbounded.

Further, starting with (50) we can explicitly calculate the non-orthogonal

spectral measure $\sum_B(\cdot)$ outside the gap $\Delta = (-\infty, 0) = \left(\sum_B(\cdot)h, h \right)$. Setting $\sum_{B,h}(\cdot)$

and $F_{B,h}(z) = 3m(M_B(z)h, h)$ we easily derive from (51) and the Fatous theorem that

$$\pi \frac{d \sum_{B,h}(x)}{dx} = F_{B,h}(x + i0) = \int_R \frac{\sqrt{x} d(E_B(t)h, h)}{(t\theta - 1)^2 + xt^2}, x > 0, h \in H \quad (64)$$

Where $\sum_{B,h}(x) = \sum_{B,h}((0, x))$, $x > 0$. straight forward computation shows that

$Supp^+(\mu_0) = (0, \infty)$. By proposition (5-2-7) we have $E_{AB}^T((0, \infty)) = 0, T = s, pp, sc$. Hence

$\mathfrak{A}(A_B) \subseteq (-\infty, 0], T = s, p, sc$. Since $\mathfrak{A}_\infty(A_0) = [0, \infty)$ we obtain from theorem (5-2-2)

that $E_{AB}(\cdot) \supseteq [0, \infty)$. Therefore, the orthogonal spectral measure $E_{AB}(\cdot)$ of \mathfrak{A} is

absolutely continuous on $(0, \infty)$ which yields that $\sum_B(\cdot)$ is absolutely continuous on $(0, \infty)$. i.e., $\sum_B^*(\delta) = \sum_B(\delta)$ for $\delta \in B((0, \infty))$. Hence

$$\begin{aligned} \sum_{B,h}((0, \infty)) &= \frac{1}{\pi} \int_0^x ds \int_R \frac{\sqrt{s} d(E_B(t)h, h)}{(t\theta - 1)^2 + \delta^2} \\ &= \int_R \Phi_\theta(x, t) d(E_B(t)h, h) \quad x > 0, h \in H \end{aligned} \quad (65)$$

Where $\Phi_\theta(x, t) = \frac{2}{\pi^2} \left(\sqrt{x} - \frac{|t\theta - 1|}{t} \arctan \left(\frac{t\sqrt{x}}{t\theta - 1} \right) \right)$ $x > 0$. Which yields

$$\sum_B((0, x))h = \int_R \Phi_\theta(x, t) dE_B(t)h, x > 0, h \in H \quad (66)$$

Thus, formulas (62) and (65) together give the explicit form for the unbounded non-orthogonal spectral measure $\sum_B(\cdot)$ of the extension \tilde{A} .

Section (5-3): Inverse Spectral Problem for Direct Sum of Symmetric Operator

We have [30,89,91,92]

Lemma (5-3-1) [133]:

Let A be a symmetric operator in the Hilbert space H . Suppose that A has a gap (α, β) . Let M be a closed subspace of H and D a self-adjoint operator in M such that $A|_M$ is a restriction of the adjoint A^* of A and

$\mathcal{D}(M) \subset \overline{J}$. Then

$$H_M = H_{ID(H)+D(M)}^* = M \oplus G_0 \quad (67)$$

For some symmetric operator A with gap (α, β) in the Hilbert space H . In particular $A|_M$ has a self-adjoint extension $A|_M$ such that

$$H_M^0 = M_j$$

Proof:

We have only to show that there exists a symmetric operator S with gap δ such that (66) holds. In fact since δ is a gap of S there exists a self-adjoint operator S_0 in \mathcal{H} such that $S(S_0)^{1/2} = 0$ and S is an extension of S_0 . Then obviously $\mathcal{H} \oplus M \oplus G$ is a self-adjoint extension of S and satisfies $\mathcal{H} \oplus M \oplus G = M \oplus G$.

Lemma (5-3-2) [133]:

Let S be a symmetric operator in the Hilbert space \mathcal{H} . Let \mathcal{H}_0 be closed subspace. \mathcal{H}_0 and S_0 a self-adjoint operator in \mathcal{H}_0 such that S_0 is a restriction of S . Then $H_M = H_{ID(H) \rightarrow D(M)}^* = M \oplus G_0$

For some symmetric operator S in \mathcal{H}

Proof:

Let $f, g \in D(H)$ and $f, g \in D(M)$. We have

$$\begin{aligned} (H_M(f+g), f+g) &= (Hf, f) + (H^*g, g) + (Mg, g) \\ &= (f, H^*(f+g)) + (g, Hf) + (g, Mg) \\ &= (f+g, H_M(f+g)) \end{aligned}$$

This H_M is a symmetric operator in the Hilbert space \mathcal{H} . Let $f \in D(H_M)$. For

every $n \in \mathbb{Z}$ let $P_n = I_{[n, n+1)}(M) P_{H_0}$.

Since H_M is a self-adjoint operator in the Hilbert space \mathcal{H} it follows from the spectral theorem that every $n \in \mathbb{C}$ the operator P_n is an orthogonal

projection in \mathcal{H} onto the closed subspace $R(P_n)$ of \mathcal{H} .

$$R(P_n) \subset D(M), \quad n \in \mathbb{C} \quad (68)$$

$$R(P_n) \perp R(P_m), \quad n \neq m \quad (69)$$

$$\sum_{n \in \mathbb{C}} P_n = P_{H_0} \quad (70)$$

$$P_n M g = M P_n g, \quad g \in D(M), \quad n \in \mathbb{C} \quad (71)$$

Thus we have

$$\begin{aligned} (P_n H_M f, g) &= (H_M f, g) = (f, M g) \\ &= (P_n f, M g) = (M P_n f, g) \end{aligned}$$

For every $g \in R(P_n)$. In the second step we have used (67) and the facts that

H_M is symmetric and P_n a restriction of H_M . In the third step we have used (70) and in the last step a gain (67).

Thus we have

$$P_n H_M f = M P_n f, \quad n \in \mathbb{C} \quad (72)$$

Since by (71), (68) and (69)

$$(M P_n f, M P_k f) = (P_n H_M f, P_k H_M f) = 0 \quad k \neq n$$

and $\sum_{n \in \mathbb{C}} \|M P_n f\|^2 = \sum_{n \in \mathbb{C}} \|P_n H_M f\|^2 = \|P_{H_0} H_M f\|^2 < \infty$

The sequence $\left\{M \sum_{n=-N}^N P_n f\right\}_{N \in \mathbb{N}}$ converges in \mathcal{H} . Since \mathcal{H} is closed and by (69),

$\lim_{N \rightarrow \infty} \sum_{n=-N}^N P_n f = P_{H_0} f$ it follows that

$$P_{H_0} f \in D(M) \quad (73)$$

and

$$MP_{H_0} f = \lim_{N \rightarrow \infty} \sum_{n=-N}^N MP_n f = \sum_{n \in \mathbb{Z}} P_n H_M f = P_{H_0} H_M f \quad (74)$$

Here again we have used (71), (69). By (72) and (73)

$$H_M = M \oplus G_0$$

$$G_0 = H_{m/D(H_M)} \upharpoonright H_0^\perp$$

\mathcal{H} is a symmetric operator in the Hilbert space \mathcal{H} since H_q is asymmetric operator in \mathcal{H} .

Proposition (5-3-3) [133]:

Let \mathcal{H} be a symmetric operator in the separable Hilbert space \mathcal{H} suppose that \mathcal{H} has a gap (λ_-, λ_+) . Let $\{\lambda_n\}_{n=1}^N$ be a (finite or infinite) sequence in (λ_-, λ_+) [54]. Then there exists a self-adjoint extension \mathcal{H}_λ of \mathcal{H} such that

$$\mathcal{S}_p(\mathcal{H}_\lambda) \upharpoonright J = \{\lambda_n : n \in \mathbb{N}, 1 \leq n \leq N\}$$

and for every eigenvalue λ_n of \mathcal{H}_λ in (λ_-, λ_+) the multiplicity of λ_n equals the number of times it occurs in the sequence $\{\lambda_n\}_{n=1}^N$ if and only if λ_n is less than or equal to the deficiency number of \mathcal{H} . In this case \mathcal{H}_λ can be chosen such that it has a pure point spectrum inside the gap (λ_-, λ_+) .

Proof:

First we shall do the "only-if-part". Trivially the assertion of this part is true provided the deficiency number of \mathcal{H} is finite. But then the "only-if-part" follows from Krein's theorem suppose now that \mathcal{H} is less than or equal to the

deficiency number of T . Then we can choose by induction an orthonormal system $\{e_n\}_{n=1}^N$ such $e_n \in N(H^* - \lambda_n), n \in \mathbb{N}, 1 \leq n \leq N$.

Due to the well known fact that the dimension of $N(H^* - \lambda)$ equal the deficiency number of T for every regular point λ of T and in particular for every $\lambda \in J$ [107,126].

Let $H_{00} = \text{Span}\{e_n : n \in \mathbb{N}, 1 \leq n \leq N\}$ and $H_0 = \overline{H_{00}}, M_0 = H_{IH_{00}}^*$

By construction $\{e_n\}_{n=1}^N$ is an orthonormal base of the Hilbert space H_{00} and for every $n \in \mathbb{N}, 1 \leq n \leq N$ is an eigenvector of M_0 corresponding to the real eigenvalue λ_n . Thus M_0 can be and will be regarded as an operator in the Hilbert space H_{00} its closure M is a self-adjoint operator in H_0 has a pure point spectrum $\sigma_p(M) = \{\lambda_n : n \in \mathbb{N}, 1 \leq n \leq N\}$ and for every eigenvalue λ_n of M the multiplicity of λ_n equals the number of times it occurs in the sequence $\{\lambda_n\}_{n=1}^N$. $M|_{H_{00}}$ is a restriction of $T|_{H_{00}}$ since the adjoint of any operator closed. Thus $M|_{H_{00}}$ has a self-adjoint extension M_1 such that $M_1^0 = M$. i.e., such that M_1 has the required properties.

Definition (5-3-4) [133]:

A symmetric operator T is significantly deficient if and only if it has a real regular point and $P_{N(H^* - z)} D(\tilde{H}) \neq N(H^* - z)$. For every regular point λ of T .

Proposition (5-3-5) [133]:

(i) Let T be a closed symmetric operator in the Hilbert space H . Let

λ be a gap of T and $0 \in J$. Let $A = P_{R(H)} H^{-1}, B = P_{R(H)^\perp} H^{-1}$, and P an

orthogonal projection in $R(H)^\perp$ such that the operator PA belongs to the trace class and let 0 be the zero-operator in the Hilbert space

$R(1-P)$, then for every invertible self-adjoint operator T in the

Hilbert space $R(P)$ the operator $L = \begin{pmatrix} A & (PB^*) \\ PB & Q \end{pmatrix} : \begin{matrix} R(H) \\ D(Q) \end{matrix} \rightarrow \begin{matrix} R(H) \\ R(P) \end{matrix}$, is

invertible and the operator $H_0 = 0 \oplus P_0$ is a self-adjoint extension of H such that $H_0^{-1} : Q_{ac,J}^{-1} \rightarrow Q_{ac,J}$.

- (ii) There exists an orthogonal projection E in the Hilbert space $R(H)^\perp$ such that the operator TE belongs to the trace class and $R(P)$ is infinite dimensional if and only if the operator T is significantly deficient in the sense of the definition (5-3-4).

Corollary (5-3-6) [133]:

Let T be a significantly deficient symmetric operator and let δ be a gap of T . Then for every self-adjoint operator M in a separable Hilbert space there exists a self-adjoint extension \tilde{T} of T such that

$$H_{ac,J}^{-1} : M_{ac,J}' \rightarrow M_{ac,J}'.$$

Definition (5-3-7) [133]:

A symmetric operator T is weakly significantly deficient if there exists a real regular point λ of T and a real number ν which is not an eigenvalue of the operator $A = P_{R(H^* - \nu)}(H_0 - \nu)^{-1}$. Such that $R(B(A - \lambda)) \neq N(H^* - \nu)$. Where

$$B = P_{N(H^* - \nu)}(H_0 - \nu)^{-1}. \text{ We may assume that the operator } T \text{ is closed and } \nu \neq 0.$$

$R(A - \lambda)$ is dense in the Hilbert space $R(H)$ since T is self-adjoint and λ is not an eigenvalue of T . Since T is bounded and $R(B)$ is dense in $R(H)^\perp$ this implies that $R(B(A - \lambda))$ is dense in $R(H)^\perp$. Thus we can replace T by $B(A - \lambda)$ in the consideration at the beginning and get that there exists an

orthogonal projection P in the Hilbert space $R(H)^\perp$ such that $\dim R(P) = \infty$.

And the operator $PB(A - \lambda)$ belongs to the trace class. Let 0 be the zero operator in $R(1-P)$. We have shown that $H_0 = H_{D(H)+R(1-P)}^*$ can be decomposed as

$$H_0 = 0 \oplus G_0 \quad (75)$$

for some continuously invertible operator G_0 in $R(H) \oplus R(P)$ such

$$G_0^{-1} = \begin{pmatrix} A \\ PB \end{pmatrix} : R(H) \rightarrow \begin{pmatrix} R(H) \\ R(P) \end{pmatrix}$$

Let Q be any invertible self-adjoint operator in the Hilbert space $R(P)$.

By the given considerations G_0 has an invertible self-adjoint extension G such that

$$G^{-1} - \lambda = \begin{pmatrix} A - \lambda & B^*P \\ PB & Q - \lambda \end{pmatrix} : \begin{pmatrix} R(H) \\ D(Q) \end{pmatrix} \rightarrow \begin{pmatrix} R(H) \\ R(P) \end{pmatrix} \quad (76)$$

The following simple lemma will play role in the investigation of the absolutely continuous spectrum of the operator G .

Lemma (5-3-8) [133]:

Let A be a bounded self-adjoint operator and B a self-adjoint operator such that $k_2 k_1 I_\Delta(k_1)$ belongs to the trace class for every bounded interval Δ . Then

$$(k_1 + k_2)_{ac} = k_1' \oplus R$$

For some self-adjoint operator R and k_1' such that $k_1' \leq k_{1ac}$.

Corollary (5-3-9) [133]:

Let H be a symmetric operator with gap Δ in the complex Hilbert space \mathcal{H} suppose that $H = \bigoplus_{n=1}^{\infty} H_n$

For some symmetric operator $H_n, n \in \mathbb{N}$ with strictly positive deficiency numbers.

Then for every self-adjoint operator A in a separable Hilbert space there exists a self-adjoint extension \tilde{A} of A such that

$$\tilde{A} \in M'_{acJ}$$

Proof:

It easily follows from the spectral theorem that $M'_{acJ} = \bigoplus_{n=1}^{\infty} Q_{\rho_n}$ for suitably chosen $\rho_n \in L^1_{loc}, \rho_n \geq 0, \int \rho_n dx = 1$. On $\mathbb{R} \setminus J, n \in \mathbb{N}$. Since A can be decomposed into infinitely sets the operator A can be decomposed as

$$H = \bigoplus_{n=1}^{\infty} H^{(n)},$$

where for every $n \in \mathbb{N}$ operator $H^{(n)}$ is the orthogonal sum of two operators with infinite deficiency numbers. For every $n \in \mathbb{N}$ there exists a self-adjoint extension $\tilde{A}^{(n)}$ of $H^{(n)}$ such that $\tilde{A}^{(n)} \in Q_{\rho_n}$

Then $\tilde{A} = \bigoplus_{n=1}^{\infty} \tilde{A}^{(n)}$ is a self-adjoint extension of A and $\tilde{A} \in \bigoplus_{n=1}^{\infty} Q_{\rho_n} = M'_{acJ}$

Proposition (5-3-10) [133]:

Let A be a symmetric operator in some Hilbert space \mathcal{H} . Suppose that the operator A has some gap (α, β) and its deficiency number is infinite. Let \mathcal{I} be an open subset of (α, β) . Then A has a self-adjoint extension \tilde{A} such that

$$\sigma_{sc}(\tilde{A}) \cap \mathcal{I} = \overline{\mathcal{I}_0} \cap \mathcal{I}$$

Corollary (5-3-11) [133]:

Let A be a symmetric operator in the Hilbert space \mathcal{H} . Suppose A has a gap (α, β) and the deficiency number of A is infinite. Let \mathcal{I} be a non-empty open subset of (α, β) . Then there exists self-adjoint extensions \tilde{A}_1 and \tilde{A}_2 and \mathcal{K} a non-empty compact subset \mathcal{K} of \mathcal{I} with lebesgue measure zero such that \tilde{A}_1 and \tilde{A}_2 have a purely singular continuous spectrum in the gap (α, β) of A

$$\sigma_{sc}(\tilde{A}_1) \cap \mathcal{I} = \overline{\mathcal{I}_0} \cap \mathcal{I} \quad \text{and} \quad \sigma_{sc}(\tilde{A}_2) \cap \mathcal{I} = \mathcal{K}$$

We shall combine studying the point singular continuous and absolutely continuous spectra.

Theorem (5-3-12) [133]:

Let T be a symmetric operator in the Hilbert space H . Suppose that T has a gap (α, β) and its deficiency number is infinite. Then for every open subset I of (α, β) and every finite or infinite sequence $\{\lambda_n\}_{n=1}^N$ in I there exists a self-adjoint extension T_0 of T with the following properties.

- (i) $\delta_p(H_0) \upharpoonright_J = \{\lambda_n : n \in \mathbb{N}, n \leq N\}$ and for every eigenvalue λ_n of T_0 in I the multiplicity of λ_n equals the number of times it occurs in the sequence $\{\lambda_n\}_{n=1}^N$.
- (ii) $\delta_{sc}(H_0) \upharpoonright_J = \overline{J_0} \upharpoonright_J$
- (iii) $\delta_{ac}(H_0) \upharpoonright_J = 0$

Proof:

Let λ_0 be any point in I and if $N < \infty$, let $\lambda_n = \lambda_0$ for all $n \in \mathbb{N}, n \leq 2N$. Since for every regular point λ of T and in particular for $\lambda = \lambda_0$, the dimension of $N(H^* - \lambda)$ equals the deficiency number of T , we can choose by induction an orthogonal system $\{e_n\}_{n \in \mathbb{N}}$ such that $e_n \in N(H^* - \lambda_n), n \in \mathbb{N}$ and $e_j \in N(H^* - \lambda_0), j = 2n+1, n \in \mathbb{N}$ as in the proof of proposition (5-3-3) we can show that there exist self-adjoint operators T_0 and T_1 in the Hilbert space

$$H_0 = [e_{2n} : n \in \mathbb{N}, n \leq N], \quad H_1 = [e_n : n \in \mathbb{N}]$$

Respectively such that

$$e_{2n} \in N(M - \lambda_n), n \in \mathbb{N}, n \leq N, \quad e_{2n} \in N(M' - \lambda_n), n \in \mathbb{N}$$

$$e_j \in N(M' - \lambda_0), j = 2n+1, n \in \mathbb{N}$$

and T_0 and T_1 are restrictions of the adjoint T^* of T . Obviously T_0 is also a restriction of the operator T_1 and therefore, $MH_M^{\lambda_0}, H_M \subset H_M^{\lambda_0}$.

Where the operators $H_{\mathcal{H}}$ and $H_{\mathcal{H}}$ are defined by (66). Moreover the self-adjoint operator H has a pure point spectrum $\sigma_p(M) = \{\lambda_n : n \in \mathbb{N}, n \leq N\}$ and for every eigenvalue λ_n of H the multiplicity of λ_n equals the number of times it occurs in the sequence $\{\lambda_n\}_{n=1}^N$. By lemma (5-3-11) the operator $H_{\mathcal{H}}$ can be decomposed as $H_M = M \oplus G_0$. Where G_0 is a symmetric operator in the Hilbert space \mathcal{H} and $\delta(G_0)$ is also a gap of G_0 . We have only to show that the deficiency number of G_0 is infinite. In fact, then corollary (5-3-18) yields that there exists a self-adjoint operator G_0 in \mathcal{H} such that G_0 is an extension of G_0 and

$$(i) \quad \mathcal{D}_e(G) \upharpoonright J = \overline{J_0} \upharpoonright J, \quad \mathcal{D}_e(G) \upharpoonright J = \mathcal{D}_{pp}(G) \upharpoonright J = 0$$

Then obviously the operator $H_M = M \oplus G_0$ is a self-adjoint extension of H with the required spectral properties. By lemma (5-3-1) the operator $H_{\mathcal{H}}$ can be decomposed as $H_M = M' \oplus G'_0$. Where G'_0 is a symmetric operator in the Hilbert space \mathcal{H} and $\delta(G'_0)$ is also a gap of G'_0 . Since the symmetric operator G'_0 has a gap it has a self-adjoint extension G'_0 . Let $H_M = M' \oplus G'_0$. Since $M \subset M' \subset H$ we have $H_M = H$. Since H is a self-adjoint operator in \mathcal{H} and $\delta(M) \subset \overline{\delta(H)}$ it follows from lemma (5-3-1) that

$$H = M \oplus G'$$

for some self-adjoint operator G' in the Hilbert space \mathcal{H} . Since

$$H_{e_j} = M_{e_j} = G'_{e_j} = \lambda_0 e_j, \quad j = 2n-1, n \in \mathbb{N} \quad \text{the point } \lambda_0 \text{ in the gap } \delta(G') \text{ of } G' \text{ is an}$$

eigenvalue of G' with infinite multiplicity. Since obviously G' is a self-adjoint extension of G_0 this implies, by Krein's theorem, i.e., that the deficiency number of G_0 is infinite.