Chapter 1

Sum Rules for Jacobi Matrices and Their Applications to Spectral Theory

We show and prove a bound of a Jacobi matrix. And we give complete description for the point and absolutely continuous spectrum, while for the singular continuous spectrum additional assumptions are needed, we prove a characterization of a characteristic function of a row contraction operator and verify its defect operator. We also prove a commutability of an operator of this row contraction.

Section (1-1): Spectral Form for Jacobi Matrices:

The case of some rules and were efficiently used to relate properties of elements of a Jacobi matrix of certain class with its special properties. For instance spectral data of Jacobi matrices being a Hilbert space-Schmeidt perturbation of the free Jacobi matrix were characterization [42,101,135] and

we suggest a modification of the method that permits us to work with higher order sum rules. We obtain sufficient conditions for a Jacobi matrix to satisfy certain constraints on its spectral measure. We consider a Jacobi matrix [129,124].



Where $a = \{a_k\}, a_{k \ge 0}$ and $b = \{b_k\}, b_k \in \mathbb{R}$, We assume that $a = \{b_k\}, b_k \in \mathbb{R}$ is a compact perturbation of the free Jacobi matrix



A scalar spectral measure $((J-z)e_0,e_0)=\int_1^d \frac{\partial \delta(x)}{x-z}$ with $z=\varepsilon$ \ , the absolutely continuous spectrum $\delta_{\infty}(J)$ of fills in [-2, 2]and the discrete spectrum consist of two sequences $[x_j^{\pm}]$ with properties $x_j = 0$, $x_j = 0$ and $x_j^{\pm} > 0$, $x_j^{\pm} = 0$. Let $\partial_{\alpha} = \{a_k - a_{k-1}\}$ for a given $x_j = 0$ and $x_j^{\pm} > 0$, we construct a sequence $x_k(\alpha)$ by

formula $\chi(a)_j = \alpha_j^k - \alpha_j ... \alpha_{j+k-1}$ where $\alpha_j = \alpha_j - \alpha_j ... \alpha_{j+k-1}$ where $\alpha_j = \alpha_j - \alpha_j ... \alpha_{j+k-1}$

Theorem (1-1-1) [87]:

Let J = J(a,b) be a Jacobi matrix described above. If

(i)
$$a \rightarrow 1, b \in L^{m-1}, \partial_a, \partial_b \in L^2$$

(ii)
$$\gamma_k(a) \in L', k = 3, \lfloor (m+1)/2 \rfloor$$
 (2)

Then

$$(i')$$
 $\int_{-2}^{2} \log \delta'(x) \cdot \left(4 - x^2\right)^{m - \frac{1}{2}} dx > -\infty$

$$(ii') \sum_{i} \left(x^{\pm 2} - 4 \right)^{m+1/2} < \infty \tag{3}$$

When m the theorem gives the fact of theorem (1-1-1)

Proof:

Define $\phi_{m}(J)$ as $\varphi_{m}(J) = \varphi_{m}(S) = \varphi_{m,1}(S) + \varphi_{m,2}(S)$

$$= \frac{1}{2\pi} \int_{-2}^{2} \log \frac{1}{\delta(x)} (4-x^{2})^{m-\frac{1}{2}} dx + \sum_{i} G_{m}(x_{j}^{\pm}) .$$

We have to show that $(J) \leftarrow a_N = \{(a_N)_k\}$ and $a'_N = \{(a'_N)_k\}$, where



Define sequences b_N, b_N' in the same way (of course, with replaced by 0.

Let $J_N = J(a_N, b_N)$. As we readily see, $\dot{a}_N \to b_N \to b_N$. $\partial b_N \to b_N$, and $\chi(a_N') \to 0$ in corresponding norms, as $N \to b$ by the Lemma (1-1-4) below, we have for N' = N - m

$$\begin{aligned} & \left| \psi_{m} \left(J \right) - \psi_{m} \left(J_{N} \right) \right| \leq \psi_{m} \left(a_{N'}', b_{N'} \right) \leq C_{1} \left(\left\| a_{N'}' - 1 \right\|_{m+1} + \left\| b_{N'} \right\|_{m+1} \\ & + \left\| \delta a_{N'} \right\|_{2} + \left\| \delta b_{N'} \right\|_{2} + \sum_{k} \left\| \gamma_{k} \left(d_{N'} \right) \right\|_{1} \right) \end{aligned}$$

or
$$\psi_m(J_N) \longrightarrow \psi_m(J)$$
, as $N \longrightarrow$

on the other hand $(J_N - z)^{-1} \longrightarrow (J - z)^{-1}$ for $z \in \mathcal{L}$ \R, and consequently weakly $\mathcal{L}_{m,1}(\mathcal{S}) = \lim_{N \to \infty} (\mathcal{S})$ and $\lim_{M \to \infty} \mathcal{L}_{m,2}(\mathcal{S}) = \mathcal{L}_{m,2}(\mathcal{S})$ we bound the latter quantity $|\mathcal{L}_{m,2}(J)| = \sum_{j} \mathcal{L}_{m}(x_{j}^{\pm})$ $= \mathcal{L}_{m,2}(\mathcal{S}) = \mathcal{L}_{m,2}(\mathcal{S})$ with some constant $\mathcal{L}_{m,2}(\mathcal{S}) = \mathcal{L}_{m,2}(\mathcal{S})$ Summing up we obtain

$$\varphi(\delta) \le \lim_N \sup \varphi(\delta_N) = \lim_N \sup \psi(J_N) = \lim_N \psi(J_N) = \psi(J)$$

The proof is complete. It is easy to give simple conditions sufficient for $\chi_k(a) \in L' \quad \text{for the instance put} \quad (A_k(a))_j = \alpha_{j+1} + \ldots + \alpha_{j+k-1} - (k-1) \alpha_j \quad \text{,then}$ relations $a \to eL^{m+1}, \ \ eL^2 \quad \text{and} \quad A_k(a) \in L^{2(k,m)} \quad 2^{(k,m)=(m+1)!(m+2-k)} \quad \text{imply that}$ $\chi_k(a) \in L'$. In particular we have the following corollary.

Corollary (1-1-2) [87]:

Theorem (1-1-1) holds if conditions (i), (ii) are replaced with

 $A_k(a) \in L^{2(k,m)}$, $2(k,m) = (m-1)^k (m-1)^k (m-1)^k$, where $k = \delta, \left[\frac{m+1}{2}\right]$ we observe that relations (i) and (ii) are trivially true in the case of discrete Schrödinger operator i.e., when $A_k(a) \in L^{2(k,m)}$.

Corollary (1-1-3) [87]:

Then inequalities (i') and (ii') let hold J = (0,b). If $b \in L^{m-a}$, $A \in L^2$, the corollary is still true if $b \in L^{m+2}$, m being even. The proof is a sum rule of a special type. First we obtain it assuming rank $(J-J_0)<\infty$. Applying methods we see that

$$\frac{1}{2\pi} \int_{-2}^{2} \log \frac{1}{\mathcal{S}'(x)} \cdot \left(4 - x^{2}\right)^{m - \frac{1}{2}} dx + \sum_{j} G_{m}(x_{j}^{\pm}) = \psi_{m}(J)$$

Where $G_m(x) = (-1)^{m+1}C_0(x^2-4)^{m+\frac{1}{2}} + o(x^2-4)^{m+\frac{3}{2}}$ with $X \in \mathbb{R} \setminus \{-\frac{1}{2}, \frac{1}{2}\}$,

being a positive constant. where

$$\psi_{m}(J) = tr \left\{ \sum_{k=1}^{m} \frac{(-1)^{k+1}}{2^{2k+1}k} \left(J^{2k} - J_{0}^{2k} \right) - \frac{(2m-1)!!}{(2m)!!} \log A \right\}$$
(4)

Where $A=diag\{a_k\}$ and $\widetilde{C}_m^k=\frac{m\,!!}{(m-k)\,!!\,k!!}$. Notation is used for "even" or "odd" factorials.

Lemma (1-1-4) [5]:

Let J = J(a,b) we have

$$\left|\psi_{m}(J)\right| \leq C_{1} \left(\left\|a-1\right\|_{m+1} + \left\|b\right\|_{m+1} + \left\|\partial_{a}\right\|_{2} + \left\|\partial b\right\|_{2} + \sum_{k=3}^{\left[(m+1)/2\right]} \left\|\gamma_{k}(a)\right\|_{1} \stackrel{:}{:} \right)$$

$$(5)$$

Where G depends on G only. Above, norms G refer to the standard G space norms. We begin with considering expressions tr $G(J^{2k} - J_0^{2k})$ a rising in (4). Defining $V = J - J_0 = J(a - 1, b)$ we have

$$tr \Big(J^{2k} - J_{\scriptscriptstyle 0}^{2k} \Big) \hspace{-0.2cm} = \hspace{-0.2cm} tr \sum_{p=1}^{2k} \sum_{i_{\scriptscriptstyle 1}+\ldots+I_{\scriptscriptstyle p}=2k-p} \hspace{-0.2cm} VJ_{\scriptscriptstyle 0}^{i_{\scriptstyle 1}} 1 \ldots VJ_{\scriptscriptstyle 0}^{i_{\scriptstyle p}} p$$

we prove the lemma in steps.

Proof:

First we bounded summands corresponding to $P = \frac{m+1}{2}, m$ in [87]. We get $|tr(V^p F_p(J_0))| = |V^p F_p(J_0)|_{s_1} = |F_p(J_0)|_{s_1} |V^p|_{s_1}$ and for these

$$\left\|V^{P}\right\|_{s_{1}} = c_{10} \left\|V^{m-a}\right\|_{s_{1}} = c_{10} \left(\left\|a-a\right\|_{m-a}^{m-a} + \left\|b\right\|_{m-a}^{m-a}\right)$$

(6)

With the constant depending on $\|\cdot\|$. Similarly $\|tr \circ \mathcal{C}\| \le C_{11} \|a - 1\|_{m+1}^{m+1}$, let $\|a - 1\|_{m+1}^{m+1}$ onw. As we already mentioned in [134]

$$V^{p} = \sum_{j=0}^{p} \left(S^{i} p_{p,j}(a,b) + p_{p,j}(a,b) \overline{S}^{j} \right)$$
.

It is easy to show by induction that the polynomials $P_{p,p}(a,b)$ are particularly simple. Namely $P_{p,p}(a,b)$ wields that

$$\begin{split} trV^{p}F_{p}(J_{0}) = & (-1)^{p} \frac{(2m-1) \, !!}{2 \, p(2m) \, !!} trV^{p}J_{0,p} \\ = & (-1)^{p} \frac{(2m-1) \, !!}{2 \, p(2m) \, !!} tr(P_{p,p}(a,p) + P_{p,p}(a,b)_{p}) \\ = & (-1)^{p} \frac{(2m-1) \, !!}{2 \, p(2m) \, !!} \sum_{j} \alpha_{j} \alpha_{j-1} \dots \alpha_{j+(p-1)} \end{split}$$

Since $trV^{p}J_{0,s} = 0$ for $s \Rightarrow -1$. Hence $tr\left(V^{p}F_{p}(J_{0}) + (-1)^{p+1}\frac{(2m+1)!!}{p(2m)!!}\alpha^{p}\right)$

 $= (-1)^{p+1} \frac{(2m-1)!!}{2p(2m)!!} \sum_{j} \alpha_{j}^{p} - \alpha_{j} \alpha_{j+1} ... \alpha_{j+(p-1)} \quad \text{and we obtain that}$

$$\left| trV^{p} F_{p}(J_{0}) + (-1)^{p+1} \frac{(2m-1)!!}{2p(2m)!!} \alpha^{p} \right| \leq C_{12} \| \gamma_{p}(a) \|_{1}$$
(7)

Where c_n depends on p,m and sequences $\gamma_k(a)$ are defined in [134] Observe that $\gamma_p(a)=0$ when p=1. Furthermore we have for p=1 that

$$\begin{split} \sum_{j} (\boldsymbol{\alpha}_{j}^{2} - \boldsymbol{\alpha}_{j} \boldsymbol{\alpha}_{j+1}) &= \frac{1}{2} \sum_{j} (\boldsymbol{\alpha}_{j}^{2} - 2\boldsymbol{\alpha}_{j} \boldsymbol{\alpha}_{j+1} + \boldsymbol{\alpha}_{j+1}^{2}) \\ &= \frac{1}{2} \sum_{j} (\boldsymbol{\alpha}_{j} - \boldsymbol{\alpha}_{j+1})^{2} = \frac{1}{2} \| \boldsymbol{\alpha}_{l} \|_{2}^{2} \end{split}$$

So the left hand-side of (7) for p=2 can be estimated by $C_{13}\|\mathbf{a}\|_{2}^{2}$. It is also clear that inclusion \mathbf{a} and $\mathbf{a} \in L^{2}$ give that $\gamma_{p}(a) \in L'$ for $\mathbf{a} \in L^{2}$. Indeed

we have
$$\alpha^p - \alpha \alpha_{(1)} ... \alpha_{p-1} = \sum_{k=1}^p \alpha^{p-k} (\alpha - \alpha_{p-k}) \alpha_{(p-(k-1))} ... \alpha_{p-1}$$

The terms in the latter sum look like $\alpha(i_1)...\alpha_{(2p-1)}(\alpha-\alpha_{(i_p)})$ for some $i=(i_1,...,i_p)$. Obviously $\alpha-\alpha_{i_p}=a-a_{i_p}=a-a_{i_p}$. Applying the Holder

inequality
$$\sum_{k} a_{k} ... a_{p+k} \leq \sum_{k} \left(\sum_{j=1}^{p} \frac{1}{q_{j}} a_{j,k}^{q_{j}} \right) \quad \text{with} \quad a_{j,k} = (c_{\mathcal{F}})_{k} |.q_{j}| = 2(p-1) \quad \text{for}$$

$$a_{p,k} \rightarrow and$$
 $a_{p,k} = |(\alpha - \alpha_{p_p})_k|, q_p = \frac{1}{2}$ we get that

$$\left\| \mathscr{E} - \mathscr{C}_{(p-1)} \right\|_1 = C_{14} \left\| \mathscr{A} \right\|_2^2 + \left\| \mathscr{C}_{(p-1)}^{2(p-1)} \right\|_2^2$$

Which is finite for property. Thus gathering the above argument which is complete(see [134])we complete the proof of the lemma

Lemma (1-1-5) [87]:

Let
$$i = (i_1, ..., i_p)$$
 and $\sum_{s} i_s = n$ then

$$+\sum_{i}^{m,p_{i}}A_{k}[V,J_{0}]B_{k}[V,J_{0}]C_{k}$$

Where $P = (P_1, P_2, P_3), 1 = (L_1, L_2, L_3)$ and A_k, B_k, C_k are some bounded operators

Lemma (1-1-6) [87]

Let
$$\sum_{s} i = 2k - p$$
 we have $|tr(VJ_0^{i_1}...VJ_0^{i_p} - V^pJ_0^{k-p})| \leq c_3(|\partial u|_2 + |\partial v|_2)$

With $^{\varsigma}$ depending on $^{|\varsigma|}$ only. The lemma exactly bounded ,we may assume that operators $^{-}$ and $^{-}$ to commute we estimating $^{-}$ $\psi_{\alpha}(J)$

$$\psi_{m}'(J) = tr\left\{ \sum_{p=1}^{2m} V^{p} F_{p}(J_{0}) - \frac{(2m-1)!!}{(2m)!!} \log(I + \mathcal{X}_{p}) \right\}$$
(8)

Where $\mathscr{E}=diag\{a_k\}=A-I$ and

$$F^{p}(J_{0}) = \sum_{k=\lfloor (p+1)/2\rfloor}^{m} \frac{(-1)^{k+1}}{2^{2k+1}} \tilde{C}_{2m-1}^{2k-1} C_{2k}^{p} J_{0}^{2k-p}$$

Here ^q is a usual binomial coefficient, observe that for ^p we have

$$|tr(V^pF_p(J_0))| = F_p(J_0) ||V^p||_{\mathcal{S}} = C_4 (|a-1|_{m-1}^{m-1} - |b|_{m-1}^{m-1})$$

Where \square is the norm in the class of nuclear operators, hence it remains to bound the first m terms in (8) we have

$$\log(1+\widetilde{\alpha}) = \sum_{p=1}^{2m} \frac{(-1)^{p+1}}{p} \widetilde{\alpha} + o(\widetilde{\alpha}^{2m+1})$$

Set obe a symmetric matrix with 1's on p-th auxiliary diagonals and o's elsewhere the following lemma holds.

Lemma (1-1-7) [87]:

We have
$$F_p(J_0) = (-1)^{p+1} \frac{(2m-1)!!}{2p(2m)!!} J_{0,p}$$

Combining this with explicit form of $^{\prime\prime}$ and the series expansion for we get the required bound (7).

Section (1-2): Spectral Properties of Self-adjoint Extensions

Let A be a closed symmetric operator on a separable Hilbert space \cdot . If A has equal deficiency indices $\cdot^{n_k(A)} = \dim (h \in ran(A \pm iI))$, then A has a lot of self-adjoint extensions. These self-adjoint extensions can be labeled by the so-called Weyl function $\cdot^{M(i)}$ [82, 83, 84]. The generalization is based on concept of a boundary triple $\cdot^{\Pi=\{H, \Gamma_0, \Gamma_i\}}$ for \cdot^{K} being an abstract generalization of the Green's identity. Here \cdot^{K} is a separable Hilbert space with dim \cdot^{K} and \cdot^{K} and \cdot^{K} are linear mapping from dom \cdot^{K} to \cdot^{K} so that Green's identity is satisfied [108,119].

The problem is the following. Let M(x) be the Weyl function of a certain self-adjoint extensions A of A introducing the associated scalar Weyl function $M_h(x) = (M(x)h, h), h = 1$ is it possible to localize the different spectral subsets of A knowing the boundary values $M_h(x+i0), x = 1$ of the associated scalar Weyl function. Let A be separable Hilbert space. Recall that an operator function A with values in A is said to be a Hirglotz or Nevanlina function or R-function if holomorphic in A and for every A

the operator F(z) in F(z) is dissipative i.e., $Sm(F(z)) = \frac{\left(F(z) - F(z)^*\right)}{2i} \ge 0$. In the following we prefer the notion R-function. The class of R-functions with values in F(z) is denoted by F(z) is denoted by F(z) then there exist bounded self-adjoint operator F(z) in F(z) in F(z) then there exist bounded self-adjoint operator F(z) in F(z) in F(z) with F(z) such that

$$F(z) = C_0 + C_1 z + R^{\frac{1}{2}} (I_k + zL) (L - z)^{-1} R^{\frac{1}{2}} | H, \quad z \in \mathcal{E}_+$$
 (9)

Denoting by $E_L()$ the spectral measure of the self-adjoint operator one immediately obtains from (9) the representation

$$F(z) = C_0 + C_1 z + \int_{-\infty}^{\infty} \left(\frac{1}{t - z} - \frac{t}{1 + t^2} \right) \int_{F} (t), z \in \mathbf{f}_{+}$$
 (10)

Where $\sum_{F}^{\sum(i)}$ is an operator valued Borel measure on i given by

$$d\sum_{F}(t) = (1-t^{2})R^{\frac{1}{2}}dE_{L}(.)R^{\frac{1}{2}}, t \in i$$
(11)

the measure $\sum_{F}^{(j)}$ is self-adjoint and obeys

$$\int_{-\infty}^{+\infty} \frac{1}{1+t^2} d\sum_{F} (t) \in [H]$$

(12)

In contrast to spectral measures of self-adjoint operators it is not necessary true that ran Σ^g is orthogonal to ran (Σ^g) for adjoint Borel sets Γ and Γ . However the measure $\Gamma^{(j)}$ is uniquely determined by the R-function $\Gamma^{(j)}$. The integral in (10) is understood in the strong sense in the following $\Gamma^{(j)}$ is called the spectral measure of $\Gamma^{(j)}$ defined by

$$\int_{F} |t| \cdot \left| \int_{F} |0,t| : t \rangle 0 \right| \\
\cdot \int_{F} |t,0| : t \langle 0 \rangle$$
(13)

The distribution function $\sum_{F}^{(.)}$ is strongly left continuous and satisfies the condition

$$\sum_{F}(t) = \sum_{F}(t)^{*}, \sum_{F}(s) \leq \sum_{F}(t), -\infty \leq t < \infty$$

The distribution function $\sum_{E}^{(i)}$ is called the spectral function of $E^{(i)}$.

We note that the spectral function $\sum_{F}^{(j)}$ can be obtained by the Stieltjes transformation:

$$\frac{1}{2} \sum_{F} (t+0) + \sum_{F} (t) - \frac{1}{2} \sum_{F} (s+0) + \sum_{F} (s) = w - \lim_{y \to 0} \frac{1}{\pi} \int_{s}^{t} Sm(F(x+iy)) dx, \quad ts \in i$$
 (14)

Where it is used that the spectral function is strongly left continuous.

A will always denote a closed symmetric operator with equal deficiency indices $n_+(A) = n_-(A)$ [97,140,147,148].

We can assume that A is simple. This means that A has no self-adjoint parts. **Definition (1-2-1) [96]:**

A triple $\Pi=\{H,\Gamma_0,\Gamma_1\}$ consisting of an auxiliary Hilbert space and linear mapping $\Gamma_i:dom(A^*) \longrightarrow H$, i=0,1 is called a boundary triple for the adjoint

operator , is called a boundary triple for the adjoint operator of if

(i) The second Green's formula takes place

$$(A^*, f) - (f, A^*g) = (\mathbf{\Gamma}_1 f, \mathbf{\Gamma}_0 g) - (\mathbf{\Gamma}_0 f, \mathbf{\Gamma}_1 g), \quad f, g \in dom(A^*)$$

$$(15)$$

(ii) The mapping $\mathbf{T} = \{\mathbf{T}_0, \mathbf{T}_1\} : dom(A^*) \longrightarrow \mathbf{T}$ is surjective

Definition (1-2-2) [96]:

- (i) A closed linear relation in is closed subspace of H .
- (ii) The closed linear relation is symmetric if $(g_1, f_2) (f_1, g_2) = 0$ for all $\{f_1, g_1\}, \{f_2, g_2\} \in \Theta$
- (iii) The closed linear relation is self-adjoint if it is maximal symmetric. **Definition (1-2-3) [96]:**

Let $\{H, \mathbf{I}_0, \mathbf{I}_1\}$ be a boundary triple for

(i) for every self-adjoint relation in we put

$$D^{\mathscr{O}}\{f \in dom(A^*): \mathbf{I}_0^{-}f, \mathbf{I}_1^{-}f \in \mathcal{A}\}A^{\mathscr{O}} = A^*|D^{\mathscr{O}}$$

$$\tag{16}$$

- (ii) In particular we set $A_i = A^q$, i = 0,1, if Q, i = 0,1
- (iii) If $e^{G(B)}$ where B is an operator on ,then we set A^BA^O

Proposition (1-2-4) [96]:

Let {H, T, T} be a boundary triple for then for every self-adjoint relation in the operator given by definition (1-2-3) is self-adjoint extension of the mapping from the set of self-adjoint extensions in onto the set factor of self-adjoint extensions of is bijective. It is well known that Weyl function are an important tool in the direct and inverse spectral theory of singular Sturm-Liouville operators.

Definition (1-2-5) [96]:

Let $\{H, \mathbf{I}_{\overline{0}}, \mathbf{I}_{\overline{1}}\}$ be a boundary triple for the operator \mathcal{I} . The Weyl function of A corresponding to the boundary triple $\{H, \mathbf{I}_{\overline{0}}, \mathbf{I}_{\overline{1}}\}$ is the unique mapping

$$M(.): A(A_0) \longrightarrow Satisfying$$

$$\mathbf{I}_1 f_z = M(z) \mathbf{I}_0 f_z, \quad f_z \in N_z, \quad z \in A_0$$
(17)

Where $N_z = \ker(A^* - zI)$ above implicit definition of the Weyl function is correct and the Weyl function is a R-function obeying

$$o \in \rho(Sm(M(i)))$$

Definition (1-2-6)[96]:

Proposition (1-2-7) [96]:

Let A be a simple closed densely defined symmetric operator in h. Suppose that $\{H, T_0, T_1\}$ is a boundary triple for M(A) is the corresponding Weyl function , a self-adjoint relation in and $A \in \rho(A_0)$. Then the following holds.

- (i) $\lambda \in A^{(a)}$ if and only if $0 \in A^{(a)}$.
- (ii) $\lambda \in \delta_r(A^{\theta})$ if and only if $0 \in \delta_r(\Theta M(\lambda)), \tau = p, c$

If A is a simple symmetric operator then the Weyl function M(x) determines the pair $\{A_1,A_0\}$ up to unitary equivalence. We shall often say that M(x) is the Weyl function of the pair $\{A_1,A_0\}$ or simply of A. We can prove $M_1(x)$ and $M_2(x)$ with values in A and A are connected via

$$M_2(z) = K^* M_1(z) K + D \tag{18}$$

Where $D = D^* \in [H_2]$ and $K \in [H_2, H_1]$ is boundedly invertible. With each boundary triple we can associate a so-called -field (.) corresponding to is defined by

$$\gamma(z) = \left(\Gamma_{\scriptscriptstyle 0} | N_{\scriptscriptstyle z}\right)^{-1} : \mathsf{H} \longrightarrow N_{\scriptscriptstyle 2}, z \in \rho(A_{\scriptscriptstyle 0}) \tag{19}$$

One can easily check that

$$\gamma(z) = (A_0 - z_0)(A_0 - z)^{-1} \gamma(z_0), z.z_0 \in \rho(A_0)$$
(20)

And consequently is a field. The field and the Weyl function are related by

$$M(z) - M(z_0)^* = (z - \overline{z}_0) \gamma(z_0)^* \gamma(z), \quad z, z_0 \in \rho(A_0)$$
 (21)

The relation (21) means the M(x) is a function of a pair $\{A,A_0\}$. Further we note that if A is simple then $N_z,z\in\mathcal{A}(A_0)$ is generating with respect to A too.

Let be a Borel measure on . A support of is a set such that $\mu(i \setminus S) = 0$ we note that the closed supports $\mu(i \setminus S) = 0$ and $\mu(i \setminus S) = 0$ we denote that the closed supports $\mu(i \setminus S) = 0$ are not generally mutually disjoint to obtain mutually disjoint supports we introduce the following sets .

$$S_0'(\mu) = \{t \in \mathbf{i} : d\mu(t) \text{ dt exists and } d\mu(t) = \emptyset \}$$
 (22)

$$S'_{pp}(\mu) = \{t \in i : \mu(\{t\}) \neq 0\}$$
 (23)

$$S'_{sc}(\mu) = \left\{ t \in \mathbf{i} : d\mu(t) d(t) \text{ exists } \frac{d\mu(t)}{dt} = \infty \text{ and } \mu(t) = 0 \right\}$$
 (24)

$$S'_{ac}(\mu) = \left\{ t \in \mathbf{i} : \frac{d\mu(t)}{dt} \text{ exists and } 0 < d\mu(t) / dt < \infty \right\}$$
 (25)

Where the distribution function is similar to (13) defined by it turns out that. Since the sets $S_{T}'(\mu)$, are of Lebesgue measure zero and mutually disjoint we find that for any Borel set one has

$$\mu(\chi \cap S_T'(\mu)) = \mu_T(\chi), T = s, pp, sc, ac$$
(26)

The sets $s'_{\varepsilon}(\mu)$, $s'_{\varepsilon}(\mu)$, $s'_{\varepsilon}(\mu)$, and $s'_{\varepsilon}(\mu)$ singular pure point ,singular continuous and absolutely continuous supports of respectively. We note that

$$S_{pp}(\mu) = \overline{S'_{pp}(\mu)} \quad \text{and} \quad S_{\tau}(\mu) \subseteq \overline{S'_{\tau}(\mu)} \subseteq S(\mu) \quad , \quad \text{ (27)}$$

In general it is not possible to replace inclusion by equalities, let now \ge be a measure with values in \bowtie the measure \ge admit a Lebesque- Jordan decomposition $\ge \ge \ge$ \ge .As above the notation

$$S_s > S > S_{pp} > S_{sc} > S_$$

stand for the singular pure point, singular continuous and absolutely. We get

$$S_{p}(\sum) = \{\tau \in i \} : \sum \{\tau\} \neq 0$$
 (28)

we have $S_p(\sum) = S_{pp}(\sum)$ and $\overline{S_p(\sum)} = S_{pp}(\sum)$ with each operator-valued measure $\sum (.) = (\sum J)h, h), h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h), h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h), h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h), h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h$, $h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h$, $h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h$, $h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h$, $h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h$, $h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h$, $h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h$, $h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h$, $h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h$, $h \in H$. In the following we are interested in the problem whether the spectral properties of the operator valued measure $\sum (.) = (\sum J)h, h$, $h \in H$. In the following we are interested in the operator valued measure \sum

$$S_s'(\sum \tau) = \cup_{k=1}^N S_s'(\sum_{k_k}) \tag{29}$$

$$S'_{pp}\left(\sum_{k=1}^{N}S'_{pp}\left(\sum_{k_{k}}\right)\right) \tag{30}$$

$$S'_{sc}\left(\sum_{t}\tau\right) = \bigcup_{k=1}^{N} S'_{sc}\left(\sum_{t_{k}}\right) |S'_{pp}\left(\sum\right)$$
(31)

$$S'_{ac}\left(\sum_{k}\tau\right) = \bigcup_{k=1}^{N} S'_{ac}\left(\sum_{k}\right) |S'_{s}\left(\sum\right)$$
(32)

Lemma (1-2-8) [96]:

Let be a separable Hilbert space and $T = \{h_k\}_{k=1}^N$, $1 \le N \le \infty$ be a total set in . Then the sets $S_s'(\sum \tau)$, $S_{pp}'(\sum \tau)$, $S_{sc}'(\sum \tau)$ and $S_{cc}'(\sum \tau)$ are singular ,pure point ,singular continuous and absolutely continuous supports of Σ respectively i.e.,

$$\sum (\chi \cap S_{\tau}'(\sum \tau)) = \sum^{\tau} (\chi), \tau = s, pp, sc, ac$$
(33)

For any Borel set Ref. . In particular the following relations hold.

$$S_p' \sum = S_{pp}' (\sum \tau)$$
 and

$$S_{p}(\sum) \subseteq \overline{S'_{p}(\sum \tau)} \subseteq S(\sum), \tau = s, sc, ac$$
 (34)

Proof:

By the Lebesgue-Jordan decomposition one easily gets that for each

We have

$$\left(\sum(\chi)h,h\right) = \sum_{\lambda,\tau}(\chi), \ \tau = s, pp, sc, ac$$
 (35)

For any Borel set $\succ^{=R}$ where $\sum_{s,\tau}(.)$ arises from the Lebesgue-Jordan decomposition of the scalar measure $\sum_{s}(.)$. Let - . Since mes $S_s'(\sum_i \tau) = 0$ We get

$$\left(\sum (\mathcal{X} \cap S_s'(\sum \tau)) h_k, h_k\right) = \sum_{k} \left(\mathcal{X} \cap S_s'(\sum \tau)\right) = \sum_{k,s} \left(\mathcal{X} \cap S_s'(\sum \tau)\right)$$
(36)

For any $h_k \in \tau$ using (35),(36) and

$$\sum_{\mathbf{k}_{s,s}} \left(\mathbf{x} \cap \mathbf{S}_{s}' \left(\sum_{\mathbf{\tau}} \mathbf{\tau} \right) \right) = \sum_{\mathbf{k}_{s,s}} \left(\mathbf{x} \cap \mathbf{S}_{s}' \left(\sum_{\mathbf{\tau}} \mathbf{\tau} \right) \cap \mathbf{S}_{s}' \sum_{\mathbf{k}_{s}} \right)$$

$$= \sum_{\mathbf{k}_{s,s}} \left(\mathbf{x} \cap \mathbf{S}_{s}' \left(\sum_{\mathbf{\tau}} \mathbf{\tau} \right) \right) = \sum_{\mathbf{k}_{s,s}} \left(\mathbf{x} \right)$$
(37)

We find $\left(\sum \left(\chi \cap S_s'\left(\sum \tau\right)\right) h_k, h_k\right) = \left(\sum_{k=1}^{s} (\chi) h_k, h_k\right)$ for any $h_k \in \tau$. Since is

total we finally obtain $\sum (x \cap S'_s(\sum \tau)) = \sum_{s=0}^s (x)$ for any Borel set $x \in I$. Similarly we prove the statements for

Let $x \in S'_{pp} \sum \tau$. Then there is $h_k \in \tau$ such that $x \in S'_{pp} (\sum_{h_k})$. Hence $\sum_k (\{x\}) = (\sum (\{x\})h_k, h_k) \neq 0$ which yields $\sum (\{x\}) \neq 0$ or $x \in S_p(\sum)$, i.e. $S'_{pp} (\sum_i \tau) \subseteq S_p(\sum)$ conversely if $x \in S_p(\sum)$ then there is a $h \in \Gamma$ such that $\sum (\{x\}) \neq 0$. If this is not the case then for each $h_k \in \Gamma$ one has $\sum (\{x\}) = (\sum (\{x\})h_k, h_k) = 0$. Since Γ is total this yields $(\sum (\{x\})h_k, h_k) = \sum (\{x\}) = 0$

For each $s \in S'$ Contrary to the assumption. however, if there is a $s \in S'$ such that $\sum_{k} (\{x\}) \neq 0$ then $x \in S'_{pp}(\sum_{k} \tau)$, $i, e.S_{p}(\sum_{k} \tau) \subseteq S'_{pp}(\sum_{k} \tau)$ hence $S_{p}(\sum_{k} \tau) = S'_{pp}(\sum_{k} \tau)$. Further from (33) we get

$$S_{T}(\sum)\subseteq \overline{S'_{T}(\sum \tau)}$$
, $\tau=s$, pp , sc , ac

Taking (27) into account we get $S_T'(\sum_k) \subseteq S(\sum_k)$, $\tau = s, sc, ac, sc$ for each $s_k \in \tau$. Since $S(\sum_{n_k}) \subseteq S(\sum_n)$ for each $s_k \in \tau$ we get $S_T'(\sum_n \tau)$ which immediately proves (34). Taking (20) and (21) into account we obtain that $s_n = t$ which leads to the representation.

$$M(z) = C_0 + \int_1 \left(\frac{1}{t-z} - \frac{t}{1+t^2}\right) d\sum(t), z \in £$$
 (38)

Lemma (1-2-9) [96]:

Let A be a simple densely defined closed symmetric operator on the a separable Hilbert space with $n_+(A) = n_-(A)$. Further, let $\Pi = [H, \Gamma_0, \Gamma_1]$ be a

boundary triple of $^{\ell}$ with Weyl function Σ . If $^{E_{A_{0}}(.)}$ is the spectral measure of $^{A_{0}} \longrightarrow ^{*}|\ker(\mathbb{T}_{0})(=E\times t_{A})|$ and Σ that of the integral representation (38) of the Weyl function Σ . Then the measure $^{E_{A_{0}}(.)}$ and Σ are equivalent. In particular one has $\mathcal{E}_{\ell}(A_{0}) = \mathcal{E}_{\ell}(\Sigma)$.

Theorem (1-2-10) [96]:

Let A be a simple densely defined closed symmetric operator on a separable Hilbert space h with $n_{\rightarrow}(A) = n_{\rightarrow}(A)$. Further let $\Pi = [H, \Gamma_0, \Gamma_1]$ be a boundary triple of K with Weyl function M(X).

If $E_{A_0}(.)$ is the spectral measure of $A_0 = A^* \ker \mathbf{T}_0 (=Ext_A)$ and Σ that of the integral representation (38) of the Weyl function M(.), then for each total set

$$\tau = \{h_k\}_{k=1}^N$$
, $1 \le N \le +\infty$ in the sets $S_s'(\sum_i \tau)$, $S_{pp}'(\sum_i \tau)$, $S_{sc}'(\sum_i \tau)$ and $S_{ac}'(\sum_i \tau)$.

Singular pure point, singular continuous, and absolutely continuous supports of $E_{A_n}(.)$ respectively, i.e. we have

$$E_{A_{0}}\left(\varkappa \cap S_{T}'\left(\sum \tau\right)\right) = E_{A_{0}}^{T}\left(\varkappa\right) \tag{39}$$

$$\delta_{\tau}(A_0) \subseteq S_{\tau} \sum \tau \subseteq \delta(A_0), \tau = s, sc, ac$$
 hold.

Proof:

Since by lemma (1-2-8) the sets $S'_{\tau}(\mathcal{S},\tau), \tau=s, pp, sc, ac$ are supports of Σ , one immediately gets from lemma (1-2-9) that the same sets are supports of $E_{A_n}(\cdot)$ of the same type, i.e., (39) holds . If $x \in S'_{pp}(\Sigma \tau)$ then there is at least one

$$(E_{A_0}(\lbrace x\rbrace) \gamma(i) h_k, \gamma(i) h_k) \neq 0$$

Hence $x \in \mathcal{F}_p(A_0)$ conversely, if $x \in \mathcal{F}_p(A_0)$ then due to the fact that $\gamma(i)\tau$ is

Generating for $E_{A_n}(.)$ then is at least one such that

$$(E_{A_0}(\lbrace x\rbrace) \gamma(i) h_k, \gamma(i) h_k) \neq 0$$

Hence $x \in S'_{pp}\Sigma_{r}^{\tau}$ which proves $\mathcal{S}_{p}(A_{0}) = S'_{pp}(\Sigma_{\tau})$ the relations $\mathcal{S}_{s}(A_{0}) \subseteq \overline{S'_{\tau}(\Sigma_{r}\tau)} \subseteq \mathcal{S}(A_{0})$, $\tau = s, sc, ac$ are consequences of lemma (1-2-8) and lemma (1-2-9) .we characterize the spectral properties of the operator-valued measure $\Sigma_{r}^{(s)}$ using the boundary behavior of the Weyl-function $\mathcal{S}(A_{0}) \subseteq \overline{S'_{\tau}(\Sigma_{r}\tau)} \subseteq \mathcal{S}(A_{0}) \subseteq \mathcal{S}($

$$\int_{R} \frac{d\mu(t)}{1+t^2} < +\infty$$

(40)

Let us associate with the Poisson integral

$$v(z) = \int_{R} \frac{yd\mu(t)}{(t-x)^2 + y^2}, z = x + iy \in \mathbf{f}_{+}$$

(41)

Which defines a positive harmonic function in $V_i(z)$. Conversely it is well known that each positive harmonic function $V_i(z)$ in admits the representation $V_i(z) = ay + V(z)$ with v(z) = ay + V(z) with v(z) = ay + V(z) of the form (40) and (41). Below we summarize some well-known facts on positive harmonic function **Proposition (1-2-11) [96]:**

Let be a positive Radan measure obeying (40) and let v(z) be a positive harmonic function in $z = x + iy \in \mathcal{E}_+$ defined by (41). Then one has.

(i) for any = the $\lim_{V(x \to io)} = \lim_{V(x \to iy)} = \text{exists and is finite}$, if and only if symmetric derivative $= D_{\mu}(x)$

$$D_{\mu}(x) = \lim_{\varepsilon \to 0} \frac{\mu(x + \varepsilon) - \mu(x - \varepsilon)}{2\varepsilon}$$
(42)

Exists and is finite. In this case one has

$$V(x+io) = \pi D_{\mu}(x) \tag{43}$$

- (ii) if the symmetric derivative $D_{\mu}(x)$ exists and is infinite the
- (iii) for each $x \in \mathbb{R}$ one has $Sm(z \to x)V(z) \to \mathbb{R}\{x\}$ as $z \to \infty$
- (iv) $v^{(z)}$ converges to a finite constant as $z \rightarrow x$, if and only if the derivative $d \mu(t) dt$ exists at t = x and is finite.

The symbol \rightarrow means that the limit $\lim_{r\to 0} V\left(x + re^{i\theta}\right), x \in \mathbb{R}$ exists uniformly in $\theta \in [\varepsilon, \pi - \varepsilon]$ for each $\varepsilon \in (0, \pi/2)$. Proposition (1-2-11) allows us to introduce measures satisfying (40) the following sets z = (x + iy)

$$S_s''(\mu) = \{x \in i : V(z) \longrightarrow as \quad z \longrightarrow x$$
 (44)

$$S_{pp}^{\prime\prime}(\mu) = x \in i : \operatorname{Sm}_{z \to x} (z - x) V(z) > 0 \tag{45}$$

$$S_{sc}''(\mu) = \{x \in \mathbf{i} : V(z) \longrightarrow and (z - x)V(z) \longrightarrow as z \longrightarrow x\}$$
(46)

$$S_{sc}^{\prime\prime}(\mu) = \left\{ x \in i : V(x + i0) \text{ exists and } 0 < V(x + i0) < \infty \right\}$$
 (47)

Obviously the sets $s_{s}^{s}(\mu)$ and $s_{sc}^{s}(\mu)$ as well as $s_{pp}^{s}(\mu)$. $s_{sc}^{s}(\mu)$, and $s_{sc}^{s}(\mu)$ are mutually disjoint. By proposition (1-2-15) one immediately gets that $s_{pp}^{s}(\mu) = s_{pp}^{s}(\mu)$ and

$$S_{\tau}'(\mu) \subseteq S_{\tau}'(\mu) \subseteq S(\mu)$$
 (48)

Indeed the relation $S_{pp}^{r}(\mu) = S_{pp}^{r}(\mu)$ is a consequence of (iii).By (ii) we get $S_{s}^{r}(\mu) = S_{s}^{r}(\mu)$

Similarly we prove $S_{sc}(\mu) = S_{sc}(\mu)$ using (ii) and(iii). Finally the relation $S_{sc}(\mu) = S_{sc}(\mu)$ follows from (i). We note that it can happen that $S_{sc}(\mu) \neq 0$ and the inclusion $S_{sc}(\mu) = S_{sc}(\mu)$ is strict even if $\mu_{sc} = 0$. Furthermore we note that from (26) and the inclusion $S_{sc}(\mu) = S_{sc}(\mu)$, $\tau = 0$, pp, sc, ac we find that

$$\mu(\chi \cap S_{\tau}'(\mu)) = \mu_{\tau}(x) \tag{49}$$

For any Borel set $x \in I$. Now we are going to characterize the spectral parts of the extension A by means of boundary values of the Weyl function A .

Using the integral representation (38) of the Weyl function we easily get that

$$V_{h}(z) = \int_{i} \frac{y}{(x-t)^{2} + y^{2}} d\sum_{h}(t) = Sm(M_{h}(z)), z \in E, \quad h \in H$$

$$(50)$$

Where
$$M_h(z) = (M(z)h,h), z \in \mathfrak{t}, h \in H$$
 (51)

The function $M_h(.)$ is a scalar R-function. Since $M_h(.)$ arises from the Weyl function we call it the associated scalar Weyl function $V_h(.)$ is imaginary part of the associated scalar Weyl function $M_h(.)$ and the theory developed we can relate the boundary behavior at the real axis the imaginary part of associated scalar Weyl functions with the spectral properties of the self-adjoint extension $A_h(.)$ To this end in addition to (29) and (32) we introduce.

$$S_s''(\sum \tau) = \bigcup_{k=1}^N S_s''(\sum_{k_k}) \tag{52}$$

$$S_{pp}^{"}\left(\sum_{k=1}^{n}S_{pp}^{"}\left(\sum_{k_{k}}\right)\right) \tag{53}$$

$$S_{sc}''\left(\sum_{k=1}^{N}\mathcal{I}\right) = \bigcup_{k=1}^{N}S_{sc}''\left(\sum_{k}\right) \setminus S_{pp}''\left(\sum\right)$$
(54)

$$S_{ac}(\sum \tau) = \bigcup_{k=1}^{N} S_{ac}(\sum_{k}) \setminus S_{s}(\sum)$$

$$(55)$$

By definition the sets $S_s''(\Sigma,\tau)$ are disjoint. They holds for $S_{pp}''(\Sigma,\tau)$.

Furthermore we denote that the sets $S_T''(\Sigma,\tau)$ have Lebesgue zero, i.e., mes $S_T''(\Sigma,\tau) = 0$, $\tau = s$, pp, sc. , it turns out that the sets $S_T''(\Sigma,\tau)$ in theorem (1-2-14) can be replaced by the sets $S_T''(\Sigma,\tau)$

Theorem (1-2-12) [96]:

Let A be a simple densely defined closed symmetric operator on a separable Hilbert space H with $n_+(A) = n_-(A)$. Further, let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple of K with Weyl function M(X). If $E_{A_0}(X)$ is the spectral measure of $A_0 = A^* |\ker \Gamma_0(E_{X} x_A)|$ and total set $T = \{h_k\}_{k=1}^N$, $1 \le N \le I_0$ in I the

sets $S_s''(\Sigma_i \tau)$, $S_{pp}''(\Sigma_i \tau)$, $S_{sc}''(\Sigma_i \tau)$ and $S_{sc}''(\Sigma_i \tau)$ are singular ,pure point , singular continuous and absolutely continuous supports of $E_{A_s}(\cdot)$ respectively , i.e., we have

$$E_{A_0}\left(\chi \cap S_{\tau}^{\tau}\left(\sum \tau\right)\right) = E_{A_0}^{\tau}\left(\chi\right), \tau = s, pp, sc, ac$$
(56)

For each Borel set $x \in I$. In particular it hold $\delta_p(A_0) = S_{pp}(\sum \tau)$ and

$$\mathcal{S}_{\tau}(A_0) \subseteq S_{\tau}(\sum \tau) \subseteq \mathcal{S}(A_0), \tau = s, sc, ac$$
.

Proposition (1-2-13) [96]:

Let (x) be a scalar R-function. Then for almost all (x) the limit (x) (x) (y) (y)

Theorem (1-2-14) [96]:

Let A be a simple densely defined closed symmetric operator on a separable Hilbert space h with $n_{+}(A) = n_{-}(A)$ Further let $\Pi = \{H\Gamma_0, \Gamma_1\}$ be a boundary triple of ℓ with Weyl function ℓ and let $E_{A_0}(0)$ be the spectral measure of the self- adjoint extension ℓ of ℓ . If $\tau = \{h_k\}_{k=1}^N$, $\iota = -\infty$ is a total set in ℓ then sets $\Omega_{\ell}(M;\tau)$, $\Omega_{pp}(M;\tau)$, $\Omega_{cc}(M;\tau)$ and $\Omega_{cc}(M;\tau)$ are supports of $E_{A_0}(\ell)$ respectively. i.e., we have

$$E_{A_0}\left(\mathcal{X} \cap \Omega_{\mathcal{A}}(M;\tau)\right) = E_{A_0}^{\tau}\left(\mathcal{X}\right), \ \tau = s, pp, sc.ac$$
 (57)

For each Borel set $\stackrel{\text{\tiny ZeR}}{=}$ In particular it holds $\stackrel{\text{\tiny ZeR}}{\circ}(A_0) = \Omega_p(M;\tau)$ and $\stackrel{\text{\tiny ZeR}}{\circ}(A_0) = \overline{\Omega_p(M;\tau)} = \mathcal{S}_p(A_0)$ for $\stackrel{\text{\tiny ZeR}}{=}$. We note that the inclusions $\stackrel{\text{\tiny ZeR}}{\circ}(A_0) = \overline{\Omega_p(M;\tau)}$ and $\stackrel{\text{\tiny ZeR}}{\circ}(A_0) = \overline{\Omega_p(M;\tau)}$ of theorem (1-2-14) may be strict even if $\stackrel{\text{\tiny ZeR}}{\circ}(A_0)$ is empty.

Let $^{\mu}$ be a Borel measure on $^{+}$ and let $^{\varkappa \subseteq i}$ be a Borel set the set

$$CL_{ac}(x) = x \subseteq i : mes((x - \varepsilon x + \varepsilon) \cap x) > 0 \forall \varepsilon > \delta$$
 (58)

(- -)

is called the absolutely continuous closure of set $^{-}$ obviously the set $^{CL_{ac}(x) \in \bar{x}}$ is always closed and one has

Lemma (1-2-15) [96]:

Let (A) be a scalar R-function which has the representation (10) then $S_{ac}(A) = CL_{ac}(\Omega_{ac}(A))$

Proof:

If
$$x \Leftarrow CL_{ac}(\Omega_{ac}(\varphi))$$
 then there is an \Rightarrow such that $mes(x - \varepsilon x + \varepsilon) \cap \Omega_{ac}(\varphi) = \theta$

$$\mathcal{H}_{ac}(x - \varepsilon x + \varepsilon) = \mathcal{H}_{ac}(x - \varepsilon x + \varepsilon) \cap \Omega_{ac}(\varphi) = 0 \tag{59}$$

Hence $x \in S(\mu_{le}) = S_{ae}(\mu)$ which yields $S_{ae}(\mu) = CL_{ae}(\Omega_{le}(\Phi))$ conversely if $x \in S_{ae}(\mu)$ then there is an e > 0 such that $(x - \varepsilon, x + \varepsilon) \cap \Omega_{le}(\mu) = 0$ then $\mu_{le}(x - \varepsilon, x + \varepsilon) = 0$ using

$$\mu_{ac}(x-\varepsilon,x+\varepsilon) = \mu_{ac}(x-\varepsilon,x+\varepsilon) \cap \Omega_{ac}(\phi) \int_{(x-\varepsilon,x+\varepsilon)\cap\Omega_{ac}(\phi)} \frac{d\mu(t)}{dt} dt = 0$$
 (60)

and proposition (1-2-11) (i) and (vi) one gets

$$\mu_{c}(x-\varepsilon,x+\varepsilon) = \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon \right\} - \frac{1}{\pi} \left\{ x - \varepsilon x + \varepsilon x +$$

Hence $Sm(\phi t \to t0)dt = 0$ for $a.e.t \in (x \to x + E) \cap \Omega_{ac}(\phi)$. However by definition of the set $\Omega_{ac}(\phi)$ one has $Sm(\phi(\tau + t0))dt > 0$ for all $\varpi\Omega_{ac}(\phi)$ which implies $mes((x \to x + E) \cap \Omega_{ac}(\phi)) = 0$

Hence $x \not\in CL_{ac}(\Omega_{ac}(\Phi))$ or equivalent $CL_{ac}(\Omega_{ac}(\Phi)) \subseteq S_{sc}(\mu)$.

Proposition (1-2-16) [96]:

Let A be a simple densely defined closed symmetric operator a separable Hilbert space with $n_+(A) = n_-(A)$. Further let $\Gamma = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple of A with Weyl function M(A) If $\tau = \{h_k\}_{k=1}^N$, $\tau = \{h_k\}_{k=1}^N$, is a

total set in then the absolutely continuous spectrum of the self-adjoint extension of is given by.

$$\delta_{ac}(A_0) = \overline{\bigcup_{k=1}^{N} CL_{ac}(\Omega_{ac}(M_{h_k}))}$$
(62)

Theorem (1-2-17)[96]:

Let be a simple densely defined closed symmetric operator on a separable Hilbert space with $n_+(A) = n_-(A)$. Further, let $\Pi = \{H, \Gamma_0, \Gamma_1\}$ be a boundary triple of K with Weyl function M(X).

If $\tau = \{h_k\}_{k=1}^N$, is a total set in , then for the self-adjoint extension of the following conclusions are valid:

(i) The self-adjoint extension A of A has no point spectrum within the interval $^{(a,b)}$. i.e., $^{B_{a}}(A_{0}) \cap ^{A}(a,b) = 0$ if and only if for each one has

$$\lim_{y\to 0} y M_{hk}(x+iy) = 0$$

(63)

for all x = (a,b). In this case the following relation holds

$$\delta(A_0) \cap (a,b) = \frac{\delta_c(A_0) \cap (a,b)}{\bigcup_{k=1}^N \Omega_{ac}(M_{bk}) \cup} = \overline{\bigcup_{k=1}^N \Omega_{ac}(M_{bk})} \cap (a,b)$$
(64)

- (ii) The self-adjoint extension A of A has no singular continuous spectrum within the interval (a,b), $i.e. \not\subset (A_0)$ $(a,b) = \emptyset$ if for each A the set A is countable in particular, if $(a,b)|\Omega_{ac}(M_{hk})$ is countable.
- (iii) The self-adjoint extension A of A has no absolutely continuous spectrum within the interval $^{(a,b)}$.i.e., $\mathscr{E}_{c}(A_{0}) \cap (a,b) = \mathscr{E}$ if and only if for each $^{A} \cap (a,b) = \mathscr{E}$ the condition

$$Sm(M_{hk}(x+i0))=0 \tag{65}$$

holds for a.e. x = (a,b) . in this case we have

$$\mathcal{S}_{s}(A_{0}) \cap (a,b) = \overline{\Omega_{s}(M;\tau)} \cap (a,b)$$

Proof:

- (i) If condition (65) is satisfied for all x = (a,b) and all k = 2....N, then a simple computation shows that $\lim_{z \to \infty} (z x) M_{hk} = 0$ holds for all x = (a,b) and each k = 2....N too. Therefore (a,b) = 0 for k = 2....N which yields (a,b) = 0 theorem (1-2-14). Implies (a,b) = 0 which yields (a,b) = 0.
- (ii) Conversely if $\mathcal{S}_{p}(A_{0}) \cap (a,b) = 0$ then $\mathcal{S}(A_{0}) \cap (a,b) = 0$ again by theorem (1-2-14) we find $\mathcal{S}_{p}(A_{0}) \cap (a,b) = 0$ therefore $\mathcal{S}_{p}(A_{0}) \cap (a,b) = 0$ for each $\mathcal{S}_{p}(A_{0}) \cap (a,b) = 0$ therefore $\mathcal{S}_{p}(A_{0}) \cap (a,b) = 0$ which yields $\lim_{y \to 0} yM_{hk}(x+iy) = 0$ for all x = a,b and each x = a,b. The first of relation (64) is consequence of $\mathcal{S}_{p}(A_{0}) \cap (a,b) = 0$. The second part of relation (64) is a consequence of theorem (1-2-18) which shows that

$$\mathcal{S}_{\epsilon}(A_{\scriptscriptstyle 0}) \subseteq \overline{\Omega_{\epsilon}(M; \tau)} = \overline{\mathbb{Q}_{\epsilon}(M; \tau)} \subseteq \mathcal{S}(A_{\scriptscriptstyle 0}), \tau = sc, ac$$
(67)

and $\mathcal{S}(A_0) = \mathcal{S}(A_0) \cup \mathcal{S}(A_0)$. Both facts imply that $\delta_c(A_0) \cap (a,b) \subseteq$

Which proves (64)

(ii) By (53) we gets that $S_{ac}(\sum) = S(\sum_{sc}) = S(\sum_{b}) = S(M_{hk})$. Therefore if $\Omega_{c}(M_{hk}) = (a,b)$ is countable, then so is $S_{ac}(\sum_{b}) = (a,b)$ this yields that the singular continuous measure $\Sigma_{sc}(a,b) = S(a,b) = S(a$

and every $h \mapsto 0$ one has $\sum_{K,xc}(a,b)=0$ which yields $\sum_{(a,b)=0}^{\infty} (a,b)=0$. Therefore by lemma(1-2-9) one gets $E_{A_c}^{\infty}(a,b)=0$ which proves $E_{A_c}^{\infty}(a,b)=0$. If $|a,b|\setminus\Omega_{\alpha}M_{bk}$ is countable, then by $\Omega_{xc}|M_{bk}|\subseteq |a,b|\setminus\Omega_{\alpha c}|M_{bk}$ the set $\Omega_{xc}(M_{bk})$ is countable too which completes the proof (ii). (iii) If for each $E_{ac}(M_{bk})=0$ the condition (65) holds for a.e. $E_{ac}(M_{bk})=0$ each one has $E_{ac}(M_{bk})=0$ hence $E_{ac}(M_{bk})=0$ taking proposition (1-2-16) into account we find $E_{ac}(A_b)=0$. Conversely if $E_{ac}(A_b)=0$ then proposition (1-2-16) for each $E_{ac}(M_{bk})=0$ which yields $E_{ac}(M_{bk})=0$ then proposition (1-2-16) for each $E_{ac}(M_{bk})=0$.

 $\mathcal{S}(A_0) = \overline{\Omega(M;\tau)} = \mathcal{S}(A_0)$ which was proved in theorem (1-2-14)

Section (1-3): Characteristic Function for Row Contraction and Factorization

We give new result for the characteristic function associated with an arbitrary row contraction and show that

$$I - \mathbf{Q} \mathbf{\ddot{Q}} = \mathbf{k}_T \mathbf{k}_T^*$$

Where is the Poisson kernel of . Consequently we will show that the curvature invariant and characteristic associated with a Hilbert model over generated by an arbitrary row contraction can be expressed only in terms of the characteristic function of [62,63].

The characteristic function associated with an arbitrary row condition $T = \{T_1, ..., T_n\}$. $T_i = \{B(H)\}$ was introduced for the classical case and it was proved to be a complete unitary invariant for completely non-coisometric (c.n.c) row contraction. Using the characteristic of multi-analytic operator on Fock spaces, one can easily see that the characteristic function of an is multi-analytic operator [64, 65].

$$\mathcal{Q}(R_1,...,R_n): F^2(H_n) \otimes \mathcal{D}_{T^*} \longrightarrow F^2(H_n) \otimes \mathcal{D}_{T}$$

With the formal Fourier representation

$$-\!I_{_{F^2(H_n)}}\!\!\circ\! T \ + \! \Big(I_{_{F^2(H_n)}}\!\!\circ\! \Delta_{_{\! T}}\Big) \! \bigg(I_{_{F^2(H_n)}\otimes\mathsf{H}} - \! \sum_{i=1}^n \! R_i \otimes\! T_i^{\,*} \frac{1}{J}^{\!-1} \big[\, R_1 \otimes\! I_{\,\mathsf{H}},, R_n \otimes\! I_{\,\mathsf{H}} \,\big] \Big(I_{_{F^2(H_n)}} \otimes\! \Delta_{_{\! T}} \,\Big) = 0$$

where R_1, \dots, R_n are the right creation operator on the full Fock space $F^2(H_n)$

are $\Delta_T = \left(I_H - \sum_{i=1}^n T_i T_i^*\right)^{\frac{1}{p}} \in B(H)$ and $\Delta = \left(I_H - T^*T\right) \in B(H^{(n)})$ while the defect spaces are $D_T = \overline{\Delta} = \overline{A} = \overline{$

Lemma (1-3-1) [61]:

If
$$\Theta(R_1,...,R_n) = R_n = \Theta(R_1,...,R_n)$$
 if $SOT - \lim_{r \to 1} \Theta(rR_1,...,rR_n)^* = \Theta(R_1,...,R_n)^*$

Proof:

We know that any multi-analytic operator $e^{R_1,...,R_n}$ with formal Fourier representation $e^{R_1,...,R_n}$ $\sim \sum_{k=0}^{\infty} \sum_{|\alpha|=k} R_{\alpha} \otimes \theta_{\alpha} \quad \theta_{\alpha} \in B(\mathsf{H},k)$ has the property that $\theta(R_1,...,R_n) = SOT - \lim_{T \to k} \sum_{k=0}^{\infty} \sum_{|\alpha|=k} r^{|\alpha|} R_{\alpha} \otimes \theta_{\alpha}$ where the series

converge in the uniform norm for each r = (0,1) now note that for every $\beta = F_n^+, h = H$, and $g = F^2(H_n) \otimes k$ we have

$$\begin{split} \left\langle \boldsymbol{\varTheta}\!\left(\boldsymbol{R}_{\!\!1},\!...,\!\boldsymbol{R}_{\!\!n}\right)^{*}\!\left(\boldsymbol{e}_{\beta}\otimes\!\boldsymbol{h}\right),\boldsymbol{\mathsf{g}}\right\rangle \!=\! &\left\langle \boldsymbol{e}_{\beta}\otimes\!\boldsymbol{h},\boldsymbol{\varTheta}\!\left(\boldsymbol{R}_{\!\!1},\!...,\boldsymbol{R}_{\!\!n}\right)\boldsymbol{\mathsf{g}}\right\rangle \\ \\ &= &\left\langle \boldsymbol{e}_{\beta}\otimes\!\boldsymbol{h}\!\left(\sum\nolimits_{\alpha\!\in\!F_{\!n}^{+},|\alpha\!|\!\leq\!|\beta|}\boldsymbol{R}_{\alpha}\otimes\!\boldsymbol{\varTheta}_{\!\alpha}\right)\boldsymbol{\mathsf{g}}\right\rangle \\ \\ &= &\left\langle \left(\sum\nolimits_{\alpha\!\in\!F_{\!n}^{+},|\alpha\!|\!\leq\!|\beta|}\boldsymbol{R}_{\alpha}^{*}\otimes\!\boldsymbol{\varTheta}_{\!\alpha}^{*}\right)\left(\boldsymbol{e}_{\beta}\otimes\!\boldsymbol{h}\right)\boldsymbol{\mathsf{g}}\right. \end{split}$$

therefore

$$(\mathbf{R}_{1},...,\mathbf{R}_{n})^{*}(\mathbf{e}_{\beta}\otimes\mathbf{h})=\left(\sum_{\mathbf{E}_{n}^{+}|\alpha|\leq\beta}\mathbf{R}_{\alpha}^{*}\otimes\mathbf{P}_{\alpha}^{*}(\mathbf{e}_{\beta}\otimes\mathbf{h})\right)$$

similarly we have

$$\mathbf{A}(rR_1,...,rR_n)^*(e_{\beta}\otimes h) = \left(\sum_{\alpha \in F_n^+ \mid \alpha \mid \alpha \mid \beta} r^{\mid \alpha \nmid} R_{\alpha}^* \otimes \mathbf{A}(e_{\beta}\otimes h)\right)$$

Using the last two equalities we obtain

$$\lim_{r \to \mathbf{1}} \mathbf{\Theta}(rR_1, ..., rR_n)^* \left(e_{\mathbf{\beta}} \otimes \mathbf{h}\right) = \mathbf{\Theta}(R_1, ..., R_n)^* \left(e_{\mathbf{\beta}} \otimes \mathbf{h}\right)$$

for any $\beta \in F_n^+$ and h on the other hand according to non commutative von Neumann inequality $\| e^{-r} e^{-r} e^{-r} \| e^{-r} e^{-r} \|$ for any $e^{-r} e^{-r} e^{-r$

$$SOT - \lim_{r \to 1} \Theta(rR_1, ..., rR_n)^* = \Theta(R_1, ..., R_n)^*$$

The proof is complete. The following factorization play an important role

Theorem (1-3-2) [5]:

Let
$$T = [T_1, ..., T_n]$$
. $T_i = B(H)$ be a row contraction, then $I = QQ = k_T k_T^*$

(69)

where is the characteristic function of and is the corresponding Poisson kernel.

Proof:

Denoting
$$\widetilde{T} = [I_{F^2(H_n)} \otimes T_1, ..., I_{F^2(H_n)} \otimes T_n]$$
 and $\widetilde{R} = [R_1 \otimes R_1, ..., R_n \otimes R_H]$

the characteristic function of has representation

$$\theta(R_{1},...,R_{n}) = SOT - \lim_{r \to 1} \left[-T^{0/2} + \Delta_{1/2} \left(I_{F^{2}(H_{n})} - e\hat{R}T^{0/2} \right)^{-1} r\hat{R} \Delta_{1/2} \right]$$
(70)

define the operator

$$A = \widetilde{T}^*$$
, $B \Longrightarrow$, $C \Longrightarrow$, $D = \widetilde{T}$, and $z \Longrightarrow$.

and note that

define

And notice that using relation (71) we have **[5]:**

$$I \longrightarrow (I \longrightarrow A^*Z)^+ = I \longrightarrow D^* \longrightarrow (I \longrightarrow A^*Z)^+ = I \longrightarrow (I \longrightarrow A)^+ = I \longrightarrow$$

$$-\mathbf{i}(I - ZA)^{+}A^{*}Z^{*} - Z^{*}Z^{*} - \mathbf{i}ZAA^{*}Z^{*}(I - A^{*}Z^{*})^{+}C^{*}$$

$$= \mathbf{i}(I - ZA)^{+}(I - ZZ^{*})(I - A^{*}Z^{*})^{+}C^{*}$$

Therefore

$$(72)$$

Therefore according to our notation for any r=(0,1) the defect operator r=(0,1) the defect operator is equal to the product **[5]**:

$$= (I \otimes \Delta_{\Gamma})^{+} (I - r^{2} \hat{R}^{-} \hat{R}^{+}) (I - r^{2} \hat{R}^{-} \hat{R}^{+}) (I - r^{2} \hat{R}^{-} \hat{R}^{+}) (I - r^{2} \sum_{i=1}^{n} R_{i}^{*} \otimes I)^{+} (I - r^{2} \sum_{i=1}^{n} R_{i}^{*} \otimes I) (I - r^{2} \sum_{i=1}^{n} R_{i}^{*} \otimes I)^{-} (I + \otimes \Delta_{\Gamma})^{-} (I + \otimes \Delta_{\Gamma})^{-} (I - r^{2} \sum_{i=1}^{n} R_{i} R_{i}^{*}) \otimes I (I - r^{2} \sum_{i=1$$

Now for every $\alpha \omega \in F_n^+$, $h \in D_T$ we have

$$\langle [I - G(rR_1,...,rR_n)G(rR_1,...,rR_n)^*] (e_{\sim}), e_{\sim} \rangle$$

$$= \sum_{i=1}^{n} R_{i}R_{i}^{*} e_{\alpha} e_{\omega} \left\langle \Delta_{\mathbf{r}} T_{\beta}^{*} \Delta_{\mathbf{r}} h, k \right\rangle$$

using lemma (1-3-1) we have

$$SOT - \lim_{r \to 1} \mathbf{G}_{1} \left(rR_{1}, ..., rR_{n} \right)^{*} = \mathbf{G}_{1} \left(R_{1}, ..., R_{n} \right)^{*}$$

therefore the above computations imply that

$$egin{align*} & \left\langle \left[I - \mathcal{C}(R_1, ..., R_n) \mathcal{C}(R_1, ..., R_n)^* \right] e \mathcal{C}(R_1, ..., R_n)^* \right\rangle & = \sum_{F_n^+, |r| \leq |\alpha|} \sum_{\mathcal{C}(F_n^+, |\beta| \leq |\alpha|)} \left\langle R_T P_{\mathbf{f}} R_{\beta}^* e_{\alpha \alpha} e_{\alpha \alpha} \right\rangle \left\langle \Delta_1 T_{\beta \alpha}^* T_{\beta \alpha} \Delta_1 h, k \right\rangle & = \sum_{F_n^+, |r| \leq |\alpha|} \left\langle R_r(1), e_{\alpha \alpha} \right\rangle \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_r^* T_{\alpha \alpha} h, k \right\rangle & = \left\langle \Delta_1 T_r^* T_r^*$$

For any $h, k \in D_T$ and CCE_n^+ . Summing up the above computation we deduce that $I = CR_1, ..., R_n$ $CR_1, ..., R_n$ which complete the proof.

We recall that the spectral radius of an n-tuple of operators $x = [x_1, ..., x_n]$ is defined by

$$r(X) = \lim_{k \to \infty} \left\| \sum_{|\alpha| = k} X_{\alpha} X_{\alpha}^{*} \right\|^{\frac{1}{2}k}$$

A closer look at the proof of theorem (1-3-2) reveals the following factorization result. We should add that operator $r = \hat{x} \hat{r}$ invertible because

Corollary (1-3-3) [61]:

Let $T = \{T_1, ..., T_n\} T_i \rightleftharpoons (H)$ be a row contraction and let be its characteristic function. If $X = \{X_1, ..., X_n\} X_i \rightleftharpoons (K)$ is a row contraction with spectral radius $I_k \bowtie I_k \bowtie I_k$

Theorem (1.3.3) [5]:

Let $T \in B(H)$ be arrow counteraction, if $QQ + k_T \cdot k_T^* = I$, where is the characteristic function of I and I is the corresponding Poisson kernel, then for $A = I^{9/6}, B = 1/6$, and C = 0 the operator $S = \begin{bmatrix} A & B \\ 0 & A \end{bmatrix}$ commute with I if and only if I commute with I and I are self-adjoint.

Proof:

If
$$S^*S = \begin{bmatrix} A^*A & A^*B + B^*A \\ 0 & A^*A \end{bmatrix}$$
 and $SS^* = \begin{bmatrix} AA^* & AB^* + BA^* \\ 0 & AA^* \end{bmatrix}$

Equating, then $S^*S = SS^*$ if and only if $A^*A = AA^*$ is normal and $A^*B = BA^*$.

Theorem (1-3-4) [61]:

Let f_{i} be the complete free semigroup algebra generated by the free semigroup f_{i} with generators g_{i}, \dots, g_{n} and natural element f_{i} any g_{i} and g_{i} and natural element g_{i} and g_{i} are g_{i} and g_{i} and g_{i} and g_{i} and g_{i} are g_{i} and g_{i} and g_{i} and g_{i} and g_{i} and g_{i} are g_{i} and g_{i} are g_{i} and g_{i} and g_{i} are g_{i} and g_{i} and g_{i} are g_{i} and g_{i} and g_{i} are g_{i} and

and studied we recall $Cuv(T) = \lim_{m \to \infty} \frac{trace[I - \varphi_T^m(I)]}{I + n + ... + n^{m-1}}$ and

 $\chi(T) = \lim_{m \to \infty} \frac{rank[I - \phi_r^m(I)]}{1 + n + ... + n^{m-1}}$ where is the completely positive map associated

with i.e., $\phi_{\tau}(X) = \sum_{i=1}^{n} T_i \times T_i^*$ using theorem (1-3-2) and some results we can show that the curvature and the Euler characteristic of arrow contraction can be expressed only in terms of the standard characteristic functions $\frac{1}{2}$.

Theorem (1-3-5) [61]:

Let $T = (T_1, ..., T_n)$, $T_i \in B(H)$ contraction with rank $A_i < \infty$ and Cur(T) and A(T) denote its curvature and Euler characteristic respectively then

$$Curv(T) = rank \ \Delta_T - \lim_{m \to \infty} \frac{trace[\theta_T \theta_T^*(p_m \otimes I)]}{n^m} \quad \text{and} \quad \chi(T) = \lim_{m \to \infty} \frac{rank[(I - \theta_T \theta_T^*)P_{\leq m} \otimes I]}{1 + n + \dots + n^{m-1}}$$

Where $P_m(resp.P_{sim})$ in the orthogonal projection of full Fock space $F^2(H_n)$ onto the subspace of all homogenous polynomials of degree (resp-polynomials of degree sim).

Proof:

Since $Curv(T) = \lim_{m \to \infty} \frac{trace\left[K_T K_T^* \left(p_m \otimes I\right)\right]}{n^m}$. Using the factorization results of theorem(1-3-2) the first result follows

$$\chi(T) = \lim_{m \to \infty} \frac{\operatorname{rank}\left[K_T^*\left(P_{\leq m} \otimes I\right)K_T\right]}{1 + n + \dots + n^{m-1}} \tag{73}$$

Since $K_T^*(P_{\leq in} \otimes I)$ has finite rank we have $rankK_T^*(P_{sin} \otimes I) = rank[ran k_T^*(P_{sin} \otimes I)]$ on the other hand ,since K_T is one-torange of $K_T^*(P_{sin} \otimes P)$ the we also on have one $rank[K_T^*(P_{sin} \bowtie)] = rank[K_T k_T^*(P_{sin} \bowtie)]$. Hence using relation (73) and theorem (1-3-2) we complete the proof.

A constrained characteristic function is associated with any constrained row contraction. For pure constrained row contractions we show that this characteristic function is a complete unitary invariant and provide a model in terms of it. We also show that Arveson's curvature invariant and Euler characteristic associated with a Hilbert model over CIZ,..., generated by a commuting row contraction can be expressed only in terms of the constrained characteristic function.

Let T be a wot-closed two —sided ideal of the non commutative analytic Toeplitz algebra T generated by a family of polynomials T we define the constrained characteristic function associated with a J-constrained row contraction $T = \{T_1, \dots, T_n\} T_1 = \{B(H)\}$ to be the multi-analytic operator (with respect to the constrained shift $\{B_1, \dots, B_n\}$

Defined by the formal Fourier representation

Taking into account that ** a co-invariant subspace under ** can see that ** is the maximal J-constrained pieced of the standard characteristic function ** of the row contraction ** . More precisely we have

$$(74)$$

$$(N_{J} \otimes D_{T}) \otimes N_{J} \otimes D_{T} \cdot and$$

$$P_{N_{J} \otimes D_{T}} \otimes (R_{1},...,R_{n}) N_{J} \otimes D_{T} \cdot \mathcal{A}_{T} ((W_{1},...,W_{n}))$$

Let us remark that the above definition of the constrained characteristic function makes sense when τ is an arbitrary wot-closed two sided ideal of and $\tau = [\tau_1, \dots, \tau_n]$ is an arbitrary c.n.c constrained row contraction.

Theorem (1-3-6) [61]:

Let be wot-closed two seded ideal of generated by a family of polynomials.

$$T = \{T_1, \dots, T_n\} T_i \in \mathcal{B}(H)$$
, a J-constrained row contraction then
$$I_{NJ \otimes D_T} - e_{J,T} = K_{J,T} K_{J,T}^*$$
(75)

where θ_{x} is the constrained characteristic function of $\frac{1}{2}$ and $\frac{K_{xx}}{2}$ is the corresponding constrained Poisson Kernel

Proof:

The constrained Poisson Kernel associated with is $\kappa_{xx} + - \kappa_{xx} + - \kappa_{$

$$K_{J,T} = \left(P_{N_J} \otimes I_{\overline{\Delta_T} H} K_T\right) \tag{76}$$

 $I_{N_J \otimes P_T} - P_{N_J \otimes P_T} \mathcal{A}(R_1, ..., R_n) \mathcal{A}(R_1, ..., R_n) | N_J \otimes P_T = P_{N_J \otimes P_T} K_T K_T K_{T^*} | N_J \otimes P_T$ taking into account relation (74) and (76) and that $W_i^* = R_i^* | N_J$, i = 1, ..., n we refer that $I_{N_J \otimes P_T} - \mathcal{A}_T (W_1, ..., W_n) \mathcal{A}_T (W_1, ..., W_n)^* = K_{J,T} K_{J,T}^*$ as in the proof of

theorem (1-3-5) one can use corollary (1-3-3) to obtain the following constrained version of it.

Corollary (1-3-7) [61]:

Let $J = F_n$ be closed sided ideal of F_n generated by a family of polynomials F_n and let $T = \{T_1, \dots, T_n\} T_i = B(H)$, be a J-constrained row contraction. If $X = \{X_1, \dots, X_n\}$, $X_i \in B(K)$, is a J-constrained row contraction with spectral radius $F(X) = \{X_1, \dots, X_n\}$, then

$$I_{k \otimes D_{T}} - \mathbf{Q}_{T}(X_{1}, ..., X_{n}) \mathbf{Q}_{T}(X_{1}, ..., X_{n})^{*} = \mathbf{\hat{X}} \left(I - \hat{X} \widetilde{T}^{*}\right)^{-1} \left(I - \widetilde{X} \hat{X}^{*}\right) \left(I - \widetilde{T} \hat{X}^{*}\right)^{-1} \mathbf{\hat{A}}$$

where $x \to x_1 \to x_2 \to x_3$ and the other notations are from the proof of theorem (1-3-2). Now we present a model for pure constrained row contraction in terms of characteristic function.

Theorem (1-3-8) [61]:

Let $J \neq F_n^-$ be a wot-closed two-sided ideal of F_n and $T = [T_1, ..., T_n]$, $T_i \in B(H)$, be a pure J-constrained row contraction. Then the constrained characteristic function $P_n = W(W_1, ..., W_n) \supseteq (D_{T^-}, D_T)$ is a partial isometry and P_n is unitarily equivalent to the row contraction

$$[P_{H_{j,T}}(B_1 \bowtie_{D_T})H_{j,T},...,P_H(B_n \bowtie_{D_T})H_{J,T}]$$

(77)

where $P_{H_{J,T}}$ is the orthogonal projection of $N_J \otimes D_T$ on the Hilbert space $H_{J,T} = (N_J \otimes D_T) \otimes Q_T (N_J \otimes D_T)$

Theorem (1-3-9) [61]:

Let $T = [T_1, ..., T_n]$, be a wot-closed two-sided ideal of T_i and let $T = [T_1, ..., T_n]$, $T_i \in B(H), T = [T_1, ..., T_n]$ be two J-constrained pure contractions then $T_i = [T_1, ..., T_n]$ are unitarily equivalent if and only if their constrained characteristic function $T_i = [T_1, ..., T_n]$ and $T_i = [T_1, ..., T_n]$.

Proof:

Assume that $T_i = U^*T_iU$ for any $T_i = U^*T_iU$ simple computation

reveal that $U \rightarrow -\infty U$ and $\left| \left(\left(\left(\int_{i:1}^{n} U \right) \left| \left(\int_{T^*}^{n} \cdot \cdot \cdot \cdot \cdot \right) \left| \left(\int_{T^*}^{n} \cdot \cdot \cdot \cdot \cdot \cdot \cdot \right) \right| \right| \right|$ define the unitary operator

and by setting
$$T = U \mid D_T : D_T \longrightarrow T$$
 and $T' = \left| \emptyset \mid_{i=1}^n U \middle| D_{T^*} \rightarrow D_{T^*} \right|$ Taking

into account the definition of the constrained characteristic function it is easy to see that

$$(I_{N_{I}} \otimes T)\theta_{I,T} = \theta_{I,T}(I_{N_{I}} \otimes T')$$

Conversely assume that the characteristic function of and coincide. According to the remarks preceding the theorem there exist unitary operators

 $r: D_r \longrightarrow D_{r'}$ and $r_*: D_{r'} \to D_{r'}$ such that the following diagram

$$\begin{array}{ccc} N_{j} \otimes D_{T^{*}} & \xrightarrow{\Phi_{j,T}} & N_{j} \otimes D_{T} \\ & \downarrow I_{N_{j}} \otimes \tau^{*} & \downarrow I_{N_{j}} \otimes \tau \\ N_{j} \otimes D_{T^{'}} & \xrightarrow{\Phi_{j,T^{'}}} & N_{j} \otimes D_{T^{'}} \end{array}$$

Is commutative i.e.,

(78)

Hence we deduce that $\Gamma = (I_{N_j} \otimes P)H_{j,T} : H_{j,T} \longrightarrow H_{j,T'}$ is a unitary operator where $H_{j,T}$ and $H_{j,T}$ are the model spaces for and respectively. Since $(B_i^* \otimes I_{D_T})(I_{N_j} \otimes T^*) = (I_{N_j} \otimes T^*)(B_i^* \otimes I_{D_T}), i = 1,...,n$ and $H_{j,T}(resp.H_{j,T'})$ is a coinvariant subspace under $B_i \otimes I_{D_T}(respp.B_i \otimes I_{D_T}), i = 1,...,n$ we deduce that

Hence we obtain

$$\mathbf{IP}_{H_{J,T}}\left(B_{i}^{*} \otimes \mathbf{I}_{D_{T}}\right)\left(H_{J,T}\right) = \mathbf{P}_{H_{J,T}}\left(B_{i} \otimes \mathbf{I}_{D_{T'}}\right)\left(H_{J,T'}\right)\mathbf{\Gamma}, i = 1,..., n$$

Now using theorem (1-3-7) we conclude that and are unitary equivalent. The proof is complete.

Theorem (1-3-10) [61]:

Let be a WOT-closed two- sided ideal of F such that $I \in N_x$ and condition $\overline{span}\{B_xB_x: \alpha B \in F_n^+\} = c^*(B_1, ..., B_n)$ is satisfied. If $M \in N_x$ is an invariant subspace under $B_1, ..., B_n$ and $T = \{F_1, ..., T_n\} T_i = \{F_{n-1}B_i \mid M^{-\frac{1}{2}}, i = \{F_{n-1}B_i \mid M^{-$

$$K_n(z, w) = \frac{1}{1 - \langle z, w \rangle \mathbf{f}^n}, z, w \in B_n$$

The algebra $w_n \to P_{n-1}$ was proved to be the r-closed algebra generated by the operators $B_1, I \to \dots, n$ and the identity. Moreover w_n can be identified with the algebra of all multiplies of P under this identification the creation operator B_1, \dots, B_n become the multiplication operator M_{n_1}, \dots, M_{n_n} by the coordinate functions w_1, \dots, w_n of w_n . Let $w_1 \to w_1, \dots, w_n$ be a constrained row contraction i.e., w_1, \dots, w_n under the above-mentioned identifications the constrained characteristic function of w_n the multiplier (multiplication operator)

$$\Theta_{T_c}$$
, $T: H^2 \otimes D_{T^*} \longrightarrow H^2 \otimes D_T$

defined by the operator. Valued analytic function on the open unit ball

$$B_{n} = \left\{ z = (z_{1}, ..., z_{n}) \in \mathbf{f}_{n} : |z| = (|z|^{2} + ... + |z|^{2})^{\frac{1}{2}} < 1 \right\}$$

given by

$$Q_{cT}(z) = T + \Delta \left(I - z_1 T_1^* - \dots - z_n T_n^*\right)^{-1} \left[z_1 I_H, \dots, z_n I_H\right] \Delta \cdot , z \in B_n$$

We going to use the same notation for the multiplication operator $M_{\Theta_{C,T}} \in B(H^2 \otimes D_T, D_T)$ and its symbol $\theta_{I,T}$ which is a $B(D_T, D_T)$ valued bounded analytic function in $\frac{1}{2}$. All the results of this part can be written in this commutative setting. Using theorem (1-3-5) and corollary (1-3-6) and some result we show that Euler characteristic associated with a commutative row contraction $\frac{1}{2}$ with rank $\Delta_T < \infty$ can be expressed in terms of the constrained characteristic function $\theta_{I,T}$.

Theorem (1-3-11) [61]:

Let $T = \{T_1, \dots, T_n\} T_i$ \Longrightarrow (H). be a commutative row contraction with rank $A \iff$ and let K(T) and $A \iff$ denote Arveson's curvature and Euler characteristic respectively.

Then

$$\begin{split} K(T) = & \int_{\partial B_{n}} \lim_{r \to \mathbf{i}} \operatorname{trace} \left[I - \mathbf{\theta}_{J_{c},T} \left(r \, \boldsymbol{\zeta} \right) \, \mathbf{\theta}_{J_{c},T} \left(r \, \boldsymbol{\zeta} \right)^{*} \, \right] \! d \, \partial (\, \boldsymbol{\zeta}) \\ = & \operatorname{rank} \, \Delta_{T} - \left(n - 1 \right) \lim_{m \to \infty} \frac{\operatorname{trace} \left[\left(\mathbf{\theta}_{J_{c},T} \, \mathbf{\theta}_{J_{c},T}^{*} \, \left(\, Q_{m} \, \otimes I_{D_{T}} \, \right) \, \right]}{n^{m}} \end{split}$$

where a is the projection of a onto the subspace of homogeneous polynomials of degree and

$$\chi(T) = n! \lim_{m \to \infty} \Delta_{T} - (n-1) \frac{rank \left[\left(1 - \theta_{J_{C},T} \theta_{J_{C},T}^{*} \right) \left(Q_{\leq m} \otimes I_{D_{T}} \right) \right]}{m^{n}}$$

Where Q_a is the projection of $^{H^2}$ onto the subspace of all polynomials of degree sm .

Proof:

Using the factorization result of corollary (1-3-6) in our particular case we obtain

$$\left(1-\mathbf{G}_{c,T}(z)\mathbf{G}_{c,T}(z)^*\right) = \left(1-|z|^2\right) \mathbf{I}_{c,T}(z)^* \mathbf{I}_{c,T}(z)^*$$

for any $z \in B_n$ the first formula follows from the definition of the curvature and the above-mentioned factorization for the constrained characteristic function of $z \in B_n$. Using (60) and we have.

$$K(T) = (n-1) \lim_{r \to \mathbb{N}} \frac{\operatorname{trace} \left[\left(P_{\scriptscriptstyle m} \otimes I \right) K_{\scriptscriptstyle T} K_{\scriptscriptstyle T}^* \right]}{m^{n-1}}$$

Where κ_r is the Poisson on Kernel of and is the orthogonal projection of $F^2(H_n)$ onto the subspace of all homogeneous polynomials of degree. Since is a commutative row contraction—constrained we have range $\kappa_T C F_s^2 \otimes D_T$ —and the constrained Poisson Kernel satisfies the equation $\kappa_{J_C,T} = (P_{F_s^2} \otimes I) \kappa_T$ —where κ_T is the symmetric Fock space. Using the standard properties. For the trace and above relation we deduce that

$$K(T) = (n-1) \lim_{r \to 1} \frac{\operatorname{trace}\left[K_{J_{cT}} K_{J_{cT}}^* \left(Q_m \otimes I\right)\right]}{m^{n-1}}$$
(79)

where $Q_m = P_{F_s^2} P_m | F_s^2$ is the projection of F_s onto the subspace of homogenous polynomials of degree. According to theorem (1-3-5) we have

$$I - \Theta_{c_T} \Theta_{c_T} = K_{J_{c_T}} K_{J_{c_T}}^*$$

$$\tag{80}$$

taking into account relation (79),(80) we deduce the second formula for the curvature. Here of course we used Arveson's identification of symmetric Fock space with his space and Arveson's showed that his Euler

characteristic satisfies the equation $\chi(T) = n! \lim_{m \to \infty} \frac{rank \left[1 - \varphi_r^{m+1}(1)\right]}{m^n}$ we here is the completely positive map associated with we get

$$\chi(T) = n! \lim_{m \to \infty} \frac{rank \left[K_T^* \left(P_{\leq m} \otimes I \right) K_T \right]}{m^n}$$
(81)

Where P_{sm} is the orthogonal projection of $F_s^2(H_n)$ on the subspace of all polynomials of degree P_sm . Using again that range P_sm and the

contained Poisson Kernel satisfies the equation $K_{J_c,T} = (P_{F_s^2} \otimes I)K_T$ we deduce that

$$\begin{aligned} rank \left[K_T^* (P_{\leq m} \oplus I) K_T \right] &= rank \left[K_T^* (P_{\leq m} \oplus I) \left(P_{F_s^2} \otimes I \right) K_T \right] \\ &= rank \left[K_T^* (Q_{\leq m} \oplus I) \left(P_{F_s^2} \otimes I \right) K_{J_C, T} \right] \\ &= rank \left[K_{J_C, T}^* (Q_{\leq m} \oplus I) \left(P_{F_s^2} \otimes I \right) \right] \\ &= rank \left[K_{J_C, T} K_{J_C, T}^* (Q_{\leq m} \oplus I) \right] \end{aligned}$$

Where Q_{nn} is the projection of P_n onto the subspace of all polynomials of degree P_n the last two equalities hold since the operator $K_{J_c,T}^*(Q_{nn})$ has finite rank and $P_{J_c,T}$ is one -to-one on the range of $P_{J_c,T}^*(Q_{nn})$. Now using relation (81) we obtain the last formula of the theorem. The proof is complete.

Let $[T_1,...,T_n]$, $T_i \rightleftharpoons (H)$ be a pure row contraction and let I be a wot-closed two-sided ideal of I such that $4F_1,...,T_n] \rightleftharpoons for any$

$$[S_1,...,S_n] = I$$

(82)

where $A_i,...,T_n$] is defined using the F_i functional calculus for row contraction is unitary equivalent to the compression of $B_i \bowtie ... \bowtie B_k \bowtie 1$ to a co-invariant subspace under each operator $B_i \bowtie ... \bowtie n$ therefore we have the $T_i = P \in (B_i \bowtie I_k)$, i = 1,...,n following result is a commutative lifting theorem for pure constrained row contraction.

Theorem (1-3-12) [61]:

Let $J \neq F_n^-$ be a WOT-closed two- sided ideal of the non commutative analytic Toeplitz algebra F_n , and $[B_1,...,B_n]$ and $[W_1,...,W_n]$ be the corresponding constrained shifts acting on F_n . For each $J \Rightarrow 2$ let A be a

Hilbert space and $\varepsilon_j \subseteq N_j \otimes K_j$ be a co-invariant subspace under each operator $B_i \otimes R_j : i \longrightarrow N_j$ if $X : \varepsilon_i \longrightarrow \varepsilon_2$ is a bounded such operator such that

$$X_{\epsilon_{\!\scriptscriptstyle i}}\!\left[P\!\left(B_{\!\scriptscriptstyle i}\otimes\!\! I_{k_{\!\scriptscriptstyle i}}\right)\right|_{\epsilon_{\!\scriptscriptstyle i}}\left]\!\!=\!\!P_{\epsilon_{\!\scriptscriptstyle 2}}\!\left(B_{\!\scriptscriptstyle i}\otimes\!\! I_{k_{\!\scriptscriptstyle 2}}\right)\!\right|_{\epsilon_{\!\scriptscriptstyle 2}}X,i=\!\!1,...,n$$

then there exists

$$G(W_1,...,W_n) \rightleftharpoons W(W_1,...,W_n) \rightleftharpoons (k_1,k)$$

(83) such that

$$G(W_1,...,W_n)^* = X^*$$
 and $||G(W_1,...,W_n)| = X|$

In particular if \longrightarrow where is a co-invariant subspace under each operator $\xrightarrow{B_i, i \longrightarrow \dots, n}$ then the above implication becomes an equivalence.

Corollary (1-3-13) [61]:

Let $S = \mathbb{R}^n$ be a wot-closed two- sided ideal of the non commutative analytic Toeplitz algebra \mathbb{R}^n and let B_1, \dots, B_n and W_1, \dots, W_n be the corresponding constrained shift acting on \mathbb{R}^n . If \mathbb{R}^n is a Hilbert space and \mathbb{R}^n is an invariant subspace under each operator \mathbb{R}^n and \mathbb{R}^n , \mathbb{R}^n we remark that theorem (1-3-10) can be extended to the following more general setting. The proof follows exactly the same lines so we shall omit it. For each \mathbb{R}^n let \mathbb{R}^n be a wot-closed two- sided ideal of \mathbb{R}^n and let \mathbb{R}^n be an invariant subspace under each operator \mathbb{R}^n and \mathbb{R}^n . Let \mathbb{R}^n be an invariant subspace under each operator \mathbb{R}^n is a bounded operator such that \mathbb{R}^n A Hilbert space. If \mathbb{R}^n is a bounded operator such that \mathbb{R}^n \mathbb{R}^n is a bounded operator such that \mathbb{R}^n \mathbb{R}^n is a bounded operator. Such that \mathbb{R}^n is a bounded operator such that \mathbb{R}^n is a bounded operator.

Theorem (1-3-14) [61]:

Let be a wot-closed two- sided ideal of F_i and let B_1, \dots, B_n be the corresponding constrained shift acting on F_i . Let A_1, \dots, A_k be distinct point in the zero set $Z_j = \{A \in B_n : f(A) = 0 \text{ for any } f \in J\}$ and let $A_1, \dots, A_k = B(k)$ then there exists $A_1, \dots, A_k = A_k$ where $A_1, \dots, A_k = A_k$ if and only if the operator matrix

$$\left[\frac{I_{k}-\!\!A_{\!i}A_{\!j}^{*}}{I-\!\!\left\langle\boldsymbol{\lambda}_{\!r},\boldsymbol{\lambda}_{\!j}\right\rangle}\right]_{\!\!k\!>\!\!k}$$

(84)

is positive semi definite

Proof:

$$(W_1,...,W_n) (W_1,...,W_n) M = \mathbb{F}^*$$

(85)

and $\| \langle w_1, ..., w_n \rangle \| = \| v_1 \|$ one can prove that $\varphi(\lambda_j) = A_j$, j = 1, ..., k if and only if (85) holds. Moreover $\| \langle w_1, ..., w_n \rangle \| = v_1 \| v_1 \| v_2 \| v_3 \|$ which is

equivalent to the fact that the operator matrix (84) is positive semi definite. This completes the proof.