Sudan University of science and Technology College of Graduate studies


# The Analytical techniques on manifold via differential forms 



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By:
Maab Ibrahim Alamin Ali
Supervisor:
Dr. Emad Eldeen Abdallah Abdelrahim

## 

قال تعالى :


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## Dedication

I dedicate this humble work to my father who was always standing ready to satisfy all my needs, and to my mother who covered me with her tenderness and love, and to all other members of my family for their patience and encouragement. I will never forget to whisper in the ears of every person who taught me even a letter through my life that I am greatly indebted to all of them.

## Acknowledgement

Thanks God, the graceful who had given me the will and talent to carry on till I reached this successful situation. Moreover, I send my thanks to my supervisor, Dr. Emad Eldeen Abdallah Abdelrahim, to whom I didn't find the suitable words to express my sincere respect and appreciation. Lastly my thanks extended to all my lecturers of Sudan university of science and technology to whom I owe this exerted successful effort .


#### Abstract

In this research, we show by using integration of different form that, for a compact orientable manifold of dimension n the De Rham cohomology group $H^{n}(M)$ is non zero, we also show that the group for a compact, connected, orientable manifold is just one-dimensional.

We discuss the metric tensor and Riemannian metric on a manifold in informal terms. We also illustrate the relation between the integral curve of geodesic flow and geodesic with some examples and applications.


## الخلاصة

فى هذا البحث بينا بإستخدام تكامل الصيغة التفاضليه ان زمره كو همولوجيا ديرام ${ }^{\text {ان }}$ (Mn الدورانى المتر اص من البعد (n) تكون غير صفريه. و اوضحنا ان هذه الزمر لمتعدد الطيات الدورانى المنصل المتراص هي فقط من البعد واحد ـ نـاقثنـا الممتد المتري و مترك ريمان علي متعدد الطيات و ذللك بإسلوب غير رسمى . أيضا أوضحنا العلاقه بين منحنى التكامل لانسياب الجيدوسيك و الجيدوسيك مع بعض الامثلهه و التطبيقات.

## Contents

| subject | Number |
| :---: | :---: |
| الاية | I |
| Dedication | II |
| Acknowledgement | III |
| Abstract (English) | IV |
| Abstract (Arabic) | V |
| Contents | VI |
| Introduction | VII |
| Chapter (1) <br> Manifold, Tangent and Cotangent Spaces |  |
| Section(1.1):Basic Concept of Manifold | 1 |
| Section(1.2):Tangent Vector and Cotangent Vectors | 13 |
| Chapter (2)Vector fields and Lie Bracket |  |
| Section(2.1):Vector Fields and Lie Bracket | 24 |
| Section(2.2):Tensor products and Exterior Algebra | 38 |
| Chapter (3) <br> Differential Forms, De Rham Cohomology and Stokes' Theorem |  |
| Section(3.1):Differential Forms and De Rham Cohomology | 46 |
| Section(3.2):Forms Integration and Stokes' Theorem | 71 |
| Chapter (4) <br> Smooth Map and Rimannian Metric |  |
| Section(4.1):The Degree of Smooth Map | 87 |
| Section(4.2):The Metric Tensor | 97 |
| Section(4.3):The Geodesic Flow | 102 |
| References | 111 |

## Introduction

Today the area of differential geometry becomes one in which recent developments have effected great changes. However, the studying of its material has been relatively little affected by many modern developments. this research deal with some of this developments, and organized as follows:-

In chapter(1), we discuss the basic concepts of manifold and its general mathematical structures with some examples and applications. Also, we study the tangent vector, tangent spaces, cotangent vector and cotangent spaces with some remarks and examples .

In chapter(2), we deal with a brief introduction to tangent bundle via our definition of vector fields and their associated mathematical structures. Also we discuss the Lie bracket .

In chapter(3), we study the differential forms and De Rham cohomology groups. And to get more informations on de Rham cohomology, we study the orientation, integration of forms and Stokes' Theorem with some examples and applications.

In chapter(4), we study the degree of smooth maps, De Rham cohomology in the top dimension, and the concepts of the orientable $n$-dimensional manifold. Also we discuss the metric tensor and Geodesic flow with some examples and applications.

## Chapter (1)

## Manifolds, Tangent and Contingent Spaces

## Section (1.1): Basic Concept of Manifold

The concept of a manifold is a bit complicated, but it starts with defining the notion of a coordinate chart.

## Definition (1.1.1):

A coordinate chart on a set X is a subset $U \subseteq X$ together with a bijection

$$
\varphi: U \rightarrow \varphi(U) \subseteq R^{n}
$$

onto an open set $\varphi(\mathrm{U})$ in $R^{n}$.
Thus we can parametrize points $x$ of $U$ by n coordinates $\varphi(x)=\left(x_{1}, \ldots, x_{n}\right)$.
We now want to consider the situation where $X$ is covered by such charts and satisfies some consistency conditions. We have

## Definition (1.1.2):

An $n$-dimensional atlas on $X$ is a collection of coordinate charts $\left\{U_{\alpha}, \varphi_{\alpha}\right\}_{\alpha \in I}$ such that
i. $X$ is covered by the $\left\{U_{\alpha}\right\}_{\alpha \in I}$
ii. for each $\alpha \in I, \varphi_{\alpha}\left(U_{\alpha} \cap U_{\beta}\right)$ is open in $R^{n}$
iii. the map

$$
\varphi_{\beta} \varphi_{\alpha}^{-1}: \varphi_{\alpha}\left(U_{\alpha} \cap U_{\beta}\right) \rightarrow \varphi_{\beta}\left(U_{\alpha} \cap U_{\beta}\right)
$$

is $C^{\infty}$ with $C^{\infty}$ inverse.

Recall that $F\left(x_{1}, \ldots, x_{n}\right) \in R^{n}$ is $C^{\infty}$ if it has derivatives of all orders. We shall also say that $F$ is smooth in this case. It is perfectly possible to develop the theory of manifolds with less differentiability than this, but this is the normal procedure.

## Examples (1.1.3):

1. Let $X=R^{n}$ and take $U=X$ with $\varphi=i d$. We could also take $X$ to be any open set in $R^{n}$.
2. Let $X$ be the set of straight lines in the plane:


Each such line has an equation $A x+B y+C=0$ where if we multiply $A, B, C$ by a non-zero real number we get the same line. Let $U_{0}$ be the set of non-vertical lines.For each line $\ell \in U_{0}$ we have the equation

$$
y=m x+c
$$

where $m, c$ are uniquely determined. So $\varphi_{0}(\ell)=(m, c)$ defines a coordinate chart $\varphi_{0}: U_{0} \rightarrow R^{2}$. Similarly if $U_{1}$ consists of the non-horizontal lines with equation

$$
x=\tilde{m} y+\tilde{c}
$$

we have another chart $\varphi_{1}: U_{1} \rightarrow R^{2}$

Now $U_{0} \cap U_{1}$ is the set of lines $y=m x+c$ which are not horizontal, so $m \neq 0$. Thus

$$
\varphi_{0}\left(U_{0} \cap U_{1}\right)=\left\{(m, c) \in R^{2}: m \neq 0\right\}
$$

which is open. Moreover, $y=m x+c$ implies $x=m^{-1} y-c m^{-1}$ and so

$$
\varphi_{1} \varphi_{0}^{-1}(m, c)=\left(m^{-1},-c m^{-1}\right)
$$

which is smooth with smooth inverse. Thus we have an atlas on the space of lines.
3. Consider $R$ as an additive group, and the subgroup of integers $Z \subset R$. Let $X$ be the quotient group $R / Z$ and $p: R \rightarrow R / Z$ the quotient homomorphism.

Set $U_{0}=p(0,1)$ and $U_{1}=p(-1 / 2,1 / 2)$. Since any two elements in the subset $p^{-1}(a)$ differ by an integer, $p$ restricted to $(0,1)$ or $(-1 / 2,1 / 2)$ is injective and so we have coordinate charts

$$
\varphi_{0}=p^{-1}: U_{0} \rightarrow(0,1), \varphi_{1}=p^{-1}: U_{1} \rightarrow(-1 / 2,1 / 2)
$$

Clearly $U_{0}$ and $U_{1}$ cover $\mathrm{R} / \mathrm{Z}$ since the integer $0 \in U_{1}$.

We check:

$$
\varphi_{0}\left(U_{0} \cap U_{1}\right)=(0,1 / 2) \cup(1 / 2,1), \varphi_{1}\left(U_{0} \cap U_{1}\right)=(-1 / 2,0) \cup(0,1 / 2)
$$

Which are open sets . Finally, if $x \in(0,1 / 2), \varphi_{1} \varphi_{0}^{-1}(x)=x$ and if $\mathrm{x} \in$ $(1 / 2,1), \varphi_{1} \varphi_{0}^{-1}(x)=x-1$. These maps are certainly smooth with smooth inverse so we have an atlas on $X=R / Z$.
4. Let X be the extended complex plane $X=C \cup\{\infty\}$. Let $U_{0}=C$ with $\varphi_{0}(z)=z \in C \cong R^{2}$. Now take

$$
U_{1}=C \backslash\{0\} \cup\{\infty\}
$$

and define $\varphi_{1}(\tilde{z})=\tilde{z}^{-1} \in C$ if $\tilde{z} \neq \infty$ and $\varphi_{1}(\infty)=0$. Then

$$
\varphi_{0}\left(U_{0} \cap U_{1}\right)=C \backslash\{0\}
$$

which is open, and

$$
\varphi_{1} \varphi_{0}^{-1}(z)=z^{-1}=\frac{x}{x^{2}+y^{2}}-i \frac{y}{x^{2}+y^{2}}
$$

This is a smooth and invertible function of $(x, y)$. We now have a 2-dimensional atlas for $X$, the extended complex plane.
5. Let $X$ be $n$-dimensional real projective space, the set of 1 -dimensional vector subspaces of $R^{n+1}$. Each subspace is spanned by a non-zero vector $v$, and we define $U_{i} \subset R P^{n}$ to be the subset for which the $i$-th component of $v \in R^{n+1}$ is non-zero. Clearly $X$ is covered by $U_{1}, \ldots, U_{n+1}$. In $U_{i}$ we can uniquely choose $v$ such that the $i$ th component is 1 , and then $U_{i}$ is in one-to-one correspondence with the hyperplane $x_{i}=1$ in $R^{n+1}$, which is a copy of $\boldsymbol{R}^{n}$. This is therefore a coordinate chart

$$
\varphi_{i}: U_{i} \rightarrow R^{n}
$$

The set $\varphi_{i}\left(U_{i} \cap U_{j}\right)$ is the subset for which $x_{j} \neq 0$ and is therefore open. Furthermore

$$
\varphi_{i} \varphi_{j}^{-1}:\left\{x \in R^{n+1}: x_{j}=1, x_{i} \neq 0\right\} \rightarrow\left\{x \in R^{n+1}: x_{i}=1, x_{j} \neq 0\right\}
$$

is

$$
v \mapsto \frac{1}{x_{i}} v
$$

which is smooth with smooth inverse. We therefore have an atlas for $\boldsymbol{R} P^{n}$.

Now we will discuss the definition of a manifold, all the examples above are actually manifolds, and the existence of an atlas is sufficient to establish that, but there is a minor subtlety in the actual definition of a manifold due to the fact that there are lots of choices of atlases. If we had used a different basis for $\boldsymbol{R}^{2}$, our charts on the space $X$ of straight lines would be different, but we would like to think of $X$ as an object independent of the choice of atlas. That's why we make the following definitions:

## Definition (1.1.4):

Two atlases $\left\{\left(U_{\alpha}, \varphi_{\alpha}\right)\right\},\left\{\left(V_{i}, \psi_{i}\right)\right\}$ are compatible if their union is an atlas.

What this definition means is that all the extra maps $\psi_{i} \varphi_{\alpha}^{-1}$ must be smooth. Compatibility is clearly an equivalence relation, and we then say that:

## Definition (1.1.5):

A differentiable structure on $X$ is an equivalence class of atlases.

Finally we come to the definition of a manifold:

## Definition (1.1.6):

An $n$-dimensional differentiable manifold is a space $X$ with a differentiable structure.

The upshot is this: to prove something is a manifold, all you need is to find one atlas. The definition of a manifold takes into account the existence of many more atlases.

Many books give a slightly different definition -they start with a topological space, and insist that the coordinate charts are homeomorphisms. This is fine if we see the
world as a hierarchy of more and more sophisticated structures but it suggests that in order to prove something is a manifold we first have to define a topology. As we'll see now, the atlas does that for us.

First recall what a topological space is: a set $X$ with a distinguished collection of subsets $V$ called open sets such that
i. $\quad$ and $X$ are open
ii. An arbitrary union of open sets is open
iii. A finite intersection of open sets is open

Now suppose $M$ is a manifold. We shall say that a subset $V \subseteq M$ is open if, for each $\alpha, \varphi_{\alpha}\left(V \cap U_{\alpha}\right)$ is an open set in $R^{n}$. One thing which is immediate is that $V=U_{\beta}$ is open, from Definition (1.1.2).

We need to check that this gives a topology. Condition 1 holds because $\varphi_{\alpha}(\phi)=$ $\phi$ and $\varphi_{\alpha}\left(M \cap U_{\alpha}\right)=\varphi_{\alpha}\left(U_{\alpha}\right)$ which is open by Definition (1.1.1) for the other two, if $V_{i}$ is a collection of open sets then because $\varphi_{\alpha}$ is bijective

$$
\begin{aligned}
\varphi_{\alpha}\left(\left(\cup V_{i}\right) \cap U_{\alpha}\right) & =\cup \varphi_{\alpha}\left(V_{i} \cap U_{\alpha}\right) \\
\varphi_{\alpha}\left(\left(\cap V_{i}\right) \cap U_{\alpha}\right) & =\cap \varphi_{\alpha}\left(V_{i} \cap U_{\alpha}\right)
\end{aligned}
$$

and then the right hand side is a union or intersection of open sets. Slightly less obvious is the following:

## Proposition (1.1.7):

With the topology above $\varphi_{\alpha}: U_{\alpha} \rightarrow \varphi_{\alpha}\left(U_{\alpha}\right)$ is a homeomorphism.

## Proof:

If $V \subseteq U_{\alpha}$ is open in the induced topology on $U_{\alpha}$ then since $U_{\alpha}$ itself is open, $V$ is open in $M$. Then $\varphi_{\alpha}(V)=\varphi_{\alpha}\left(V \cap U_{\alpha}\right)$ is open by the definition of the topology, so $\varphi_{\alpha}^{-1}$ is certainly continuous.

Now let $W \subset \varphi_{\alpha}\left(U_{\alpha}\right)$ be open, then $\varphi_{\alpha}^{-1}(W) \subseteq U_{\alpha}$ so we need to prove that $\varphi_{\alpha}^{-1}(W)$ is open in . But

$$
\begin{equation*}
\varphi_{\beta}\left(\varphi_{\alpha}^{-1}(W) \cap U_{\beta}\right)=\varphi_{\beta} \varphi_{\alpha}^{-1}\left(W \cap \varphi_{\alpha}\left(U_{\alpha} \cap U_{\beta}\right)\right) \tag{1.1}
\end{equation*}
$$

From Definition (1.1.2) the set $\varphi_{\alpha}\left(U_{\alpha} \cap U_{\beta}\right)$ is open and hence its intersection with the open set $W$ is open. Now $\varphi_{\beta} \varphi_{\alpha}^{-1}$ is $C^{\infty}$ with $C^{\infty}$ inverse and so certainly a homeomorphism, and it follows that the right hand side of (1.1) is open. Thus the left hand side $\varphi_{\beta}\left(\varphi_{\alpha}^{-1} W \cap U_{\beta}\right)$ is open and by the definition of the topology this means that $\varphi_{\alpha}^{-1}(W)$ is open, i.e. $\varphi_{\alpha}$ is continuous.

To make any reasonable further progress, we have to make two assumptions about this topology which will hold for the rest of search:
i. the manifold topology is Hausdorff
ii. in this topology we have a countable basis of open sets

Without these assumptions, manifolds are not even metric spaces, and there is not much analysis that can reasonably be done on them.

Now we will discuss further examples of manifolds, we need better ways of recognizing manifolds than struggling to find explicit coordinate charts. For example, the sphere is a manifold

and although we can use stereographic projection to get an atlas:

there are other ways. Here is one.

## Theorem (1.1.8):

Let $F: U \rightarrow R^{m}$ be a $C^{\infty}$ function on an open set $U \subseteq R^{n+m}$ and take $c \in R^{m}$. Assume that for each $a \in F^{-1}(c)$, the derivative

$$
D F_{a}: R^{n+m} \rightarrow R^{m}
$$

is surjective. Then $F^{-1}(c)$ has the structure of an n -dimensional manifold which is Hausdorff and has a countable basis of open sets.

## Proof:

Recall that the derivative of F at a is the linear map $D \mathrm{~F}_{\mathrm{a}}: R^{n+m} \rightarrow R^{m}$ such that

$$
F(a+h)=F(a)+D F_{a}(h)+R(a, h)
$$

Where

$$
R(a, h) /\|h\| \rightarrow 0 \text { as } h \rightarrow 0
$$

If we write $F\left(x_{1}, \ldots, x_{n+m}\right)=\left(F_{1}, \ldots, F_{m}\right)$ the derivative is the Jacobian matrix

$$
\frac{\partial F_{i}}{\partial x_{j}}(a) \quad 1 \leq i \leq m, 1 \leq j \leq n+m
$$

Now we are given that this is surjective, so the matrix has rank m. Therefore by reordering the coordinates $x_{1}, \ldots, x_{n+m}$ we may assume that the square matrix

$$
\frac{\partial F_{i}}{\partial x_{j}}(a) \quad 1 \leq i \leq m, 1 \leq j \leq m
$$

is invertible.

Now define

$$
G: U \rightarrow R^{n+m}
$$

by

$$
\begin{equation*}
G\left(x_{1}, \ldots, x_{n+m}\right)=\left(F_{1}, \ldots, F_{m}, x_{m+1} \ldots, x_{n+m}\right) \tag{1.2}
\end{equation*}
$$

Then $D G_{a}$ is invertible.

We now apply the inverse function theorem to $G$, a proof of which is given in the Appendix. It tells us that there is a neighbourhood $V$ of $a$, and $W$ of $G(a)$ such that $G: V \rightarrow W$ is invertible with smooth inverse. Moreover, the formula (1.2) shows that $G$ maps $V \cap F^{-1}(c)$ to the intersection of $W$ with the copy of $R^{n}$ given $\operatorname{by}\left\{x \in R^{n+m}: x_{i}=c_{i}, 1 \leq i \leq m\right\}$. This is therefore a coordinate chart $\varphi$. If we take two such charts $\varphi_{\alpha}, \varphi_{\beta}$ then $\varphi_{\alpha} \varphi_{\beta}^{-1}$ is a map from an open set in $\left\{x \in R^{n+m}: x_{i}=c_{1}, 1 \leq i \leq m\right\}$ to another one which is the restriction of the $\operatorname{map} G_{\alpha} G_{\beta}^{-1}$ of (an open set in) $R^{n+m}$ to itself. But this is an invertible $C^{\infty}$ map and so we have the requisite conditions for an atlas.

Finally, in the induced topology from $R^{n+m}, G_{\alpha}$ is a homeomorphism, so open sets in the manifold topology are the same as open sets in the induced topology. Since $R^{n+m}$ is Hausdorff with a countable basis of open sets, so is $F^{-1}(c)$.

Effectively (1.2) gives a coordinate chart on $R^{n+m}$ such that $F^{-1}(c)$ is a linear subspace there: we are treating $R^{n+m}$ as a manifold in its own right.

We can now give further examples of manifolds:

## Examples (1.1.9):

1. Let

$$
S^{n}=\left\{x \in R^{n+1}: \sum_{1}^{n+1} x_{i}^{2}=1\right\}
$$

be the unit n-sphere. Define $F: R^{n+1} \rightarrow R$ by

$$
F(x)=\sum_{1}^{n+1} x_{i}^{2}
$$

This is a $C^{\infty}$ map and

$$
D F_{a}(h)=2 \sum_{i} a_{i} h_{i}
$$

is non-zero (and hence surjective in the 1 -dimensional case) so long as $a$ is not identically zero. If $F(a)=1$, then

$$
\sum_{1}^{n+1} a_{i}^{2}=1 \neq 0
$$

so $a \neq 0$ and we can apply Theorem (1.1.8) and deduce that the sphere is a manifold.
2. Let $O(n)$ be the space of $n \times n$ orthogonal matrices: $A A^{T}=I$. Take the vector space $M_{n}$ of dimension $n^{2}$ of all real $n \times n$ matrices and define the function

$$
F(A)=A A^{T}
$$

to the vector space of symmetric $n \times n$ matrices. This has dimension $n(n+1) / 2$ then $O(n)=F^{-1}(I)$.

Differentiating F we have

$$
D F_{A}(H)=H A^{T}+A H^{T}
$$

and putting $H=K A$ this is

$$
K A A^{T}+A A^{T} K^{T}=K+K^{T}
$$

if $A A^{T}=I$, i.e. if $A \in F^{-1}(I)$. But given any symmetric matrix S , taking $K=S / 2$ shows that $D F_{I}$ is surjective and so, applying Theorem (1.1.8) we find that $O(n)$ is a manifold. Its dimension is

$$
n^{2}-n(n+1) / 2=n(n-1) / 2
$$

Now we will discuss maps between manifolds, we need to know what a smooth map between manifolds is. Here is the definition:

## Definition (1.1.10):

A map $F: M \rightarrow N$ of manifolds is a smooth map if for each point $x \in M$ and chart $\left(U_{\alpha}, \varphi_{\alpha}\right)$ in M with $x \in U_{\alpha}$ and chart $\left(V_{i}, \psi_{i}\right)$ of N with $F(x) \in V_{i}$, the set $F^{-1}\left(V_{i}\right)$ is open and the composite function

$$
\psi_{i} F \varphi_{\alpha}^{-1}
$$

on $\varphi_{\alpha}\left(F^{-1}\left(V_{i}\right) \cap U_{\alpha}\right)$ is a $C^{\infty}$ function.

Note that it is enough to check that the above holds for one atlas - it will follow from the fact that $\varphi_{\alpha} \varphi_{\beta}^{-1}$ is $C^{\infty}$ that it then holds for all compatible atlases.

The natural notion of equivalence between manifolds is the following:

## Definition (1.1.11):

A diffeomorphism $F: M \rightarrow N$ is a smooth map with smooth inverse.

## Example (1.1.12)

Take two of our examples above - the quotient group $R / Z$ and the 1 -sphere, the circle, $S^{1}$. We shall show that these are eomorphic. First we define a map

$$
G: R / Z \rightarrow S^{1}
$$

by

$$
G(x)=(\cos 2 \pi x, \sin 2 \pi x)
$$

This is clearly a bijection. Take $x \in U_{0} \subset R / Z$ then we can represent the point by $x \in(0,1)$. Within the range $(0,1 / 2), \sin 2 \pi x \neq 0$, so with $F=x_{1}^{2}+x_{2}^{2}$, we have $\partial F / \partial x_{2} \neq 0$. The use of the inverse function theorem in Theorem (1.1.8) then says that $x_{1}$ is a local coordinate for $S^{1}$, and in fact on the whole of $(0,1 /$ 2) $\cos 2 \pi x$ is smooth and invertible. We proceed by taking the other similar open sets to check fully that $G$ is a smooth, bijective map. To prove that its inverse is smooth, we can rely on the inverse function theorem, since $\sin 2 \pi x \neq 0$ in the interval.

## Section (1.2): Tangent Vectors and Cotangent Vectors

We begin this section by study the existence of smooth functions, the most fundamental type of map between manifolds is a smooth map

$$
f: M \rightarrow R
$$

We can add these and multiply by constants so they form a vector space $C^{\infty}(M)$, the space of $C^{\infty}$ functions on $M$. In fact, under multiplication it is also a commutative ring. So far, all we can assert is that the constant functions lie in this space, so let's see why there are lots and lots of global $C^{\infty}$ functions. We shall use bump functions and the Hausdorff property.

First note that the following function of one variable is $C^{\infty}$ :

$$
\begin{aligned}
f(t) & =e^{-\frac{1}{t}} & & t>0 \\
& =0 & & t \leq 0
\end{aligned}
$$

Now form

$$
g(t)=\frac{f(t)}{f(t)+f(1-t)}
$$

so that g is identically 1 when $t \geq 1$ and vanishes if $t \leq 0$. Next write

$$
h(t)=g(t+2) g(2-t)
$$

This function vanishes if $|t| \geq 2$ and is 1 where $|t| \leq 1$ : it is completely at on top.


Finally make an $n$-dimensional version

$$
k\left(x_{1}, \ldots, x_{n}\right)=h\left(x_{1}\right) h\left(x_{2}\right) \ldots h\left(x_{n}\right)
$$

In the sup norm, this is 1 if $|x| \leq 1$, so $k\left(r^{-1} x\right)$ is identically 1 in a ball of radius r and 0 outside a ball of radius 2 r .

We shall use this construction several times later on. For the moment, let $M$ be any manifold and $\left(U, \varphi_{U}\right)$ a coordinate chart. Choose a function $k$ of the type above whose support (remember supp $f=\overline{\{x: f(x) \neq 0\}})$ lies in $\varphi_{U}(U)$ and define

$$
f: M \rightarrow R
$$

by

$$
\begin{array}{r}
f(x)=k \circ \varphi_{U}(x) \quad x \in U \\
=0 \quad x \in M \backslash U
\end{array}
$$

is this a smooth function? The answer is yes: by definition $f$ is smooth for points in the coordinate neighbourhood $U$. But supp $k$ is closed and bounded in $R^{n}$ and so
compact and since $\varphi_{U}$ is a homeomorphism, $f$ is zero on the complement of a compact set in $M$. But a compact set in a Hausdorff space is closed, so its complement is open. If $y \neq U$ then there is a neighbourhood of $y$ on which f is identically zero, in which case clearly $f$ is smooth at $y$.

Now we will discuss the derivative of a function, smooth functions exist in abundance. The question now is: we know what a diferentiable function is - so what is its derivative? We need to give some coordinate independent definition of derivative and this will involve some new concepts. The derivative at a point $a \in M$ will lie in a vector space $T_{a}^{*}$ called the cotangent space.

First let's address a simpler question - what does it mean for the derivative to vanish? This is more obviously a coordinate-invariant notion because on a compact manifold any function has a maximum, and in any coordinate system in a neighbourhood of that point, its derivative must vanish. We can check that: if $f: M \rightarrow R$ is smooth then the composition

$$
g=f \varphi_{\alpha}^{-1}
$$

is $\mathrm{a} C^{\infty}$ function of $x_{1}, \ldots, x_{n}$. Suppose its derivative vanishes at $\varphi_{\alpha}(a)=\left(x_{1}(a), \ldots, x_{n}(a)\right)$ and now take a erent chart $\varphi_{\beta}$ with $h=f \varphi_{\beta}^{-1}$. Then

$$
g=f \varphi_{\alpha}^{-1}=f \varphi_{\beta}^{-1} \varphi_{\beta} \varphi_{\alpha}^{-1}=h \varphi_{\beta} \varphi_{\alpha}^{-1}
$$

But from the definition of an atlas, $\varphi_{\beta} \varphi_{\alpha}^{-1}$ is smooth with smooth inverse, so

$$
g\left(x_{1}, \ldots, x_{n}\right)=h\left(y_{1}(x), \ldots, y_{n}(x)\right)
$$

and by the chain rule

$$
\frac{\partial g}{\partial x_{i}}=\sum_{j} \frac{\partial h}{\partial y_{j}}(y(x)) \frac{\partial y_{j}}{\partial x_{i}}(x) .
$$

Since $y(x)$ is invertible, its Jacobian matrix is invertible, so that $D g_{x}(a)=0$ if and only if $D h_{y}(a)=0$. We have checked then that the vanishing of the derivative at a point a is independent of the coordinate chart. We let $Z_{a} \subset C^{\infty}(M)$ be the subset of functions whose derivative vanishes at a. Since $D f_{a}$ is linear in $f$ the subset $Z_{a}$ is a vector subspace.

## Definition (1.2.1):

The cotangent space $T_{a}^{*}$ at $a \in M$ is the quotient space

$$
T_{a}^{*}=C^{\infty}(M) / Z_{a}
$$

the derivative of a function $f$ at a is its image in this space and is denoted $(d f)_{a}$.
Here we have simply defined the derivative as all functions modulo those whose derivative vanishes. It's almost a tautology, so to get anywhere we have to prove something about $T_{a}^{*}$. First note that if $\varphi$ is a smooth function on a neighbourhood of $x$, we can multiply it by a bump function to extend it to $M$ and then look at its image in $T_{a}^{*}=C^{\infty}(M) / Z_{a}$. But its derivative in a coordinate chart around $a$ is independent of the bump function, because all such functions are identically 1 in a neighbourhood of $a$. Hence we can actually define the derivative at $a$ of smooth functions which are only defined in a neighbourhood of $a$. In particular we could take the coordinate functions $x_{1}, \ldots, x_{n}$. We then have

## Proposition (1.2.2):

Let $M$ be an $n$-dimensional manifold, then
i. the cotangent space $T_{a}^{*}$ at $a \in M$ is an $n$-dimensional vector space
ii. if $(U, \varphi)$ is a coordinate chart around x with coordinates $x_{1}, \ldots, x_{n}$, then the elements $\left(d x_{1}\right) a, \ldots\left(d x_{n}\right) a$ form a basis for $T_{a}^{*}$.
iii. if $f \in C^{\infty}(M)$ and in the coordinate chart, $f \varphi^{-1}=\phi\left(x_{1}, \ldots, x_{n}\right)$ then

$$
\begin{equation*}
(d f)_{a}=\sum_{i} \frac{\partial \phi}{\partial x_{i}}(\varphi(a))\left(d x_{i}\right)_{a} \tag{1.3}
\end{equation*}
$$

## Proof:

If $f \in C^{\infty}(M)$, with $f \varphi^{-1}=\phi\left(x_{1}, \ldots, x_{n}\right)$ then

$$
f-\sum \frac{\partial \phi}{\partial x_{i}}(\varphi(a)) x_{i}
$$

is a (locally defined) smooth function whose derivative vanishes at a, so

$$
(d f)_{a}=\sum \frac{\partial f}{\partial x_{i}}(\varphi(a))\left(d x_{i}\right)_{a}
$$

and $\left(d x_{1}\right) a, \ldots\left(d x_{n}\right) a \quad \operatorname{span} T_{a}^{*}$.
If $\sum_{i} \lambda_{i}\left(d x_{i}\right)_{a}=0$ then $\sum_{i} \lambda_{i} x_{i}$ has vanishing derivative at $a$ and so $\lambda_{i}=0$ for all $i$.

## Remark (1.2.3):

It is rather heavy handed to give two symbols $f, \phi$ for a function and its representation in a given coordinate system, so often in what follows we shall use just $f$. Then we can write (1.3) as

$$
d f=\sum \frac{\partial \emptyset}{\partial x_{i}} d x_{i}
$$

With a change of coordinates $\left(x_{1}, \ldots, x_{n}\right) \rightarrow\left(y_{1}(x), \ldots, y_{n}(x)\right)$ the formalism gives

$$
d f=\sum_{j} \frac{\partial f}{\partial y_{j}} d y_{j}=\sum_{i, j} \frac{\partial f}{\partial y_{j}} \frac{\partial y_{j}}{\partial x_{i}} d x_{i}
$$

## Definition (1.2.4):

The tangent space $T_{a}$ at $a \in M$ is the dual space of the cotangent space $T_{a}^{*}$.

This is admittedly a roundabout way of defining $T_{a}$, but since the double dual $\left(V^{*}\right)^{*}$ of a finite dimensional vector space is naturally isomorphic to $V$ the notation is consistent. If $x_{1}, \ldots, x_{n}$ is a local coordinate system at a and $\left(d x_{1}\right)_{a}, \ldots\left(d x_{n}\right)_{a}$ the basis of $T_{a}^{*}$ a defined in (1.2.2) then the dual basis for the tangent space Ta is denoted

$$
\left(\frac{\partial}{\partial x_{1}}\right)_{a}, \ldots,\left(\frac{\partial}{\partial x_{n}}\right)_{a}
$$

This definition at first sight seems far away from our intuition about the tangent space to a surface in $\boldsymbol{R}^{3}$.


The problem arises because our manifold $M$ does not necessarily sit in Euclidean space and we have to define a tangent space intrinsically. There are two ways around this: one would be to consider functions $f: R \rightarrow M$ and equivalence classes of these, instead of functions the other way $f: M \rightarrow R$. Another, perhaps more useful, one is provided by the notion of directional derivative. If $f$ is a function on a surface in $\boldsymbol{R}^{3}$, then for every tangent direction $u$ at $a$ we can define the derivative of $f$ at $a$ in the direction $u$, which is a real number: $u$. $\nabla f(a)$ or $D f_{a}(u)$. Imitating this gives the following:

## Definition (1.2.5):

A tangent vector at a point $a \in M$ is a linear map $X_{a}: C^{\infty}(M) \rightarrow R$ such that

$$
X_{a}(f g)=f(a) X_{a} g+g(a) X_{a} f
$$

This is the formal version of the Leibnitz rule for differentiating a product.
Now if $\xi \in T_{a}$, it lies in the dual space of $T_{a}^{*}=C^{\infty}(M) / Z_{a}$ and so

$$
f \mapsto \xi\left((d f)_{a}\right)
$$

is a linear map from $C^{\infty}(M)$ to $R$. Moreover from (1.3)

$$
d(f g)_{a}=f(a)(d g)_{a}+g(a)(d f)_{a}
$$

and so

$$
X_{a}(f)=\xi\left((d f)_{a}\right)
$$

is a tangent vector at $a$. In fact, any tangent vector is of this form, but the price paid for the nice algebraic definition in (1.2.5) which is the usual one in textbooks is that we need a lemma to prove it.

## Lemma (1.2.6):

Let $X_{a}$ be a tangent vector at $a$ and $f$ a smooth function whose derivative at a vanishes. Then $X_{a} f=0$.

## Proof:

Use a coordinate system near . By the fundamental theorem of calculus,

$$
f(x)-f(a)=\int_{0}^{1} \frac{\partial}{\partial t} f(a+t(x-a)) d t
$$

$$
=\sum_{i}\left(x_{i}-a_{i}\right) \int_{0}^{1} \frac{\partial f}{\partial x_{i}}(a+t(x-a)) d t
$$

If $(d f)_{a}=0$ then

$$
g_{i}(x)=\int_{0}^{1} \frac{\partial f}{\partial x_{i}}(a+t(x-a)) d t
$$

vanishes at $x=a$, as does $h_{i}(x)=x_{i}-a_{i}$. Now although these functions are defined locally, using a bump function we can extend them to $M$, so that

$$
\begin{equation*}
f=f(a)+\sum_{i} g_{i} h_{i} \tag{1.4}
\end{equation*}
$$

where $g_{i}(a)=h_{i}(a)=0$.
By the Leibnitz rule

$$
X_{a}(1)=X_{a}(1.1)=2 X_{a}(1)
$$

which shows that $X_{a}$ annihilates constant functions. Applying the rule of (1.4)

$$
X_{a}(f)=X_{a}\left(\sum_{i} g_{i} h_{i}\right)=\sum_{i}\left(g_{i}(a) X_{a} h_{i}+h_{i}(a) X_{a} g_{i}\right)=0
$$

This means that $X_{a}: C^{\infty}(M) \rightarrow R$ annihilates $Z_{a}$.

Now if $V \subset W$ are vector spaces then the annihilator of $V$ in the dual space $W^{*}$ is naturally the dual of $/ V$. So a tangent vector, which lies in the dual of $C^{\infty}(M)$ is naturally a subspace of $\left(C^{\infty}(M) / Z_{a}\right)^{*}$ which is, by our definition, the tangent $\operatorname{space}_{a}$.

The vectors in the tangent space are therefore the tangent vectors as defined by (1.2.5). Locally, in coordinates, we can write

$$
X_{a}=\sum_{i}^{n} c_{i}\left(\frac{\partial}{\partial x_{i}}\right)_{a}
$$

and then

$$
\begin{equation*}
X_{a}(f)=\sum_{i} c_{i} \frac{\partial f}{\partial x_{i}}(a) \tag{1.5}
\end{equation*}
$$

Now we will study the derivatives of smooth maps, suppose $F: M \rightarrow N$ is a smooth map and $f \in C^{\infty}(N)$. Then $f \circ F$ is a smooth function on $M$.

## Definition (1.2.7):

The derivative at $a \in M$ of the smooth map $F: M \rightarrow N$ is the homomorphism of tangent spaces

$$
D F_{a}: T_{a} M \rightarrow T_{F(a)} N
$$

defined by

$$
D F_{a}\left(X_{a}\right)(f)=X_{a}(f \circ F)
$$

This is an abstract, coordinate-free definition. Concretely, we can use (1.5) to see that

$$
\begin{gathered}
D F_{a}\left(\frac{\partial}{\partial x_{i}}\right)_{a}(f)=\frac{\partial}{\partial x_{i}}(f \circ F)(a) \\
=\sum_{j} \frac{\partial F_{j}}{\partial x_{i}}(a) \frac{\partial f}{\partial y_{j}}(F(a))=\sum_{j} \frac{\partial F_{j}}{\partial x_{i}}(a)\left(\frac{\partial}{\partial y_{j}}\right)_{F(a)} f
\end{gathered}
$$

Thus the derivative of $F$ is an invariant way of defining the Jacobian matrix.

With this definition we can give a generalization of Theorem (1.1.8) - the proof is virtually the same and is omitted.

Theorem (1.2.8):

Let $F: M \rightarrow N$ be a smooth map and $c \in N$ be such that at each point $a \in F^{-1}(c)$ the derivative $D F_{a}$ is surjective. Then $F^{-1}(c)$ is a smooth manifold of dimension $\operatorname{dim} M-\operatorname{dimN}$.

In the course of the proof, it is easy to see that the manifold structure on $F^{-1}(c)$ makes the inclusion

$$
\mathrm{l}: F^{-1}(c) \subset M
$$

a smooth map, whose derivative is injective and maps isomorphically to the kernel of $D F$. So when we construct a manifold like this, its tangent space at a is

$$
T_{a} \cong \operatorname{Ker} D F_{a}
$$

This helps to understand tangent spaces for the case where $F$ is defined on $\boldsymbol{R}^{n}$ :

## Examples (1.2.9):

1. The sphere $S^{n}$ is $F^{-1}(1)$ where $F: R^{n+1} \rightarrow R$ is given by

$$
F(x)=\sum_{i} x_{i}^{2}
$$

So here

$$
D F_{a}(x)=2 \sum_{i} x_{i} a_{i}
$$

and the kernel of $D F_{a}$ consists of the vectors orthogonal to $a$, which is our usual vision of the tangent space to a sphere.
2. The orthogonal matrices $O(n)$ are given by $F^{-1}(I)$ where $F(A)=A A^{T}$. At $A=I$, the derivative is

$$
D F_{I}(H)=H+H^{T}
$$

so the tangent space to $O(n)$ at the identity matrix is $\operatorname{Ker} D F_{I}$, the space of skew symmetric matrices $H=-H^{T}$.

The examples above are of manifolds $F^{-1}(c)$ sitting inside $M$ and are examples of
submanifolds. Here we shall adopt the following definition of a submanifold, which is often called an embedded submanifold.

## Definition (1.2.10):

A manifold $M$ is a submanifold of $N$ if there is an inclusion map

$$
\mathrm{\imath}: M \rightarrow N
$$

such that
i. t is smooth.
ii. $\quad D \mathrm{l}_{x}$ is injective for each $x \rightarrow M$.
iii. the manifold topology of $M$ is the induced topology from $N$.

## Remark (1.2.11):

The topological assumption avoids a situation like this:

$$
\mathbf{l}(t)=\left(t^{2}-1, t\left(t^{2}-1\right)\right) \in R^{2}
$$

for $t \in(-1, \infty)$. This is smooth and injective with injective derivative: it is the part of the singular cubic $y^{2}=x^{2}(x+1)$ consisting of the left hand loop and the part in the first quadrant. Any open set in $R^{2}$ containing 0 intersects the curve in a $t$-interval $(-1,-1+\delta)$ and an interval $(1-\delta, 1+\delta)$. Thus $(1-\delta, 1+$ $\delta$ ) on its own is not open in the induced topology.


## Chapter (2)

## Vector fields and Tensor Product

## Section (2.1) Vector fields and Lie Bracket

We begin with a brief introduction to tangent bundle. Think of the wind velocity at each point of the earth.

This is an example of a vector field on the 2 -sphere $S^{2}$. Since the sphere sits inside $\boldsymbol{R}^{3}$, this is just a smooth map $X: S^{2} \rightarrow R^{3}$ such that $X(x)$ is tangential to the sphere at $x$.

Our problem now is to define a vector field intrinsically on a general manifold $M$, without reference to any ambient space. We know what a tangent vector at $a \in M$ is - a vector in $T_{a}$ - but we want to describe a smoothly varying family of these. To do this we need to fit together all the tangent spaces as $a$ ranges over $M$ into a single manifold called the tangent bundle. We have $n$ degrees of freedom for $a \in M$ and $n$ for each tangent space $T_{a}$ so we expect to have a $2 n$-dimensional manifold. So the set to consider is

$$
T M=\mathrm{U}_{x \in M} T_{x}
$$

the disjoint union of all the tangent spaces.
First let $\left(U, \varphi_{U}\right)$ be a coordinate chart for $M$. Then for $x \in U$ the tangent vectors

$$
\left(\frac{\partial}{\partial x_{1}}\right)_{x}, \ldots,\left(\frac{\partial}{\partial x_{n}}\right)_{x}
$$

provide a basis for each $T_{x}$. So we have a bijection

$$
\varphi_{U}: U \times R^{n} \rightarrow \cup_{x \in U} T_{x}
$$

defined by

$$
\varphi_{U}\left(x, y_{1}, \ldots, y_{n}\right)=\sum_{1}^{n} y_{i}\left(\frac{\partial}{\partial x_{i}}\right)_{x}
$$

Thus

$$
\Phi_{U}=\left(\varphi_{U}, i d\right) \circ \psi_{U}^{-1}: \cup_{x \in U} T_{x} \rightarrow \varphi_{U}(U) \times R^{n}
$$

is a coordinate chart for

$$
V=U_{x \in U} T_{x}
$$

Given $U_{\alpha}, U_{\beta}$ coordinate charts on $M$, clearly

$$
\Phi_{\alpha}\left(V_{\alpha} \cap V_{\beta}\right)=\varphi_{\alpha}\left(U_{\alpha} \cap U_{\beta}\right) \times R^{n}
$$

which is open in $R^{2 n}$. Also, if $\left(x_{1}, \ldots, x_{n}\right)$ are coordinates on $U_{\alpha}$ and $\left(\tilde{x}_{1}, \ldots, \tilde{x}_{n}\right)$ on $U_{\beta}$ then

$$
\left(\frac{\partial}{\partial x_{i}}\right)_{x}=\sum_{j} \frac{\partial \tilde{x}_{j}}{\partial x_{i}}\left(\frac{\partial}{\partial \tilde{x}_{j}}\right)_{x}
$$

the dual of $\mathrm{Eq}(1.3)$. It follows that

$$
\Phi_{\beta} \Phi_{\alpha}^{-1}\left(x_{1}, \ldots, x_{n}, y_{1}, \ldots, y_{n}\right)=\left(\tilde{x}_{1}, \ldots, \tilde{x}_{n}, \sum_{i} \frac{\partial \tilde{x}_{1}}{\partial x_{i}} y_{i}, \ldots, \sum_{i} \frac{\partial \tilde{x}_{n}}{\partial x_{i}} y_{i}\right)
$$

and since the Jacobian matrix is smooth in $x$, linear in $y$ and invertible, $\Phi_{\beta} \Phi_{\alpha}^{-1}$ is smooth with smooth inverse and so $\left(V_{\alpha}, \Phi_{\alpha}\right)$ defnes an atlas on $T M$.

## Definition (2.1.1):

The tangent bundle of a manifold $M$ is the $2 n$-dimensional differentiable structure on $T M$ defined by the above atlas.

The construction brings out a number of properties. First of all the projection map

$$
p: T M \rightarrow M
$$

which assigns to $X_{a} \in T_{a} M$ the point a is smooth with surjective derivative, because in our local coordinates it is defined by

$$
p\left(x_{1}, \ldots, x_{n}, y_{1}, \ldots, y_{n}\right)=\left(x_{1}, \ldots, x_{n}\right)
$$

The inverse image $p^{-1}(a)$ is the vector space $T_{a}$ and is called a fibre of the projection. Finally, $T M$ is Hausdorff because if $X_{a}, X_{b}$ lie in different fibres, since $M$ is Hausdorff we can separate $a, b \in M$ by open sets $U, U^{\prime}$ and then the open sets $p^{-1}(U), p^{-1}\left(U^{\prime}\right)$ separate $X_{a}, X_{b}$ in $T M$. If $X_{a}, Y_{b}$ are in the same tangent space then they lie in a coordinate neighbourhood which is homeomorphic to an open set of $R^{2 n}$ and so can be separated there. Since $M$ has a countable basis of open sets and $R^{n}$ does, it is easy to see that $T M$ also has a countable basis.

We can now define a vector field:

## Definition (2.1.2):

A vector field on a manifold is a smooth map

$$
X: M \rightarrow T M
$$

such that

$$
p \circ X=i d_{M}
$$

This is a clear global definition. What does it mean? We just have to spell things out in local coordinates. Since $p \circ X=i d_{M}$,

$$
X\left(x_{1}, \ldots, x_{n}\right)=\left(x_{1}, \ldots, x_{n}, y_{1}(x), \ldots, y_{n}(x)\right)
$$

where $y_{i}(x)$ are smooth functions. Thus the tangent vector $X(x)$ is given by

$$
X(x)=\sum_{i} y_{i}(x)\left(\frac{\partial}{\partial x_{i}}\right)_{x}
$$

which is a smoothly varying field of tangent vectors.

## Remark (2.1.3):

We shall meet other manifolds $Q$ with projections $p: Q \rightarrow M$ and the general terminology is that a smooth map $s: M \rightarrow Q$ for which $p \circ s=i d_{M}$ is called a section. When $Q=T M$ is the tangent bundle we always have the zero section given by the vector field $X=0$. Using a bump function we can easily construct other vector fields by taking a coordinate system, some locally defined smooth functions $y_{i}(x)$ and writing

$$
X(x)=\sum_{i} y_{i}(x)\left(\frac{\partial}{\partial x_{i}}\right)_{x}
$$

Multiplying by $\psi$ and extending gives a global vector field.

## Remark (2.1.4):

Clearly we can do a similar construction using the cotangent spaces $T_{a}^{*}$ instead of the tangent spaces $T_{a}$, and using the basis

$$
\left(d x_{1}\right)_{x}, \ldots,\left(d x_{n}\right)_{x}
$$

instead of the dual basis

$$
\left(\frac{\partial}{\partial x_{1}}\right)_{x}, \ldots,\left(\frac{\partial}{\partial x_{1}}\right)_{x}
$$

This way we form the cotangent bundle $T^{*} M$. The derivative of a function $f$ is then a map $d f: M \rightarrow T^{*} M$ satisfying $p \circ d f=i d_{M}$, though not every such map of this form is a derivative.

Perhaps we should say here that the tangent bundle and cotangent bundle are examples of vector bundles. Here is the general definition:

## Definition (2.1.5):

A real vector bundle of rank m on a manifold $M$ is a manifold $E$ with a smooth projection $p: E \rightarrow M$ such that
i. each fibre $p^{-1}(x)$ has the structure of an $m$-dimensional real vector space
ii. each point $x \in M$ has a neighbourhood $U$ and a diffieomorphism

$$
\psi_{U}: p^{-1}(U) \cong U \times R^{m}
$$

such that $\psi_{U}$ is a linear isomorphism from the vector space $p^{-1}(x)$ to the vector space $\{x\} \times R^{m}$
iii. on the intersection $U \cap V$

$$
\psi_{U} \psi_{v}^{-1}: U \cap V \times R^{m} \rightarrow U \cap V \times R^{m}
$$

is of the form

$$
(x, v) \mapsto\left(x, g_{U V}(x) v\right)
$$

where $g_{U V}(x)$ is a smooth function on $U \cap V$ with values in the space of invertible $m \times m$ matrices.

For the tangent bundle $g_{U V}$ is the Jacobian matrix of a change of coordinates and for the cotangent bundle, its inverse transpose.

Now we will study vector fields as derivations. The algebraic definition of tangent vector in Definition (1.2.5) shows that a vector field $X$ maps a $C^{\infty}$ function to a function on $M$ :

$$
X(f)(x)=X_{x}(f)
$$

and the local expression for $X$ means that

$$
X(f)(x)=\sum_{i} y_{i}(x)\left(\frac{\partial}{\partial x_{i}}\right)_{x}(f)=\sum_{i} y_{i}(x) \frac{\partial f}{\partial x_{i}}(x)
$$

Since the $y_{i}(x)$ are smooth, $X(f)$ is again smooth and satisfies the Leibnitz property

$$
X(f g)=f(X g)+g(X f) .
$$

In fact, any linear transformation with this property (called a derivation of the algebra $\left.C^{\infty}(M)\right)$ is a vector field:

## Proposition (2.1.6):

Let $X: C^{\infty}(M) \rightarrow C^{\infty}(M)$ be a linear map which satisfies

$$
X(f g)=f(X g)+g(X f):
$$

Then $X$ is a vector field.

## Proof:

For each $a \in M, X_{a}(f)=X(f)(a)$ satisfies the conditions for a tangent vector at a, so $X$ defines a map $X: M \rightarrow T M$ with $p \circ X=i d_{M}$, and so locally can be written as

$$
X_{x}=\sum_{i} y_{i}(x)\left(\frac{\partial}{\partial x_{i}}\right)_{x}
$$

We just need to check that the $y_{i}(x)$ are smooth, and for this it suffices to apply $X$ to a coordinate function $x_{i}$ extended by using a bump function in a coordinate neighbourhood. We get

$$
X x_{i}=y_{i}(x)
$$

and since by assumption $X$ maps smooth functions to smooth functions, this is smooth.

The characterization of vector fields given by Proposition (2.1.6) immediately leads to a way of combining two vector fields $X, Y$ to get another. Consider both $X$ and $Y$ as linear maps from $C^{\infty}(M)$ to itself and compose them. Then
$X Y(f g)=X(f(Y g)+g(Y f))=(X f)(Y g)+f(X Y g)+(X g)(Y f)+g(X Y f)$ $Y X(f g)=Y(f(X g)+g(X f))=(Y f)(X g)+f(Y X g)+(Y g)(X f)+g(Y X f)$ and subtracting and writing $[X, Y]=X Y-Y X$ we have

$$
[X, Y](f g)=f([X, Y] g)+g([X, Y] f)
$$

which from Proposition (2.1.6) means that $[X, Y]$ is a vector field.

## Definition (2.1.7):

The Lie bracket of two vector fields $X, Y$ is the vector field $[X, Y]$.

## Example (2.1.8):

If $M=R$ then $X=f d / d x, Y=g d / d x$ and so

$$
[X, Y]=\left(f g^{\prime}-g f^{\prime}\right) \frac{d}{d x}
$$

We shall later see that there is a geometrical origin for the Lie bracket.

In the following we will discuss one-parameter groups of diffeomorphisms. Think of wind velocity (assuming it is constant in time) on the surface of the earth as a vector field on the sphere $S^{2}$. There is another interpretation we can make. A particle at position $x \in S^{2}$ moves after time t seconds to a $\operatorname{position} \varphi_{t}(x) \in S^{2}$. After a further s seconds it is at

$$
\varphi_{t+s}(x)=\varphi_{s}\left(\varphi_{t}(x)\right):
$$

What we get this way is a homomorphism of groups: from the additive group $\boldsymbol{R}$ to the group of diffeomorphisms of $S^{2}$ under the operation of composition. The technical definition is the following:

## Definition (2.1.9):

A one-parameter group of diffeomorphisms of a manifold $M$ is a smooth map

$$
\varphi: M \times R \rightarrow M
$$

such that (writing $\left.\varphi_{t}(x)=\varphi(x, t)\right)$
i. $\varphi_{t}: M \rightarrow M$ is a diffeomorphism
ii. $\varphi_{0}=$ id
iii. $\varphi_{s+t}=\varphi_{s} \circ \varphi_{t}$.

We shall show that vector fields generate one-parameter groups of difffeomorphisms, but only under certain hypotheses. If instead of the whole surface of the earth our manifold is just the interior of the UK and the wind is blowing East-West, clearly after however short a time, some particles will be blown offshore, so we cannot hope for $\varphi_{t}(x)$ that works for all $x$ and $t$. The fact that the earth is compact is one reason why it works there, and this is one of the results below. The idea, nevertheless, works locally and is a useful way of
understanding vector fields as "infinitesimal diffieomorphisms" rather than as abstract derivations of functions.

To make the link with vector fields, suppose $\varphi_{t}$ is a one-parameter group of diffieomorphisms and fa smooth function. Then

$$
f\left(\varphi_{t}(a)\right)
$$

is a smooth function of $t$ and we write

$$
\left.\frac{\partial}{\partial t} f\left(\varphi_{t}(a)\right)\right|_{t=0}=X_{a}(f)
$$

It is straightforward to see that, since $\varphi_{0}(a)=a$ the Leibnitz rule holds and this is a tangent vector at $a$, and so as $a=x$ varies we have a vector field. In local coordinates we have

$$
\varphi_{t}\left(x_{1}, \ldots, x_{n}\right)=\left(y_{1}(x, t), \ldots, y_{n}(x, t)\right)
$$

and

$$
\begin{gathered}
\frac{\partial}{\partial t} f\left(y_{1}, \ldots, y_{n}\right)=\left.\sum_{i} \frac{\partial f}{\partial y_{i}}(y) \frac{\partial y_{i}}{\partial t}(x)\right|_{t=0} \\
=\sum_{i} c_{i}(x) \frac{\partial f}{\partial x_{i}}(x)
\end{gathered}
$$

which yields the vector field

$$
X=\sum_{i} c_{i}(x) \frac{\partial}{\partial x_{i}}
$$

We now want to reverse this: go from the vector field to the diffeomorphism. The first point is to track that "trajectory" of a single particle.

## Definition (2.1.10):

An integral curve of a vector field $X$ is a smooth map $\varphi:(\alpha, \beta) \subset R \rightarrow M$ such that
$D \varphi_{t}\left(\frac{d}{d t}\right)=X_{\varphi(t)}$.

## Example (2.1.11):

Suppose $M=R^{2}$ with coordinates $(x, y)$ and $X=\partial / \partial x$. The derivative $D_{\varphi}$ of the smooth function $\varphi(t)=(x(t), y(t))$ is

$$
D \varphi\left(\frac{d}{d t}\right)=\frac{d x}{d t} \frac{\partial}{\partial x}+\frac{d y}{d t} \frac{\partial}{\partial y}
$$

so the equation for an integral curve of $X$ is

$$
\begin{aligned}
& \frac{d x}{d t}=1 \\
& \frac{d y}{d t}=0
\end{aligned}
$$

which gives

$$
\varphi(t)=\left(t+a_{1}, a_{2}\right)
$$

In our wind analogy, the particle at $\left(a_{1}, a_{2}\right)$ is transported to $\left(t+a_{1}, a_{2}\right)$.
In general we have:

## Theorem (2.1.12):

Given a vector field $X$ on a manifold $M$ and $a \in M$ there exists a maximal integral curve of $X$ through $a$.

By "maximal" we mean that the interval $(\alpha, \beta)$ is maximal - as we saw above it may not be the whole of the real numbers.

## Proof:

First consider a coordinate chart $\left(U_{\gamma}, \psi_{\gamma}\right)$ around $a$ then if

$$
X=\sum_{i} c_{i}(x) \frac{\partial}{\partial x_{i}}
$$

the equation

$$
D \varphi_{t}\left(\frac{d}{d t}\right)=X_{\varphi(t)}
$$

can be written as the system of ordinary diferential equations

$$
\frac{d x_{i}}{d t}=c i\left(x_{1}, \ldots, x_{n}\right)
$$

The existence and uniqueness theorem for ODE's asserts that there is some interval on which there is a unique solution with initial condition

$$
\left(x_{1}(0), \ldots, x_{n}(0)=\psi_{\gamma}(a)\right.
$$

Suppose $\varphi:(\alpha, \beta) \rightarrow M$ is any integral curve with $\varphi(0)=a$. For each $x \in(\alpha, \beta)$ the subset $\varphi([0, x]) \subset M$ is compact, so it can be covered by a finite number of coordinate charts, in each of which we can apply the existence and uniqueness theorem to intervals $\left[0, \alpha_{1}\right],\left[\alpha_{1}, \alpha_{2}\right], \ldots,\left[\alpha_{n}, x\right]$. Uniqueness implies that these local solutions agree with $\varphi$ on any subinterval containing 0 .

We then take the maximal open interval on which we can define '.

To find the one-parameter group of diffeomorphisms we now let $a \in M$ vary. In the example above, the integral curve through $\left(a_{1}, a_{2}\right)$ was $t \mapsto\left(t+a_{1}, a_{2}\right)$ and this defines the group of diffeomorphisms

$$
\varphi_{t}\left(x_{1}, x_{2}\right)=\left(t+x_{1}, x_{2}\right)
$$

## Theorem (2.1.13):

Let $X$ be a vector field on a manifold $M$ and for $(t, x) \in R \times M$, let $\varphi(t, x)=\varphi_{t}(x)$ be the maximal integral curve of $X$ through $x$. Then
i. The map $(t, x) \mapsto \varphi_{t}(x)$ is smooth.
ii. $\varphi_{t} \circ \varphi_{s}=\varphi_{t+s}$ wherever the maps are defined.
iii. If M is compact, then $\varphi_{t}(x)$ is defined on $R \times M$ and gives a one-parameter group of diffeomorphisms.

## Proof:

The previous theorem tells us that for each $a \in M$ we have an open interval $(\alpha(a), \beta(a))$ on which the maximal integral curve is defined. The local existence theorem also gives us that there is a solution for initial conditions in a neighbourhood of $a$ so the set

$$
\{(t, x) \in R \times M: t \in(\alpha(x), \beta(x))\}
$$

is open. This is the set on which $\varphi_{t}(x)$ is maximally defined.

The theorem on smooth dependence on initial conditions tells us that $(t, x) \mapsto$ $\varphi_{t}(x)$ is smooth.

Consider $\varphi_{t} \circ \varphi_{s}(x)$. If we fix $s$ and vary $t$, then this is the unique integral curve of $X$ through $\varphi_{s}(x)$. But $\varphi_{t+s}(x)$ is an integral curve which at $\mathrm{t}=0$ passes through
$\varphi_{s}(x)$. By uniqueness they must agree so that $\varphi_{t} \circ \varphi_{s}=\varphi_{t+s}$. (Note that $\varphi_{t} \circ \varphi_{-t}=i d$ shows that we have a diffeomorphism wherever it is defined).

Now consider the case where $M$ is compact. For each $x \in M$, we have an open interval $(\alpha(x), \beta(x))$ containing 0 and an open set $U_{x} \subseteq M$ on which $\varphi_{t}(x)$ is defined. Cover M by $\left\{U_{x}\right\}_{x \in M}$ and take a finite subcovering $U_{x_{1}}, \ldots, U_{x_{N}}$, and set

$$
I=\bigcap_{1}^{N}\left(\alpha\left(x_{i}\right), \beta\left(x_{i}\right)\right)
$$

which is an open interval containing 0 . By construction, for $t \in I$ we get

$$
\varphi_{t}: I \times M \rightarrow M
$$

which defines an integral curve (though not necessarily maximal) through each point $x \in M$ and with $\varphi_{0}(x)=x$. We need to extend to all real values of $t$.

If $s, t \in R$, choose $n$ such that $(|s|+|t|) / n \in I$ and define (where multiplication is composition)

$$
\varphi_{t}=\left(\varphi_{t / n}\right)^{n}, \quad \varphi_{s}=\left(\varphi_{s / n}\right)^{n}
$$

Now because $t / n, s / n$ and $(s+t) / n$ lie in $I$ we have

$$
\varphi_{t / n} \varphi_{s / n}=\varphi_{(s+t) / n}=\varphi_{s / n} \varphi_{t / n}
$$

and so because $\varphi_{t / n}$ and $\varphi_{s / n}$ commute, we also have

$$
\begin{aligned}
\varphi_{t} \varphi_{s} & =\left(\varphi_{t / n}\right)^{n}\left(\varphi_{s / n}\right)^{n} \\
= & \left(\varphi_{(s+t) / n}\right)^{n} \\
& =\varphi_{s+t}
\end{aligned}
$$

which completes the proof.

Now we will discuss the Lie bracket. All the objects we shall consider will have the property that they can be transformed naturally by a diffeomorphism, and the link between vector fields and diffeomorphisms we have just observed provides an "infinitesimal" version of this.

Given a diffeomorphism $F: M \rightarrow M$ and a smooth function $f$ we get the transformed function $f \circ F$. When $F=\varphi_{t}$, generated according to the theorems above by a vector field $X$, we then saw that

$$
\left.\frac{\partial}{\partial t} f\left(\varphi_{t}\right)\right|_{t=0}=X(f)
$$

So: the natural action of diffeomorphisms on functions specializes through oneparameter groups to the derivation of a function by a vector field.

Now suppose $Y$ is a vector field, considered as a map $Y: M \rightarrow T M$. With a diffeomorphism $F: M \rightarrow M$, its derivative $D F_{x}: T_{x} \rightarrow T_{F(x)}$ gives

$$
D F_{x}\left(Y_{x}\right) \in T_{F(x)}
$$

This defines a new vector field $\tilde{Y}$ by

$$
\begin{equation*}
Y F_{(x)}=D F_{x}\left(\tilde{Y}_{x}\right) \tag{2.1}
\end{equation*}
$$

Thus for a function $f$,

$$
\begin{equation*}
(\tilde{Y})(f \circ F)=(Y f) \circ F \tag{2.2}
\end{equation*}
$$

Now if $F=\varphi_{t}$ for a one-parameter group, we have $\tilde{Y}_{\mathrm{t}}$ and we can differentiate to get

$$
\dot{Y}=\left.\frac{\partial}{\partial t} \tilde{Y}_{\mathrm{t}}\right|_{t=0}
$$

From (2.2) this gives

$$
\dot{Y} f+Y(X f)=X Y f
$$

so that $\dot{Y}=X Y-Y X$ is the natural derivative defined above. Thus the natural action of diffeomorphisms on vector fields specializes through one-parameter groups to the Lie bracket [X, Y].

## Section (2.2): Tensor Products and Exterior Algebra

We begin this section by studding Tensor products, we have so far encountered vector fields and the derivatives of smooth functions as analytical objects on manifolds. These are examples of a general class of objects called tensors which we shall encounter in more generality. The starting point is pure linear algebra.

Let $V, W$ be two finite-dimensional vector spaces over $R$. We are going to define a new vector space $V \otimes W$ with two properties:
i. if $v \in V$ and $w \in W$ then there is a product $v \otimes w \in V \otimes W$
ii. the product is bilinear:

$$
\begin{aligned}
& \left(\lambda v_{1}+\mu v_{2}\right) \otimes w=\lambda v_{1} \otimes w+\mu v_{2} \otimes w \\
& v \otimes\left(\lambda w_{1}+\mu w_{2}\right)=\lambda v \otimes w_{1}+\mu v \otimes w_{2}
\end{aligned}
$$

In fact, it is the properties of the vector space $V \otimes W$ which are more important than what it is (and after all what is a real number? Do we always think of it as an equivalence class of Cauchy sequences of rationals?).

## Proposition (2.2.1):

The tensor product $V \otimes W$ has the universal property that if $B: V \times W \rightarrow U$ is a bilinear map to a vector space $U$ then there is a unique linear map

$$
\beta: V \otimes W \rightarrow U
$$

such that $B(v, w)=\beta(v \otimes w)$.
There are various ways to define $V \otimes W$. In the finite-dimensional case we can say that $V \otimes W$ is the dual space of the space of bilinear forms on $V \times W$ : i.e. maps $B: V \times W \rightarrow R$ such that

$$
\begin{aligned}
& B\left(\lambda v_{1}+\mu v_{2}, w\right)=\lambda B\left(v_{1}, w\right)+\mu B\left(v_{2}, w\right) \\
& B\left(v, \lambda w_{1}+\mu w_{2}\right)=\lambda B\left(v, w_{1}\right)+\mu B\left(v, w_{2}\right)
\end{aligned}
$$

Given $v, w \in V, W$ we then define $v \otimes w \in V \otimes W$ as the map

$$
(v \otimes w)(B)=B(v, w)
$$

This satisfies the universal property because given $B: V \times W \rightarrow U$ and $\xi \in U^{*}, \xi \circ B$ is a bilinear form on $V \times W$ and defines a linear map from $U^{*}$ to the space of bilinear forms. The dual map is the required homomorphism $\beta$ from $V \otimes W$ to $\left(U^{*}\right)^{*}=U$.

A bilinear form B is uniquely determined by its values $B\left(v_{i}, w_{j}\right)$ on basis vectors $v_{1}, \ldots, v_{m}$ for V and $w_{1}, \ldots, w_{n}$ for W which means the dimension of the vector space of bilinear forms is $m n$, as is its dual space $V \otimes W$. In fact, we can easily see that the $m n$ vectors

$$
v_{i} \otimes w_{j}
$$

form a basis for $V \otimes W$. It is important to remember though that a typical element of $V \otimes W$ can only be written as a sum

$$
\sum_{i, j} a_{i j} v_{i} \otimes w_{j}
$$

and not as a pure product $V \otimes W$.

Taking $W=V$ we can form multiple tensor products

$$
V \otimes V, V \otimes V \otimes V=\otimes^{3} V, \ldots
$$

We can think of $\otimes^{p} V$ as the dual space of the space of $p$-fold multilinear forms on $V$.

Mixing degrees we can even form the tensor algebra:

$$
T(V)=\bigotimes_{k=0}^{\infty}\left(\otimes^{k} V\right)
$$

An element of $T(V)$ is a finite sum

$$
\lambda 1+v_{0}+\sum v_{i} \otimes v_{j}+\cdots+\sum v_{i_{1}} \otimes v_{i_{2}} \ldots \otimes v_{i_{p}}
$$

of products of vectors $v_{i} \in V$. The obvious multiplication process is based on extending by linearity the product

$$
\left(v_{1} \otimes \ldots \otimes v_{p}\right)\left(u_{1} \otimes \ldots \otimes u_{q}\right)=v_{1} \otimes \ldots \otimes v_{p} \otimes u_{1} \otimes \ldots \otimes u_{q}
$$

It is associative, but noncommutative.
For the most part we shall be interested in only a quotient of this algebra, called the exterior algebra.

Now we will discuss the exterior algebra. Let $T(V)$ be the tensor algebra of a real vector space $V$ and let $I(V)$ be the ideal generated by elements of the form

$$
v \otimes v
$$

where $v \in V$. So $I(V)$ consists of all sums of multiples by $T(V)$ on the left and right of these generators.

## Definition (2.2.2):

The exterior algebra of $V$ is the quotient

$$
\wedge^{*} V=T(V) / I(V)
$$

If $\pi: T(V) \rightarrow \wedge^{*} V$ is the quotient projection then we set

$$
\Lambda^{p} V=\pi\left(\otimes^{p} V\right)
$$

and call this the p-fold exterior power of $V$. We can think of this as the dual space of the space of multilinear forms $M\left(v_{1}, \ldots, v_{p}\right)$ on V which vanish if any two arguments coincide the so called alternating multilinear forms. If $a \in \bigotimes^{p} V, b \in \bigotimes^{q} V$ then $a \otimes b \in \bigotimes^{p+q} V$ and taking the quotient we get a product called the exterior product:

## Definition (2.2.3):

The exterior product of $\alpha=\pi(a) \in \wedge^{p} V$ and $\beta=\pi(b) \in \wedge^{q} V$ is

$$
\alpha \wedge \beta=\pi(a \otimes b)
$$

## Remark (2.2.4):

If $v_{1}, \ldots, v_{p} \in V$ then we define an element of the dual space of the space of alternating multilinear forms by

$$
v_{1} \wedge v_{2} \wedge \ldots \wedge v_{p}(M)=M\left(v_{1}, \ldots, v_{p}\right):
$$

The key properties of the exterior algebra follow:

## Proposition (2.2.5):

If $\alpha \in \wedge^{p} V, \beta \in \wedge^{q} V$ then

$$
\alpha \wedge \beta=(-1)^{p q} \beta \wedge \alpha
$$

## Proof:

Because for $v \in V, v \otimes v \in I(V)$, it follows that $v \wedge v=0$ and hence

$$
0=\left(v_{1}+v_{2}\right) \wedge\left(v_{1}+v_{2}\right)=0+v_{1} \wedge v_{2}+v_{2} \wedge v_{1}+0
$$

So interchanging any two entries from V in an expression like

$$
v_{1} \wedge \ldots \wedge v_{k}
$$

changes the sign.
Write $\alpha$ as a linear combination of terms $v_{1} \wedge \ldots \wedge v_{p}$ and $\beta$ as a linear combination of $w_{1} \wedge \ldots \wedge w_{q}$ and then, applying this rule to bring $w_{1}$ to the front we see that

$$
\left(v_{1} \wedge \ldots \wedge v_{p}\right) \wedge\left(w_{1} \wedge \ldots \wedge w_{q}\right)=(-1)^{p} w_{1} \wedge v_{1} \wedge \ldots v_{p} \wedge w_{2} \wedge \ldots \wedge w_{q}
$$

For each of the $\mathrm{q} \mathrm{w}_{\mathrm{i}}$ 's we get another factor $(-1)^{p}$ so that in the end

$$
\left(w_{1} \wedge \ldots \wedge w_{q}\right)\left(v_{1} \wedge \ldots \wedge v_{p}\right)=(-1)^{p q}\left(v_{1} \wedge \ldots \wedge v_{p}\right)\left(w_{1} \wedge \ldots \wedge w_{q}\right)
$$

## Proposition (2.2.6):

If $\operatorname{dim} V=n$ then $\operatorname{dim} \wedge^{n} V=1$.

## Proof:

Let $w_{1} \ldots w_{n}$ be n vectors in V and relative to some basis let $M$ be the square matrix whose columns are $w_{1} \ldots w_{n}$. then

$$
B\left(w_{1} \ldots w_{n}\right)=\operatorname{det} M
$$

is a non-zero n -fold multilinear form on $V$. Moreover, if any two of the $w_{i}$ coincide, the determinant is zero, so this is a non-zero alternating n-linear form - an element in the dual space of $\Lambda^{n} V$.

On the other hand, choose a basis $v_{1} \ldots v_{n}$ for V , then anything in $\otimes^{n} V$ is a linear combination of terms like $v_{i_{1}} \otimes \ldots \otimes v_{i_{n}}$ and so anything in $\Lambda^{n} V$ is, after using Proposition (2.2.5), a linear combination of $v_{1} \wedge \ldots \wedge v_{n}$.

Thus $\Lambda^{n} V$ is non-zero and at most one-dimensional hence is one-dimensional.

## Proposition (2.2.7)

let $v_{1}, \ldots, v_{n}$ be a basis for V , then the $\binom{n}{p}$ elements $v_{i_{1}} \wedge v_{i_{2}} \wedge \ldots \wedge v_{i_{p}}$ for $i_{1}<i_{2}<\cdots<i_{p}$ form a basis for $\wedge^{n} V$.

## Proof:

By reordering and changing the sign we can get any exterior product of the $v_{i}{ }^{\prime} s$ so these elements clearly span $\wedge^{n} V$. Suppose then that

$$
\sum a_{i_{1} \ldots i_{p}} v_{i_{1}} \wedge v_{i_{2}} \wedge \ldots \wedge v_{i_{p}}=0
$$

Because $i_{1}<i_{2}<\cdots<i_{p}$, each term is uniquely indexed by the $\operatorname{subset}\left\{i_{1}, i_{2}, \ldots, i_{p}\right\}=I \subseteq\{1,2, \ldots, n\}$, and we can write

$$
\begin{equation*}
\sum_{I} a_{I} v_{I}=0 \tag{2.3}
\end{equation*}
$$

If $I$ and $J$ have a number in common, then $v_{I} \wedge v_{J}=0$, so if $J$ has $n-p$ elements, $v_{I} \wedge v_{J}=0$ unless $J$ is the complementary subset $I^{\prime}$ in which case the product is a multiple of $v_{1} \wedge v_{2} \ldots \wedge v_{n}$ and by Proposition (2.2.6) this is non-
zero. Thus, multiplying (2.3) by each term $v_{I^{\prime}}$ we deduce that each coefficient $a_{I}=0$ and so we have linear independence.

## Proposition (2.2.8):

The vector $v$ is linearly dependent on the linearly independent vectors $v_{1}, \ldots, v_{p}$ if and only if $v_{1} \wedge v_{2} \wedge \ldots \wedge v_{p} \wedge v=0$.

## Proof:

If $v$ is linearly dependent on $v_{1}, \ldots, v_{p}$ then $v=\sum a_{i} v_{i}$ and expanding

$$
v_{1} \wedge v_{2} \wedge \ldots \wedge v_{p} \wedge v=v_{1} \wedge v_{2} \wedge \ldots \wedge v_{p} \wedge\left(\sum_{1}^{p} a_{i} v_{i}\right)
$$

gives terms with repeated $v_{i}$, which therefore vanish. If not, then $v_{1}, v_{2}, \ldots, v_{p}, v$ can be extended to a basis and Proposition (2.2.7) tells us that the product is nonzero.

## Proposition (2.2.9):

If $A: V \rightarrow W$ is a linear transformation, then there is an induced linear transformation

$$
\wedge^{p} A: \wedge^{p} V \rightarrow \wedge^{p} W
$$

such that

$$
\wedge^{p} A\left(v_{1} \wedge \ldots \wedge v_{p}\right)=A v_{1} \wedge A v_{2} \wedge \ldots \wedge A v_{p}
$$

## Proof:

From Proposition (2.2.7) the formula

$$
\wedge^{p} A\left(v_{1} \wedge \ldots \wedge v_{p}\right)=A v_{1} \wedge A v_{2} \wedge \ldots \wedge A v_{p}
$$

actually defines what $\Lambda^{p} A$ is on basis vectors but doesn't prove it is independent of the choice of basis. But the universal property of tensor products gives us

$$
\otimes^{p} A: \otimes^{p} V \rightarrow \otimes^{p} W
$$

and $\otimes^{p} A$ maps the ideal $I(V)$ to $I(W)$ so defines $\wedge^{p} A$ invariantly.

## Proposition (2.2.10):

If $\operatorname{dim} V=n$, then the linear transformation $\wedge^{p} A: \wedge^{n} V \rightarrow \Lambda^{n} V$ is given by $\operatorname{det} A$.

## Proof:

From Proposition (2.2.7), $\wedge^{n} V$ is one-dimensional and so $\wedge^{n} A$ is multiplication by a real number $\lambda(A)$. So with a basis $v_{1}, \ldots, v_{n}$,

$$
\wedge^{n} A\left(v_{1} \wedge \ldots \wedge v_{n}\right)=A v_{1} \wedge A v_{2} \wedge \ldots A v_{n}=\lambda(A) v_{1} \wedge \ldots \wedge v_{n}
$$

But

$$
A v_{i}=\sum_{j} A_{j i} v_{j}
$$

and so

$$
\begin{gathered}
A v_{1} \wedge A v_{2} \wedge \ldots \wedge A v_{n}=\sum A_{j_{1}, 1} v_{j_{1}} \wedge A_{j_{2}, 2} v_{j_{2}} \wedge \ldots \wedge A_{j_{n}, n} v_{j_{n}} \\
=\sum_{\sigma \in S_{n}} A_{\sigma 1,1} v_{\sigma 1} \wedge A_{\sigma 2,2} v_{\sigma 2} \wedge \ldots \wedge A_{\sigma n, n} v_{\sigma n}
\end{gathered}
$$

where the sum runs over all permutations $\sigma$. But if $\sigma$ is a transposition then the term
$v_{\sigma 1} \wedge v_{\sigma 2} \ldots \wedge v_{\sigma n}$ changes sign, so

$$
A v_{1} \wedge A v_{2} \wedge \ldots \wedge A v_{n}=\sum_{\sigma \in S_{n}} \operatorname{sgn} \sigma A_{\sigma 1,1} A_{\sigma 2,2} \ldots A_{\sigma n, n} v_{1} \wedge \ldots \wedge v_{n}
$$

which is the definition of $(\operatorname{det} A) v_{1} \wedge \ldots \wedge v_{n}$.

## Chapter (3)

## Differential Forms ,De Rham Cohomology and S'tokes Theorem

## Section (3.1): Differential Forms and De Rham Cohomology

We begin this section by studding the bundle of p-forms. Now let $M$ be an $n$ dimensional manifold and $\mathrm{T}_{\mathrm{x}}^{*}$ the cotangent space at $x$. We form the $p$-fold exterior power

$$
\wedge^{p} T_{x}^{*}
$$

and, just as we did for the tangent bundle and cotangent bundle, we shall make

$$
\wedge^{p} T^{*} M=\cup_{x \in M} \wedge^{p} T_{x}^{*}
$$

into a vector bundle and hence a manifold.

If $x_{1}, \ldots, x_{n}$ are coordinates for a chart $\left(U, \varphi_{U}\right)$ then for $x \in U$, the elements

$$
d_{x_{i_{1}}} \wedge d_{x_{i_{2}}} \wedge \ldots \wedge d_{x_{i_{p}}}
$$

for $i_{1}<i_{2}<\ldots<i_{p}$ form a basis for $\wedge^{p} T_{x}^{*}$. The $\binom{n}{p}$ coefficients of $\alpha \in \wedge^{p} T_{x}^{*}$ then give a coordinate chart $\Psi_{U}$ mapping to the open set

$$
\varphi_{U}(U) \times \wedge^{p} R^{n} \subseteq R^{n} \times R^{\binom{n}{p}} .
$$

When $p=1$ this is just the coordinate chart we used for the cotangent bundle:

$$
\Phi_{U}\left(x, \sum y_{i} d x_{i}\right)=\left(x_{1}, \ldots, x_{n}, y_{1}, \ldots, y_{n}\right)
$$

and on two overlapping coordinate charts we there had

$$
\Phi_{\beta} \Phi_{\alpha}^{-1}\left(x_{1}, \ldots, x_{n}, y_{1}, \ldots, y_{n}\right)=\left(\tilde{x}_{1}, \ldots, \tilde{x}_{n}, \sum_{j} \frac{\partial \tilde{x}_{i}}{\partial x_{1}} y_{i}, \ldots, \sum_{j} \frac{\partial \tilde{x}_{i}}{\partial x_{n}} y_{n}\right):
$$

For the $p$-th exterior power we need to replace the Jacobian matrix

$$
J=\frac{\partial \tilde{x}_{i}}{\partial x_{j}}
$$

by its induced linear map

$$
\wedge^{p} J: \Lambda^{p} R^{n} \rightarrow \Lambda^{p} R^{n}
$$

It's a long and complicated expression if we write it down in a basis but it is invertible and each entry is a polynomial in $C^{\infty}$ functions and hence gives a smooth map with smooth inverse. In other words,

$$
\Psi_{\beta} \Psi_{\alpha}^{-1}
$$

satisfies the conditions for a manifold of dimension $n+\binom{n}{p}$.

## Definition (3.1.1):

The bundle of $p$-forms of a manifold $M$ is the differentiable structure on $\wedge^{p} T^{*} M$ defined by the above atlas. There is a natural projection $p: \wedge^{p} T^{*} M \rightarrow M$ and a section is called a differential $p$-form

## Examples (3.1.2):

1. A zero-form is a section of $\Lambda^{p} T^{*}$ which by convention is just a smooth function $f$.
2. A 1 -form is a section of the cotangent bundle $T^{*}$. From our definition of the derivative of a function, it is clear that $d f$ is an example of a 1 -form. We can write in a coordinate system

$$
d f=\sum_{j} \frac{\partial f}{\partial x_{j}} d x_{j}
$$

By using a bump function we can extend a locally-defined p-form like $d x_{1} \wedge d x_{2} \wedge \ldots \wedge d x_{p}$ to the whole of $M$, so sections always exist. In fact, it will be convenient at various points to show that any function, form, or vector field can be written as a sum of these local ones. This involves the concept of partition of unity.

Now we will illustrate the partition of unity.

## Definition (3.1.3):

A partition of unity on $M$ is a collection $\left\{\varphi_{i}\right\}_{i \in I}$ of smooth functions such that
i. $\varphi_{i} \geq 0$
ii. $\left\{\operatorname{supp} \varphi_{i}: i \in I\right\}$ is locally finite.
iii. $\sum_{i} \varphi_{i}=1$

Here locally finite means that for each $x \in M$ there is a neighbourhood U which intersects only finitely many supports $\operatorname{supp} \varphi_{i}$.

## Theorem (3.1.4):

Given any open covering $\left\{V_{\alpha}\right\}$ of a manifold $M$ there exists a partition of unity $\left\{\varphi_{i}\right\}$ on $M$ such that $\operatorname{supp} \varphi_{i} \subset V_{\alpha(i)}$ for some $\alpha(i)$.

We say that such a partition of unity is subordinate to the given covering.

Here let us just note that in the case when $M$ is compact, life is much easier: For each point $x \in\left\{V_{\alpha}\right\}$ we take a coordinate neighbourhood $U_{x} \subset\left\{V_{\alpha}\right\}$ and a bump function which is 1 on a neighbourhood $V_{x}$ of $x$ and whose support lies in $U_{x}$. Compactness says we can extract a finite subcovering of the $\left\{V_{x}\right\}_{x \in X}$ and so we get smooth functions $\Psi_{i} \geq 0$ for $i=1, \ldots, N$ and equal to 1 on $V_{x_{i}}$. In particular the sum is positive, and defining

$$
\varphi_{i}=\frac{\psi_{i}}{\sum_{1}^{N} \Psi_{i}}
$$

gives the partition of unity.
Now, not only can we create global p-forms by taking local ones, multiplying by $\varphi_{i}$ and extending by zero, but conversely if $\alpha$ is any $p$-form, we can write it as

$$
\alpha=\left(\sum_{i} \varphi_{i}\right) \alpha=\sum_{i}\left(\varphi_{i} \alpha\right)
$$

which is a sum of extensions of locally defined ones.
At this point, it may not be clear why we insist on introducing these complicated exterior algebra objects, but there are two motivations. One is that the algebraic theory of determinants is, as we have seen, part of exterior algebra, and multiple integrals involve determinants. We shall later be able to integrate p -forms over p dimensional manifolds.

The other is the appearance of the skew-symmetric cross product in ordinary three dimensional calculus, giving rise to the curl differential operator taking vector fields to vector fields. As we shall see, to do this in a coordinate-free way, and in all dimensions, we have to dispense with vector fields and work with differential forms instead.

In the following we discuss the working with differential forms. We defined a differential form in Definition (3.1.1) as a section of a vector bundle. In a local coordinate system it looks like this:

$$
\begin{equation*}
\alpha=\sum_{i_{1}<i_{2}<\ldots<i_{p}} a_{i_{1} i_{2} \ldots i_{p}}(x) d x_{i_{1}} \wedge d x_{i_{2}} \ldots \wedge d x_{i_{p}} \tag{3.1}
\end{equation*}
$$

where the coeffcients are smooth functions. If $x(y)$ is a different coordinate system, then we write the derivatives

$$
d x_{i_{k}}=\sum_{j} \frac{\partial x_{i_{k}}}{\partial y_{j}} d y_{j}
$$

and substitute in (3.1) to get

$$
\alpha=\sum_{j_{1<j_{2<\ldots<j_{p}}}} \tilde{a}_{j_{1} j_{2} \ldots j_{p}}(y) d y_{j_{1}} \wedge d y_{j_{2}} \ldots \wedge d y_{j_{p}}
$$

## Example (3.1.5):

Let $M=R^{2}$ and consider the 2-form $\omega=d x_{1} \wedge d x_{2}$. Now change to polar coordinates on the open set $\left(x_{1}, x_{2}\right) \neq(0,0)$ :

$$
x_{1}=r \cos \theta, x_{2}=r \sin \theta .
$$

We have

$$
\begin{aligned}
& d x_{1}=\cos \theta d r-r \sin \theta d \theta \\
& d x_{2}=\sin \theta d r+r \cos \theta d \theta
\end{aligned}
$$

so that

$$
\omega=(\cos \theta d r-r \sin \theta d \theta) \wedge(\sin \theta d r+r \cos \theta d \theta)=r d r \wedge d \theta:
$$

We shall often write

$$
\Omega^{p}(M)
$$

as the infinite-dimensional vector space of all p -forms on $M$.
Although we first introduced vector fields as a means of starting to do analysis on manifolds, in many ways differential forms are better behaved. For example, suppose we have a smooth map

$$
F: M \rightarrow N:
$$

The derivative of this gives at each point $x \in M$ a linear map

$$
D F_{x}: T_{x} M \rightarrow T_{F_{(x)}} N
$$

but if we have a section of the tangent bundle $T M$ - a vector field $X$ - then $D F_{x}\left(X_{x}\right)$ doesn't in general define a vector field on $N$ - it doesn't tell us what to choose in $T_{a} N$ if $a \in N$ is not in the image of $F$.

On the other hand suppose $\alpha$ is a section of $\wedge^{p} T^{*} \mathrm{~N}$ - a p-form on $N$. Then the dual map

$$
D F_{x}^{\prime}: T_{F(x)}^{*} N \rightarrow T_{x}^{*} M
$$

defines

$$
\wedge^{p}\left(D F_{x}^{\prime}\right): \wedge^{p} p T_{F(x) N} \rightarrow \Lambda^{p} T_{x}^{*} M
$$

and then

$$
\wedge^{p}\left(D F_{x}^{\prime}\right)\left(\alpha_{F(x)}\right)
$$

is defined for all $x$ and is a section of $\wedge^{p} T^{*} \mathrm{M}$ - a $p$-form on $M$.

## Definition (3.1.6):

The pull-back of a $p$-form $\alpha \in \Omega^{p}(N)$ by a smooth map $F: M \rightarrow N$ is the $p$ form $F^{*} \alpha \in \Omega^{p}(M)$ defined by

$$
\left(F^{*} \alpha\right)_{x}=\wedge^{p}\left(D F_{x}^{\prime}\right)\left(\alpha_{F(x)}\right)
$$

## Examples (3.1.7):

1. The pull-back of a 0 -form $f \in C^{\infty}(N)$ is just the composition $f \circ F$.
2. By the definition of the dual map $D F_{x}^{\prime}$ we have

$$
\begin{gathered}
D F_{x}^{\prime}(\alpha)\left(X_{x}\right)=\alpha_{F(x)}\left(D F_{x}\left(X_{x}\right)\right), \text { so if } \alpha=d f, \\
D F_{x}^{\prime}(d f)\left(X_{x}\right)=d f_{F(x)}\left(D F_{x}\left(X_{x}\right)\right)=X_{x}(f \circ F)
\end{gathered}
$$

by the definition of $D F_{x}$. This means that $F^{*}(d f)=d(f \circ F)$.
3. Let $F: R^{3} \rightarrow R^{2}$ be given by

$$
F\left(x_{1}, x_{2}, x_{3}\right)=\left(x_{1} x_{2}, x_{2}+x_{3}\right)=(x, y)
$$

and take

$$
\alpha=x d x \wedge d y
$$

Then, using the definition of $\wedge^{p}\left(D F_{x}^{\prime}\right)$ and the previous example,

$$
\begin{gathered}
F^{*} \alpha=(x \circ F) d(x \circ F) \wedge d(y \circ F)=x_{1} x_{2} d\left(x_{1} x_{2}\right) \wedge d\left(x_{2}+x_{3}\right)= \\
x_{1} x_{2}\left(x_{1} d x_{2}+x_{2} d x_{1}\right) \wedge d\left(x_{2}+x_{3}\right)=x_{1}^{2} x_{2} d x_{2} \wedge d x_{3}+x_{1} x_{2}^{2} d x_{1} \wedge d x_{2}+ \\
x_{1} x_{2}^{2} d x_{1} \wedge d x_{3}
\end{gathered}
$$

From the algebraic properties of the maps

$$
\Lambda^{p} A: \Lambda^{p} V \rightarrow \Lambda^{p} V
$$

We have the following straightforward properties of the pull-back:
i. $(F \circ G)^{*} \alpha=G^{*}\left(F^{*} \alpha\right)$
ii. $F^{*}(\alpha+\beta)=F^{*} \alpha+F^{*} \beta$
iii. $F^{*}(\alpha \wedge \beta)=F^{*} \alpha \wedge F^{*} \beta$

Now we will discuss the exterior derivative. We now come to the construction of the basic differential operator on forms - the exterior derivative which generalizes the grads, divs and curls of three-dimensional calculus. The key feature it has is that it is defined naturally by the manifold structure without any further assumptions.

## Theorem (3.1.8):

On any manifold $M$ there is a natural linear map

$$
d: \Omega^{p}(M) \rightarrow \Omega^{p+1}(M)
$$

called the exterior derivative such that

1. if $f \in \Omega^{0}(M)$, then $d f \in \Omega^{1}(M)$ is the derivative of $f$
2. $d^{2}=0$
3. $d(\alpha \wedge \beta)=d \alpha \wedge \beta+(-1)^{p} \alpha \wedge d \beta$ if $\alpha \in \Omega^{p}(M)$

## Examples (3.1.9):

Before proving the theorem, let's look at $M=R^{3}$, following the rules of the theorem, to see d in all cases $p=0,1,2$.
$p=0$ : by definition

$$
d f=\frac{\partial f}{\partial x_{1}} d x_{1}+\frac{\partial f}{\partial x_{2}} d x_{2}+\frac{\partial f}{\partial x_{3}} d x_{3}
$$

which we normally would write as $\operatorname{grad} f$.
$p=1$ : take a 1 -form

$$
\alpha=a_{1} d x_{1}+a_{2} d x_{2}+a_{3} d x_{3}
$$

then applying the rules we have

$$
\begin{gathered}
d\left(a_{1} d x_{1}+a_{2} d x_{2}+a_{3} d x_{3}\right)=d a_{1} \wedge d x_{1}+d a_{2} \wedge d x_{2}+d a_{3} \wedge d x_{3} \\
=\left(\frac{\partial a_{1}}{\partial x_{1}} d x_{1}+\frac{\partial a_{1}}{\partial x_{2}} d x_{2}+\frac{\partial a_{1}}{\partial x_{3}} d x_{3}\right) \wedge d x_{1}+\cdots \\
=\left(\frac{\partial a_{1}}{\partial x_{3}}-\frac{\partial a_{3}}{\partial x_{1}}\right) d x_{3} \wedge d x_{1}+\left(\frac{\partial a_{2}}{\partial x_{1}}-\frac{\partial a_{1}}{\partial x_{2}}\right) d x_{1} \wedge d x_{2}+\left(\frac{\partial a_{3}}{\partial x_{2}}-\frac{\partial a_{2}}{\partial x_{3}}\right) d x_{2} \wedge d x_{3} .
\end{gathered}
$$

The coefficients of this define what we would call the curl of the vector field a but $\alpha$ has now become a 1 -form $\alpha$ and not a vector field and $\mathrm{d} \alpha$ is a 2 -form, not a vector field. The geometrical interpretation has changed. Note nevertheless that the invariant statement $d^{2}=0$ is equivalent to curl grad $f=0$.
$p=2$ : now we have a 2 -form

$$
\beta=b_{1} d x_{2} \wedge d x_{3}+b_{2} d x_{3} \wedge d x_{1}+b_{3} d x_{1} \wedge d x_{2}
$$

and

$$
\begin{gathered}
d \beta=\frac{\partial b_{1}}{\partial x_{1}} d x_{1} \wedge d x_{2} \wedge d x_{3}+\frac{\partial b_{2}}{\partial x_{2}} d x_{1} \wedge d x_{2} \wedge d x_{3}+\frac{\partial b_{3}}{\partial x_{3}} d x_{1} \wedge d x_{2} \wedge d x_{3} \\
=\left(\frac{\partial b_{1}}{\partial x_{1}}+\frac{\partial b_{2}}{\partial x_{2}}+\frac{\partial b_{3}}{\partial x_{3}}\right) d x_{1} \wedge d x_{2} \wedge d x_{3}
\end{gathered}
$$

which would be the divergence of $a$ vector field $b$ but in our case is applied to a 2form $\beta$. Again $d^{2}=0$ is equivalent to div curl $b=0$.

Here we see familiar formulas, but acting on unfamiliar objects. The fact that we can pull differential forms around by smooth maps will give us a lot more power, even in three dimensions, than if we always considered these things as vector fields.

Let us return to the Theorem (3.1.8) now and give its proof.

## Proof:

We shall define $d \alpha$ by first breaking up $\alpha$ as a sum of terms with support in a local coordinate system (using a partition of unity), define a local d operator using a coordinate system, and then show that the result is independent of the choice.

So, to begin with, write a p-form locally as

$$
\alpha=\sum_{i_{1}<i_{2}<\ldots<i_{p}} a_{i_{1} i_{2} \ldots i_{p}}(x) d x_{i_{1}} \wedge d x_{i_{2}} \ldots \wedge d x_{i_{p}}
$$

and define

$$
d \alpha=\sum_{i_{1}<i_{2}<\ldots<i_{p}} d a_{i_{1} i_{2} \ldots i_{p}} \wedge d x_{i_{1}} \wedge d x_{i_{2}} \wedge \ldots \wedge d x_{i_{p}}
$$

When $p=0$, this is just the derivative, so the first property of the theorem holds.

For the second part, we expand

$$
d \alpha=\sum_{j, i_{1}<i_{2}<\ldots<i_{p}} \frac{\partial a_{i_{1} i_{2} \ldots i_{p}}}{\partial x_{j}} d x_{j} \wedge d x_{i_{1}} \wedge d x_{i_{2}} \wedge \ldots \wedge d x_{i_{p}}
$$

and then calculate

$$
d^{2} \alpha=\sum_{j, k, i_{1}<i_{2}<\ldots<i_{p}} \frac{\partial^{2} a_{i_{1} i_{2} \ldots i_{p}}}{d x_{j} \partial x_{k}} d x_{k} \wedge d x_{j} \wedge d x_{i_{1}} \wedge d x_{i_{2}} \ldots \wedge d x_{i_{p}}
$$

The term

$$
\frac{\partial^{2} a_{i_{1} i_{2} \ldots i_{p}}}{\partial x_{j} \partial x_{k}}
$$

is symmetric in $j, k$ but it multiplies $d x_{k} \wedge d x_{j}$ in the formula which is skewsymmetric in $j$ and $k$, so the expression vanishes identically and $d^{2} \alpha=0$ as required.

For the third part, we check on decomposable forms

$$
\begin{aligned}
& \alpha=f d x_{i_{1}} \wedge \ldots \wedge d x_{i_{p}}=f d x_{I} \\
& \beta=g d x_{j_{1}} \wedge \ldots \wedge d x_{j_{q}}=g d x_{J}
\end{aligned}
$$

and extend by linearity. So

$$
\begin{aligned}
& d(\alpha \wedge \beta)=d\left(f g d x_{I} \wedge d x_{J}\right)=d(f g) \wedge d x_{I} \wedge d x_{J} \\
& =(f d g+g d f) \wedge d x_{I} \wedge d x_{J}=(-1)^{p} f d x_{I} \wedge d g \wedge d x_{J}+d f \wedge d x_{I} \wedge g d x_{J} \\
& \quad=(-1)^{p} \alpha \wedge d \beta+d \alpha \wedge \beta
\end{aligned}
$$

So, using one coordinate syste 2 m we have defined an operation $d$ which satisfies the three conditions of the theorem. Now represent $\alpha$ in coordinates $y_{1}, \ldots, y_{n}$ :

$$
\alpha=\sum_{i_{1}<i_{2}<\ldots<i_{p}} b_{i_{1} i_{2} \ldots i_{p}} d y_{i_{1}} \wedge d y_{i_{2}} \wedge \ldots \wedge d y_{i_{p}}
$$

and define in the same way

$$
d^{\prime} \alpha=\sum_{i_{1}<i_{2}<\ldots<i_{p}} d b_{i_{1} i_{2} \ldots i_{p}} \wedge d y_{i_{1}} \wedge d y_{i_{2}} \wedge \ldots \wedge d y_{i_{p}}
$$

We shall show that $d=d^{\prime}$ by using the three conditions.

From (1) and (3),

$$
\begin{gathered}
d \alpha=d\left(\sum b_{i_{1} i_{2} \ldots i_{p}} d y_{i_{1}} \wedge d y_{i_{2}} \ldots \wedge d y_{i_{p}}\right) \\
=\sum d b_{i_{1} i_{2} \ldots i_{p}} \wedge d y_{i_{1}} \wedge d y_{i_{2}} \wedge \ldots \wedge d y_{i_{p}}+\sum b_{i_{1} i_{2} \ldots i_{p}} d\left(d y_{i_{1}} \wedge d y_{i_{2}} \wedge \ldots \wedge d y_{i_{p}}\right)
\end{gathered}
$$

and from (3)

$$
\begin{gathered}
d\left(d y_{i_{1}} \wedge d y_{i_{2}} \wedge \ldots \wedge d y_{i_{p}}\right)=d\left(d y_{i_{1}}\right) \wedge d y_{i_{2}} \wedge \ldots \wedge d y_{i_{p}}-d y_{i_{1}} \wedge \\
d\left(d y_{i_{2}} \wedge \ldots \wedge d y_{i_{p}}\right)
\end{gathered}
$$

From (1) and (2) $d^{2} y_{i_{1}}=0$ and continuing similarly with the right hand term, we get zero in all terms.

Thus on each coordinate neighbourhood
$U d \alpha=\sum_{i_{1}<i_{2}<\ldots<i_{p}} d b_{i_{1} i_{2} \ldots i_{p}} \wedge d y_{i_{1}} \wedge d y_{i_{2}} \wedge \ldots \wedge d y_{i_{p}}=d^{\prime} \alpha$ and $d \alpha$ is thus globally well-defined.

One important property of the exterior derivative is the following:

## Proposition (3.1.10):

Let $F: M \rightarrow N$ be a smooth map and $\alpha \in \Omega^{p}(N)$. Then

$$
d\left(F^{*} \alpha\right)=F^{*}(d \alpha)
$$

## Proof:

Recall that the derivative $D F_{x}: T_{x} M \rightarrow T_{F(x)} N$ was defined in (1.2.7) by

$$
D F_{x}\left(X_{x}\right)(f)=X_{x}(f \circ F)
$$

so that the dual map $D F_{x}^{\prime}: T_{F(x)}^{*} N \rightarrow T_{x}^{*} M$ satisfies

$$
D F_{x}^{\prime}(d f)_{F(x)}=d(f \circ F)_{x}
$$

From the definition of pull-back this means that

$$
\begin{equation*}
F^{*}(d f)=d(f \circ F)=d\left(F^{*} f\right) \tag{3.2}
\end{equation*}
$$

Now if

$$
\begin{gathered}
\alpha=\sum_{i_{1}<i_{2}<\ldots<i_{p}} a_{i_{1} i_{2} \ldots i_{p}}(x) d x_{i_{1}} \wedge d x_{i_{2}} \wedge \ldots \wedge d x_{i_{p}}, \\
F^{*} \alpha=\sum_{i_{1}<i_{2}<\ldots<i_{p}} a_{i_{1} i_{2} \ldots i_{p}}(F(x)) F^{*} d x_{i_{1}} \wedge F^{*} d x_{i_{2}} \wedge \ldots \wedge F^{*} d x_{i_{p}}
\end{gathered}
$$

by the multiplicative property of pull-back and then using the properties of $d$ and (3.2)

$$
\begin{gathered}
d\left(F^{*} \alpha\right)=\sum_{i_{1}<i_{2}<\cdots<i_{p}} d\left(a_{i_{1} i_{2} \ldots i_{p}}(F(x))\right) \wedge F^{*} d x_{i_{1}} \wedge F^{*} d x_{i_{2}} \wedge \ldots \wedge F^{*} d x_{i_{p}} \\
=\sum_{i_{1}<i_{2}<\cdots<i_{p}} F^{*} d a_{i_{1} i_{2} \ldots i_{p}} \wedge F^{*} d x_{i_{1}} \wedge F^{*} d x_{i_{2}} \wedge \ldots \wedge F^{*} d x_{i_{p}}=F^{*}(d \alpha)
\end{gathered}
$$

In the following we will study the lie derivative of a differential form. Suppose $\varphi_{t}$ is the one-parameter (locally defined) group of diffeomorphisms defined by a vector field $X$. Then there is a naturally defined Lie derivative

$$
\mathcal{L}_{X} \alpha=\left.\frac{\partial}{\partial t} \varphi_{t}^{*} \alpha\right|_{t=0}
$$

of a $p$-form $\alpha$ by $X$. It is again a p-form. We shall give a useful formula for this involving the exterior derivative.

## Proposition (3.1.11):

Given a vector field $X$ on a manifold $M$, there is a linear map
$i_{X}: \Omega^{p}(M) \rightarrow \Omega^{p-1}(M)$
(called the interior product) such that
i. $i_{X} d f=X(f)$
ii. $i_{X}(\alpha \wedge \beta)=i_{X} \alpha \wedge \beta+(-1)^{p} \alpha \wedge i_{X} \beta$ if $\alpha \in \Omega^{p}(M)$

The proposition tells us exactly how to work out an interior product: if

$$
X=\sum_{i} a_{i} \frac{\partial}{\partial x_{i}}
$$

and $\alpha=d x_{1} \wedge d x_{2} \wedge \ldots \wedge d x_{p}$ is a basic $p-$ form then

$$
\begin{equation*}
i_{X} \alpha=a_{1} d x_{2} \wedge \ldots \wedge d x_{p}-a_{2} d x_{1} \wedge d x_{3} \wedge \ldots \wedge d x_{p}+\ldots \tag{3.3}
\end{equation*}
$$

In particular

$$
i_{X}\left(i_{X} \alpha\right)=a_{1} a_{2} d x_{3} \wedge \ldots \wedge d x_{p}-a_{2} a_{1} d x_{3} \wedge \ldots \wedge d x_{p}+\ldots=0
$$

## Example (3.1.12):

Suppose

$$
\alpha=d x \wedge d y, \quad X=x \frac{\partial}{\partial x}+y \frac{\partial}{\partial y}
$$

then

$$
i_{X} \alpha=x d y-y d x
$$

The interior product is just a linear algebra construction. Above we have seen how to work it out when we write down a form as a sum of basis vectors. We just need to prove that it is well-defined and independent of the way we do that, which motivates the following, more abstract proof:

## Proof:

In Remark (2.2.4) we defined $\wedge^{p} V$ as the dual space of the space of alternating p-multilinear forms on $V$. If M is an alternating $(p-1)$-multilinear form on $V$ and $\boldsymbol{\xi}$ a linear form on $V$ then

$$
\begin{equation*}
(\xi M)\left(v_{1}, \ldots, v_{p}\right)=\xi\left(v_{1}\right) M\left(v_{2}, \ldots, v_{p}\right)-\xi\left(v_{2}\right) M\left(v_{1}, v_{3}, \ldots, v_{p}\right)+\ldots \tag{3.4}
\end{equation*}
$$

is an alternating p-multilinear form. So if $\alpha \in \wedge^{p} V$ we can define $i_{\xi} \alpha \in \wedge^{p-1} V$ by $\left(i_{\xi} \alpha\right)(M)=\alpha(\xi M)$

Taking $V=T^{*}$ and $\xi=X \in V^{*}=\left(T^{*}\right)^{*}=T$ gives the interior product. Equation (3.4) gives us the rule (3.3) for working out interior products.

Here then is the formula for the Lie derivative:

## Proposition (3.1.13):

The Lie derivative $\mathcal{L}_{X} \alpha$ of a $p$-form $\alpha$ is given by

$$
\mathcal{L}_{X} \alpha=d\left(i_{X} \alpha\right)+i_{X} d \alpha
$$

## Proof:

Consider the right hand side

$$
R_{X}(\alpha)=d\left(i_{X} \alpha\right)+i_{X} d \alpha
$$

Now $i_{X}$ reduces the degree $p$ by 1 but d increases it by 1 , so $R_{X}$ maps p-forms to $p$-forms. Also,

$$
d\left(d\left(i_{X} \alpha\right)+i_{X} d \alpha\right)=d i_{X} d \alpha=\left(d i_{X}+i_{X} d\right) d \alpha
$$

because $d^{2}=0$, so $R_{X}$ commutes with d. Finally, because

$$
\begin{aligned}
& i_{X}(\alpha \wedge \beta)=i_{X} \alpha \wedge \beta+(-1)^{p} \alpha \wedge i_{X} \beta \\
& d(\alpha \wedge \beta)=d \alpha \wedge \beta+(-1)^{p} \alpha \wedge d \beta
\end{aligned}
$$

we have

$$
R_{X}(\alpha \wedge \beta)=\left(R_{X} \alpha\right) \wedge \beta+\alpha \wedge R_{X}(\beta)
$$

On the other hand

$$
\varphi_{t}^{*}(d \alpha)=d\left(\varphi_{t}^{*} \alpha\right)
$$

so differentiating at $t=0$, we get

$$
\mathcal{L}_{X} d \alpha=d\left(\mathcal{L}_{X} \alpha\right)
$$

and

$$
\varphi_{t}^{*}(\alpha \wedge \beta)=\varphi_{t}^{*} \alpha \wedge \varphi_{t}^{*} \beta
$$

and differentiating this, we have

$$
\mathcal{L}_{X}(\alpha \wedge \beta)=\mathcal{L}_{X} \alpha \wedge \beta+\alpha \wedge \mathcal{L}_{X} \beta
$$

Thus both $\mathcal{L}_{X}$ and $R_{X}$ preserve degree, commute with d and satisfy the same Leibnitz identity. Hence, if we write a p-form as

$$
\begin{aligned}
\alpha=\sum_{-}\left(i_{-} 1\right. & <i_{-} 2<\cdots \\
& \left.<i_{-} p\right) \text {. }{ }^{\text {in }}\left(a_{-}\left(i_{-} 1 i_{-} 2 \ldots i_{-} p\right)(x) d x_{-}\left(i_{-} 1\right) \wedge d x_{-}\left(i_{-} 2\right) \wedge \ldots\right. \\
& \wedge d x_{-}\left(i_{-} p\right) \rrbracket
\end{aligned}
$$

$\mathcal{L}_{X}$ and $R_{X}$ will agree so long as they agree on functions. But

$$
R_{X} f=i_{X} d f=X(f)=\left.\frac{\partial}{\partial t} f(\varphi t)\right|_{t=0}=\mathcal{L}_{X} f
$$

so they do agree.
Now we will study the de Rham cohomology. In textbooks on vector calculus, one may read not only that curl $\operatorname{grad} f=0$, but also that if a vector field $a$ satisfies curl $a=0$, then it can be written as $a=\operatorname{grad} f$ for some function $f$. Sometimes the
statement is given with the proviso that the open set of $R^{3}$ on which $a$ is defined satisfies the topological condition that it is simply connected (any closed path can be contracted to a point).

In the language of differential forms on a manifold, the analogue of the above statement would say that if a 1 -form $\alpha$ satisfies $d \alpha=0$, and $M$ is simplyconnected, there is a function f such that $d f=\alpha$.

While this is true, the criterion of simply connectedness is far too strong. We want to know when the kernel of

$$
d: \Omega^{1}(M) \rightarrow \Omega^{2}(M)
$$

is equal to the image of

$$
d: \Omega^{0}(M) \rightarrow \Omega^{1}(M):
$$

Since $d^{2} f=0$, the second vector space is contained in the first and what we shall do is simply to study the quotient, which becomes a topological object in its own right, with an algebraic structure which can be used to say many things about the global topology of a manifold.

## Definition (3.1.14)

The p-th de Rham cohomology group of a manifold $M$ is the quotient vector space:

$$
H^{p}(M)=\frac{\operatorname{Kerd} d: \Omega^{p}(M) \rightarrow \Omega^{p+1}(M)}{\operatorname{Imd} d: \Omega^{p-1}(M) \rightarrow \Omega^{p}(M)}
$$

## Remark (3.1.15):

1. Although we call it the cohomology group, it is simply a real vector space. There are analogous structures in algebraic topology where the additive group structure is more interesting.
2. Since there are no forms of degree -1 , the group $H^{0}(M)$ is the space of functions $f$ such that $d f=0$. Now each connected component $M_{i}$ of $M$ is an open set of $M$ and hence a manifold.

The mean value theorem tells us that on any open ball in a coordinate neighbourhood of $M_{i}, d f=0$ implies that $f$ is equal to a constant $c$, and the subset of Mi on which $f=c$ is open and closed and hence equal to $M_{i}$.

Thus if $M$ is connected, the de Rham cohomology group $H^{0}(M)$ is naturally isomorphic to $\boldsymbol{R}$ : the constant value $c$ of the function $f$. In general $H^{0}(\mathrm{M})$ is the vector space of real valued functions on the set of components. Our assumption that $M$ has a countable basis of open sets means that there are at most countably many components. When $M$ is compact, there are only finitely many, since components provide an open covering. In fact, the cohomology groups of a compact manifold are finite-dimensional vector spaces for all $p$, though we shall not prove that here.

It is convenient in discussing the exterior derivative to introduce the following terminology:

## Definition (3.1.16)

A form $\alpha \in \Omega^{p}(M)$ is closed if $d \alpha=0$.

## Definition (3.1.17):

A form $\alpha \in \Omega^{p}(M)$ is exact if $\alpha=d \beta$ for some $\beta \in \Omega^{p-1}(M)$.
The de Rham cohomology group $H^{p}(M)$ is by definition the quotient of the space of closed $p$-forms by the subspace of exact $p$-forms. Under the quotient map, a closed $p$-form $\alpha$ defines a cohomology class $[\alpha] \in H^{p}(M)$, and $\left[\alpha^{\prime}\right]=[\alpha]$ if and only if $\alpha^{\prime}-\alpha=d \beta$ for some $\beta$.

Here are some basic features of the de Rham cohomology groups.

## Proposition (3.1.18):

The de Rham cohomology groups of a manifold M of dimension n have the following properties:
i. $H^{p}(M)=0$ if $p>n$
ii. for $a \in H^{p}(M), b \in H^{q}(M)$ there is a bilinear product $a b \in H^{p+q}(M)$ which satisfies

$$
a b=(-1)^{p q} b a
$$

iii. if $F: M \rightarrow N$ is a smooth map, it defines a natural linear map

$$
F^{*}: H^{p}(N) \rightarrow H^{p}(M)
$$

which commutes with the product.

## Proof:

The first part is clear since $\wedge^{p} T^{*}=0$ for $p>n$.

For the product, this comes directly from the exterior product of forms. If $a=[\alpha], b=[\beta]$ we define

$$
a b=[\alpha \wedge \beta]
$$

but we need to check that this really does define a cohomology class. Firstly, since $\alpha, \beta$ are closed,

$$
d(\alpha \wedge \beta)=d \alpha \wedge \beta+(-1)^{p} \alpha \wedge d \beta=0
$$

so there is a class defined by $\alpha \wedge \beta$. Suppose we now choose a different representative $\alpha^{\prime}=\alpha+d \gamma$ for a. Then

$$
\alpha^{\prime} \wedge \beta=(\alpha+d \gamma) \wedge \beta=\alpha \wedge \beta+d(\gamma \wedge \beta)
$$

using $d \beta=0$, so $\mathrm{d}(\gamma \wedge \beta)=d \gamma \wedge \beta$. Thus $\alpha^{\prime} \wedge \beta$ and $\alpha \wedge \beta$ differ by an exact form and define the same cohomology class. Changing $\beta$ gives the same result.

The last part is just the pull-back operation on forms. Since

$$
d F^{*} \alpha=F^{*} d \alpha
$$

$F^{*}$ defines a map of cohomology groups. And since

$$
F^{*}(\alpha \wedge \beta)=F^{*} \alpha \wedge F^{*} \beta
$$

it respects the product.

Perhaps the most important property of the de Rham cohomology, certainly the one that links it to algebraic topology, is the deformation invariance of the induced maps $F$. We show that if $F_{t}$ is a smooth family of smooth maps, then the effect on cohomology is independent of $t$. As a matter of terminology (because we have only defined smooth maps of manifolds) we shall say that a map

$$
F: M \times[a, b] \rightarrow N
$$

is smooth if it is the restriction of a smooth map on the product with some slightly bigger open interval $M \times(a-\epsilon, b+\epsilon)$.

## Theorem (3.1.19):

Let $F: M \times[0,1] \rightarrow N$ be a smooth map. Set $F_{t}(x)=F(x, t)$ and consider the induced map on de Rham cohomology $F_{t}^{*}: H^{p}(N) \rightarrow H^{p}(M)$ : Then

$$
F_{1}^{*}=F_{0}^{*}
$$

## Proof:

Represent $a \in H^{p}(N)$ by a closed $p$-form $\alpha$ and consider the pull-back form $F^{*} \alpha$ on $M \times[0,1]$. We can decompose this uniquely in the form

$$
\begin{equation*}
F^{*} \alpha=\beta+d t \wedge \gamma \tag{3.5}
\end{equation*}
$$

where $\beta$ is a $p$-form on $M$ (also depending on t ) and $\gamma$ is a ( $\mathrm{p}-1$ )-form on $M$, depending on $t$. In a coordinate system it is clear how to do this, but more invariantly, the form $\beta$ is just $F_{t}^{*} \alpha$. To get $\gamma$ in an invariant manner, we can think of

$$
(x, s) \mapsto(x, s+t)
$$

as a local one-parameter group of diffeomorphisms of $M \times(a, b)$ which generates a vector field $X=\partial / \partial t$. Then

$$
\gamma=i_{X} F^{*} \alpha
$$

Now $\alpha$ is closed, so from (3.5),

$$
0=d_{M} \beta+d t \wedge \frac{\partial \beta}{\partial t}-d t \wedge d_{M} \gamma
$$

where $d_{M}$ is the exterior derivative in the variables of $M$. It follows that

$$
\frac{\partial \beta}{\partial t}=d_{M} \gamma
$$

Now integrating with respect to the parameter $t$, and using

$$
\frac{\partial}{\partial t} F_{t}^{*} \alpha=\frac{\partial \beta}{\partial t}
$$

we obtain

$$
F_{1}^{*} \alpha-F_{0}^{*} \alpha=\int_{0}^{1} \frac{\partial}{\partial t} F_{t}^{*} \alpha d t=d \int_{0}^{1} d t
$$

So the closed forms $F_{1}^{*} \alpha$ and $F_{0}^{*} \alpha$ differ by an exact form and

$$
F_{1}^{*}(a)=F_{0}^{*}(a):
$$

Here is an immediate corollary:

## Proposition (3.1.20):

The de Rham cohomology groups of $M=R^{n}$ are zero for $p>0$.

## Proof

Define $F: R^{n} \times[0,1] \rightarrow R^{n}$ by

$$
F(x, t)=t x
$$

Then $F_{1}(x)=x$ which is the identity map, and so

$$
F_{1}^{*}: H^{p}\left(R^{n}\right) \rightarrow H^{p}\left(R^{n}\right)
$$

is the identity.

But $F_{0}(x)=0$ which is a constant map. In particular the derivative vanishes, so the pull-back of any $p$-form of degree greater than zero is the zero map. So for $p>0$

$$
F_{0}^{*}: H^{p}\left(R^{n}\right) \rightarrow H^{p}\left(R^{n}\right)
$$

vanishes.

From Theorem (3.1.19) $F_{0}^{*}=F_{1}^{*}$ and we deduce that $H^{p}\left(R^{n}\right)$ vanishes for $p>$ 0 . Of course $\boldsymbol{R}^{\boldsymbol{n}}$ is connected so $H^{0}\left(\boldsymbol{R}^{\boldsymbol{n}}\right) \cong \boldsymbol{R}$.

We are in no position yet to calculate many other de Rham cohomology groups, but here is a first non-trivial example. Consider the case of $\boldsymbol{R} / \boldsymbol{Z}$, diffeomorphic to the circle. In the atlas given earlier, we had $\varphi_{1} \varphi_{0}^{-1}(x)=x$ or $\varphi_{1} \varphi_{0}^{-1}(x)=x-1$ so the 1-form $d x=d(x-1)$ is well-defined, and nowhere zero. It is not the derivative of a function, however, since $\boldsymbol{R} / \boldsymbol{Z}$ is compact and any function must have a minimum where $d f=0$. We deduce that

$$
H^{1}(\boldsymbol{R} / \boldsymbol{Z}) \neq 0
$$

On the other hand, suppose that $\alpha=g(x) d x$ is any 1-form (necessarily closed because it is the top degree). Then $g$ is a periodic function: $g(x+1)=g(x)$. To solve $d f=\alpha$ means solving $f^{\prime}(x)=g(x)$ which is easily done on $\boldsymbol{R}$ by:

$$
f(x)=\int_{0}^{x} g(s) d s
$$

But we want $f(x+1)=f(x)$ which will only be true if

$$
\int_{0}^{1} g(x) d x=0
$$

Thus in general

$$
\alpha=g(x) d x=\left(\int_{0}^{1} g(s) d s\right) d x+d f
$$

and any 1-form is of the form $c d x+d f$. Thus $H^{1}(\boldsymbol{R} / \mathbf{Z}) \cong \boldsymbol{R}$.
We can use this in fact to start an inductive calculation of the de Rham cohomology of the $n$-sphere.

## Theorem (3.1.21):

For $n>0, H^{p}\left(S^{n}\right) \cong \boldsymbol{R}$ if $p=0$ or $p=n$ and is zero otherwise.

## Proof:

We have already calculated the case of $n=1$ so suppose that $n>1$.
Clearly the group vanishes when $p>n$, the dimension of $S^{n}$, and for $n>0, S^{n}$ is connected and so $H^{0}\left(S^{n}\right) \cong \boldsymbol{R}$.

Decompose $S^{n}$ into open sets, $V$, the complement of closed balls around the North and South poles respectively. By stereographic projection these are diffeomorphic to open balls in $\boldsymbol{R}^{n}$. If $\alpha$ is a closed $p$-form for $l<p<n$, then by the Poincare lemma $\alpha=d u$ on $U$ and $\alpha=d v$ on $V$ for some $(p-1)$ forms $u, v$. On the intersection $U \cap V$,

$$
d(u-v)=\alpha-\alpha=0
$$

so $(u-v)$ is closed. But

$$
U \cap V \cong S^{n-1} \times R
$$

so

$$
H^{p-1}(U \cap V) \cong H^{p-1}\left(S^{n-1}\right)
$$

and by induction this vanishes, so on $U \cap V, u-v=d w$.
Now look at $U \cap V$ as a product with a finite open interval: $S^{n-1} \times(-2,2)$. We can find a bump function $\varphi(s)$ which is 1 for $s \in(-1,1)$ and has support in $(-2, \quad 2) . \quad$ Take slightly smaller sets $U^{\prime} \subset U, V^{\prime} \subset V$ such that $U^{\prime} \cap V^{\prime}=S^{n-1} \times(-1,1)$. Then $\varphi w$ extends by zero to define a form on $S^{n}$ and we have $u$ on $U^{\prime}$ and $\quad v+d(\varphi w)$ on $V^{\prime}$ with $u=v+d w=v+d(\varphi w)$ on $U^{\prime} \cap V^{\prime}$. Thus we have defined a ( $\mathrm{p}-1$ ) form $\beta$ on $S^{n}$ such that $\beta=u$ on $U^{\prime}$ and $v+d(\varphi w)$ on $V^{\prime}$ and $\alpha=d \beta$ on $U^{\prime}$ and $V^{\prime}$ and so globally $\alpha=d \beta$. Thus the cohomology class of $\alpha$ is zero.

This shows that we have vanishing of $H^{p}\left(S^{n}\right)$ for $1<p<n$.

When $p=1$, in the argument above $u-v$ is a function on $U \cap V$ and since $d(u-v)=0$ it is a constant c if $U \cap V$ is connected, which it is for $n>1$. Then $d(v+c)=\alpha$ and the pair of functions $u$ on $U$ and $v+c$ on $V$ agree on the overlap and define a function f such that $d f=\alpha$.

When $p=n$ the form $u-v$ defines a class in $H^{n-1}(U \cap V) \cong H^{n-1}\left(S^{n-1}\right) \cong \boldsymbol{R}$. So let $\omega$ be an $(n-1)$ form on $S^{n-1}$ whose cohomology class is non-trivial and pull it back to $S^{n-1} \times(-2,2)$ by the projection onto the first factor. Then $H^{n-1}\left(S^{n-1} \times(-2,2)\right)$ is generated by $[\omega]$ and we have

$$
u-v=\lambda \omega+d w
$$

for some $\lambda \in \boldsymbol{R}$. If $\lambda=0$ we repeat the process above, so $H^{n}\left(S^{n}\right)$ is at most onedimensional. Note that $\lambda$ is linear in $\alpha$ and is independent of the choice of $u$ and $v$ - if we change $u$ by a closed form then it is exact since $H^{p-1}(U)=0$ and we can incorporate it into $w$.

All we need now is to find a class in $H^{n}\left(S^{n}\right)$ for which $\lambda \neq 0$. To do this consider

$$
\varphi d t \wedge \omega
$$

extended by zero outside $U \cap V$. Then

$$
\left(\int_{-2}^{t} \varphi(s) d s\right) \omega
$$

vanishes for $t<-2$ and so extends by zero to define a form $u$ on $U$ such that $d u=\alpha$. When $t>2$ this is non-zero but we can change this to

$$
v=\left(\int_{-2}^{t} \varphi(s) d s\right) \omega-\left(\int_{-2}^{2} \varphi(s) d s\right) \omega
$$

which does extend by zero to $V$ and still satisfies $d v=\alpha$. Thus taking the difference, $\lambda$ above is the positive number

$$
\lambda=\int_{-2}^{2} \varphi(s) d s
$$

To get more information on de Rham cohomology we need to study the other aspect of differential forms: integration.

## Section (3.2): Forms Integration and Stokes' Theorem

We will begin this section by studding the orientation. Recall the change of variables formula in a multiple integral:

$$
\int f\left(y_{1}, \ldots, y_{n}\right) d y_{1} d y_{2} \ldots d y_{n}=\int f\left(y_{1}(x), \ldots, y_{n}(x)\right)\left|\operatorname{det} \partial y_{i} / \partial x_{j}\right| d x_{1} d x_{2} \ldots d x_{n}
$$ and compare to the change of coordinates for an $n$-form on an $n$-dimensional manifold:

$$
\begin{gathered}
\theta=f\left(y_{1}, \ldots, y_{n}\right) d y_{1} \wedge d y_{2} \wedge \ldots \wedge d y_{n}=f\left(y_{1}(x), \ldots, y_{n}(x)\right) \sum_{i} \frac{\partial y_{1}}{\partial x_{i}} d x_{i} \wedge \\
\ldots \wedge \sum_{p} \frac{\partial y_{n}}{\partial x_{p}} d x_{p}=f\left(y_{1}(x), \ldots, y_{n}(x)\right)\left(\operatorname{det} \partial y_{i} / \partial x_{j}\right) d x_{1} \wedge d x_{2} \ldots \wedge d x_{n}
\end{gathered}
$$

The only difference is the absolute value, so that if we can sort out a consistent sign, then we should be able to assign a coordinate-independent value to the integral of an $n$-form over an n-dimensional manifold. The sign question is one of orientation.

## Definition (3.2.1):

An n-dimensional manifold is said to be orientable if it has an everywhere nonvanishing $n$-form $\omega$.

## Definition (3.2.2):

Let $M$ be an n-dimensional orientable manifold. An orientation on $M$ is an equivalence class of non-vanishing n-forms $\omega$ where $\omega \sim \omega^{\prime}$ if $\omega^{\prime}=f \omega$ with $f>0$.

Clearly a connected orientable manifold has two orientations: the equivalence classes of $\pm \omega$.

## Example (3.2.3):

1. Let $M \subset \boldsymbol{R}^{n+1}$ be defined by $f(x)=c$, with $d f(a) \neq 0$ if $f(a)=c$. By Theorem (1.1.8), M is a manifold and moreover, if $\partial f / \partial x_{i} \neq 0, x_{1}, \ldots, x_{i-1}, x_{i+1}, x_{n+1}$ are local coordinates. Consider, on such a coordinate patch,

$$
\begin{equation*}
\omega=(-1)^{i} \frac{1}{\partial f / \partial x_{i}} d x_{1} \wedge \ldots \wedge d x_{i-1} \wedge d x_{i+1} \ldots \wedge d x_{n+1} \tag{3.6}
\end{equation*}
$$

This is non-vanishing.

Now $M$ is defined by $f(x)=c$ so that on $M$

$$
\sum_{j} \frac{\partial f}{\partial x_{j}} d x_{j}=0
$$

and if $\partial f / d x_{j} \neq 0$

$$
d x_{j}=-\frac{1}{\partial f / \partial x_{j}}\left(\partial f / \partial x_{i} d x_{i}+\cdots\right)
$$

Substituting in (3.6) we get

$$
\omega=(-1)^{j} \frac{1}{\partial f / \partial x_{j}} d x_{1} \wedge \ldots \wedge d x_{j-1} \wedge d x_{j+1} \ldots \wedge d x_{n+1}
$$

The formula (3.6) therefore defines for all coordinate charts a non-vanishing n form, so $M$ is orientable.

The obvious example is the sphere $S^{n}$ with

$$
\omega=(-1)^{i} \frac{1}{x_{i}} d x_{1} \wedge \ldots \wedge d x_{i-1} \wedge d x_{i+1} \ldots \wedge d x_{n+1}
$$

2. Consider real projective space $R P^{n}$ and the smooth map

$$
p: S^{n} \rightarrow R P^{n}
$$

which maps a unit vector in $R^{n+1}$ to the one-dimensional subspace it spans. Concretely, if $x_{1} \neq 0$, we use $x=\left(x_{2}, \ldots, x_{n+1}\right)$ as coordinates on $S^{n}$ and the usual coordinates $\left(x_{2} / x_{1}, \ldots, x_{n+1} / x_{1}\right)$ on $R P^{n}$, then

$$
\begin{equation*}
p(x)=\frac{1}{\sqrt{1+\|x\|^{2}}} x \tag{3.7}
\end{equation*}
$$

This is smooth with smooth inverse

$$
q(y)=\frac{1}{\sqrt{1+\|y\|^{2}}} y
$$

so we can use $\left(x_{2}, \ldots, x_{n+1}\right)$ as local coordinates on $R P^{n}$.

Let $\sigma: S^{n} \rightarrow S^{n}$ be the di_eomorphism $\sigma(x)=-x$. Then

$$
\sigma^{*} \omega=(-1)^{i} \frac{1}{-x_{i}} d\left(-x_{i}\right) \wedge \ldots \wedge d\left(-x_{i-1}\right) \wedge d\left(-x_{i+1}\right) \ldots \wedge d\left(-x_{n+1}\right)=(-1)^{n-1} \omega
$$

Suppose $R P^{n}$ is orientable, then it has a non-vanishing n-form $\theta$. Since the map (3.7) has a local smooth inverse, the derivative of $p$ is invertible, so that $p^{*} \theta$ is a non-vanishing n-form on $S^{n}$ and so

$$
p^{*} \theta=f \omega
$$

for some non-vanishing smooth function $f$. But $p \circ \sigma=p$ so that

$$
f \omega=p^{*} \theta=\sigma^{*} p^{*} \theta=(f \circ \sigma)(-1)^{n-1} \omega
$$

Thus, if $n$ is even,

$$
f \circ \sigma=-f
$$

and if $f(a)>0, f(-a)<0$. But $\boldsymbol{R} P^{n}=p\left(S^{n}\right)$ and $S^{n}$ is connected so $\boldsymbol{R} P^{n}$ is connected. This means that f must vanish somewhere, which is a contradiction. Hence $R P^{2 m}$ is not orientable.

There is a more sophisticated way of seeing the non-vanishing form on $S^{n}$ which gives many more examples. First note that a non-vanishing $n$-form on an $n$ dimensional manifold is a non-vanishing section of the rank 1 vector bundle $\Lambda^{n} T^{*} M$. The top exterior power has a special property: suppose $U \subset V$ is an mdimensional vector subspace of an $n$-dimensional space V , then $V / U$ has dimension $n-m$. There is then a natural isomorphism

$$
\begin{equation*}
\Lambda^{m} U \otimes \wedge^{n-m}(V / U) \cong \Lambda^{n} V \tag{3.8}
\end{equation*}
$$

To see this let $u_{1}, \ldots, u_{m}$ be a basis of $U$ and $v_{1}, \ldots, v_{n-m}$ vectors in $V / U$. By definition there exist vectors $\tilde{v}_{1}, \ldots, \tilde{v}_{n-\mathrm{m}}$ such that $v_{i}=\tilde{v}_{i}+U$. Consider

$$
u_{1} \wedge u_{2} \ldots \wedge u_{m} \wedge \tilde{v}_{1} \wedge \ldots \wedge \tilde{v}_{n-m}
$$

This is independent of the choice of $\tilde{v}_{i}$ since any two choices differ by a linear combination of $u_{i}$, which is annihilated by $u_{1} \wedge \ldots u_{m}$. This map defines the isomorphism. Because it is natural it extends to the case of vector bundles.

Suppose now that $M$ of dimension n is defined as the subset $f^{-1}(c)$ of $R^{n}$ where $f: \boldsymbol{R}^{n} \rightarrow \boldsymbol{R}^{m}$ has surjective derivative on $M$. This means that the 1 -forms $d f_{1}, \ldots d f_{m}$ are linearly independent at the points of $M \subset \boldsymbol{R}^{n}$. We saw that in this situation, the tangent space $T_{a} M$ of $M$ at $a$ is the subspace of $T_{a} \boldsymbol{R}^{n}$ annihilated by the derivative of $f$, or equivalently the 1 -forms $d f_{i}$. Another way of saying this is that the cotangent space $T_{a}^{*} M$ is the quotient of $T_{a}^{*} \boldsymbol{R}^{n}$ by the subspace $U$ spanned by $d f_{1}, \ldots, d f_{m}$. From (3.8) we have an isomorphism

$$
\Lambda^{m} U \otimes \wedge^{n-m}\left(T^{*} M\right) \cong \wedge^{n} T^{*} \boldsymbol{R}^{n}
$$

Now $d f_{1} \wedge d f_{2} \wedge \ldots \wedge d f_{m}$ is a non-vanishing section of $\wedge^{m} U$ and $d x_{1} \wedge \ldots \wedge$ $d x_{n}$ is a non-vanishing section of $\Lambda^{n} T R^{n}$ so the isomorphism defines a nonvanishing section $\omega$ of $\wedge^{n-m} T^{*} M$.

All such manifolds, and not just the sphere, are therefore orientable. In the case $m=1$, where M is defined by a single real-valued function $f$, we have

$$
d f \wedge \omega=d x_{1} \wedge d x_{2} \ldots \wedge d x_{n}
$$

If $\partial f / \partial x_{n} \neq 0$, then $x_{1}, \ldots, x_{n-1}$ are local coordinates and so from this formula we see that

$$
\omega=(-1)^{n-1} \frac{1}{\partial f / \partial x_{n}} d x_{1} \wedge \ldots \wedge d x_{n-1}
$$

as above.

## Remark (3.2.4):

Any compact manifold $M^{m}$ can be embedded in $R^{N}$ for some $N$, but the argument above shows that M is not always cut out by $N-m$ globally defined functions with linearly independent derivatives, because it would then have to be orientable.

Orientability helps in integration through the following:

## Proposition (3.2.5):

A manifold is orientable if and only if it has a covering by coordi-nate charts such that

$$
\operatorname{det}\left(\frac{\partial y_{i}}{\partial x_{j}}\right)>0
$$

on the intersection.

## Proof:

Assume $M$ is orientable, and let $\omega$ be a non-vanishing $n$-form. In a coordinate chart

$$
\omega=f\left(x_{1}, \ldots, x_{n}\right) d x_{1} \wedge \ldots d x_{n}
$$

After possibly making $a$ coordinate change $x_{1} \mapsto c-x_{1}$, we have coordinates such that $f>0$.

Look at two such overlapping sets of coordinates. Then

$$
\begin{gathered}
\omega=g\left(y_{1}, \ldots, y_{n}\right) d y_{1} \wedge \ldots \wedge d y_{n}=g\left(y_{1}(x), \ldots, y_{n}(x)\right)\left(\operatorname{det} \partial y_{i} / \partial x_{j}\right) d x_{1} \wedge \\
d x_{2} \ldots \wedge d x_{n}=f\left(x_{1}, \ldots, x_{n}\right) d x_{1} \wedge \ldots d x_{n}
\end{gathered}
$$

Since $f>0$ and $g>0$, the determinant $\operatorname{det} \partial y_{i} / \partial x_{j}$ is also positive.
Conversely, suppose we have such coordinates. Take a partition of unity $\left\{\varphi_{\alpha}\right\}$ subordinate to the coordinate covering and put

$$
\omega=\sum \varphi_{\alpha} d y_{1}^{\alpha} \wedge d y_{2}^{\alpha} \wedge \ldots \wedge d d y_{n}^{\alpha}
$$

Then on a coordinate neighbourhood $U_{\beta}$ with coordinates $x_{1}, \ldots, x_{n}$ we have

$$
\left.\omega\right|_{U_{\beta}}=\sum \varphi_{\alpha} \operatorname{det}\left(\partial y_{i}^{\alpha} / \partial x_{j}\right) d x_{1} \wedge \ldots d x_{n}
$$

Since $\varphi_{\alpha} \geq 0$ and $\operatorname{det}\left(\partial y_{i}^{\alpha} / \partial x_{j}\right)$ is positive, this is non-vanishing.
Now suppose $M$ is orientable and we have chosen an orientation. We shall define the integral

$$
\int_{M} \theta
$$

of any $n$-form $\theta$ of compact support on $M$.
We first choose a coordinate covering as in Proposition (3.2.5). On each coordinate neighbourhood $U_{\alpha}$ we have

$$
\left.\theta\right|_{U_{\alpha}}=f_{\alpha}\left(x_{1}, \ldots, x_{n}\right) d x_{1} \wedge \ldots \wedge d x_{n}
$$

Take a partition of unity $\varphi_{i}$ subordinate to this covering. Then

$$
\left.\varphi_{i} \theta\right|_{U_{\alpha}}=\mathrm{g}_{\mathrm{i}}\left(x_{1}, \ldots, x_{n}\right) \mathrm{d} x_{1} \wedge \ldots \wedge d x_{n}
$$

where $g_{i}$ is a smooth function of compact support on the whole of $\mathrm{R}^{\mathrm{n}}$. We then define
$\int_{\mathrm{M}} \theta=\sum_{\mathrm{i}} \int_{\mathrm{M}} \varphi_{\mathrm{i}} \theta=\sum_{\mathrm{i}} \int_{\mathrm{R}^{\mathrm{n}}} g_{\mathrm{i}}\left(x_{1}, \ldots, x_{\mathrm{n}}\right) d x_{1} \ldots d x_{\mathrm{n}}$.
Note that since $\theta$ has compact support, its support is covered by finitely many open sets on which $\varphi_{i} \neq 0$, so the above is a finite sum.

The integral is well-defined precisely because of the change of variables formula in integration, and the consistent choice of sign from the orientation.

Now we will study Stokes' theorem. The theorems of Stokes and Green in vector calculus are special cases of a single result in the theory of differential forms, which by convention is called Stokes' theorem. We begin with a simple version of it:

## Theorem (3.2.6):

Let $M$ be an oriented $n$-dimensional manifold and $\omega \in \Omega^{\mathrm{n}-1}(\mathrm{M})$ be of compact support. Then

$$
\int_{\mathrm{M}} \mathrm{~d} \omega=0
$$

## Proof:

Use a partition of unity subordinate to a coordinate covering to write

$$
\omega=\sum \varphi_{i} \omega
$$

Then on a coordinate neighbor hood

$$
\varphi_{i} \omega=a_{1} d x_{2} \wedge \ldots \wedge d x_{n}-a_{2} d x_{1} \wedge d x_{3} \wedge \ldots \wedge d x_{n}+\cdots
$$

and

$$
d\left(\varphi_{i} \omega\right)=\left(\frac{\partial a_{1}}{\partial x_{1}}+\ldots+\frac{\partial a_{n}}{\partial x_{n}}\right) d x_{1} \wedge d x_{2} \wedge \ldots \wedge d x_{n}
$$

From the definition of the integral, we need to sum each

$$
\int_{R^{n}}\left(\frac{\partial a_{1}}{\partial x_{1}}+\ldots+\frac{\partial a_{n}}{\partial x_{n}}\right) d x_{1} d x_{2} \ldots d x_{n}
$$

Consider

$$
\int_{R^{n}} \frac{\partial a_{1}}{\partial x_{1}} d x_{1} d x_{2} \ldots d x_{n}
$$

By Fubini's theorem we evaluate this as a repeated integral

$$
\int_{R} \int_{R} \ldots\left(\int \frac{\partial a_{1}}{\partial x_{1}} d x_{1}\right) d x_{2} d x_{3} \ldots d x_{n}
$$

But $a_{1}$ has compact support, so vanishes if $\left|x_{1}\right| \geq N$ and thus

$$
\int_{R} \frac{\partial a_{1}}{\partial x_{1}} d x_{1}=\left[a_{1}\right]_{-N}^{N}=0
$$

The other terms vanish in a similar way.

Theorem (3.2.6) has an immediate payoff for de Rham cohomology:

## Proposition (3.2.7):

Let $M$ be a compact orientable $n$-dimensional manifold. Then the de Rham cohomology group $H^{n}(M)$ is non-zero.

## Proof:

Since $M$ is orientable, it has a non-vanishing $n$-form $\theta$. Because there are no $n+1$-forms, it is closed, and defines a cohomology class $[\theta] \in H^{n}(M)$.

Choose the orientation defined by $\theta$ and integrate: we get

$$
\int_{\mathrm{M}} \theta=\sum \int f_{i} d x_{1} d x_{2} \ldots d x_{n}
$$

which is positive since each $f_{i} \geq 0$ and is positive somewhere.
Now if the cohomology class $[\theta]=0, \theta=d \omega$, but then Theorem (3.2.6) gives

$$
\int_{\mathrm{M}} \theta=\int_{\mathrm{M}} d \omega=0
$$

a contradiction.

Here is a topological result which follows directly from the proof of the above fact:

## Theorem (3.2.8):

Every vector field on an even-dimensional sphere $S^{2 m}$ vanishes somewhere.

## Proof:

Suppose for a contradiction that there is a non-vanishing vector field. For the sphere, sitting inside $R^{2 m+1}$, we can think of a vector field as a smooth map

$$
v: S^{2 m} \rightarrow R^{2 m+1}
$$

such that $(x, v(x))=0$ and if $v$ is non-vanishing we can normalize it to be a unit vector. So assume $(v(x), v(x))=1$.

Now define Ft : $S^{2 m} \rightarrow \boldsymbol{R}^{2 m+1}$ by
$F_{t}(x)=\cos t x+\sin t v(x):$
Since $(x, v(x))=0$, we have

$$
(\cos t x+\sin t v(x), \cos t x+\sin t v(x))=1
$$

so that $F_{t}$ maps the unit sphere to itself. Moreover,

$$
F_{0}(x)=x, \quad F_{\pi}(x)=-x
$$

Now let $\omega$ be the standard orientation form on S2m:

$$
\omega=d x_{1} \wedge d x_{2} \wedge \ldots \wedge d x_{2 m} / x_{2 m+1}
$$

We see that

$$
F_{0}^{*} \omega=\omega, F_{\pi}^{*} \omega=-\omega
$$

But by Theorem (3.1.19), the maps $F_{0}^{*}, F_{\pi}^{*}$ on $H^{2 m}\left(S^{2 m}\right)$ are equal. We deduce that the de Rham cohomology class of $\omega$ is equal to its negative and so must be zero, but this contradicts that fact that its integral is positive. Thus the vector field must have $a$ zero.

Green's theorem relates a surface integral to a volume integral, and the full version of Stokes' theorem does something similar for manifolds. The manifolds we have defined are analogues of a surface - the sphere for example. We now need to find analogues of the solid ball that the sphere bounds. These are still called manifolds, but with a boundary.

## Definition (3.2.9):

An n-dimensional manifold with boundary is a set $M$ with a collection of subsets $U_{\alpha}$ and maps

$$
\varphi_{\alpha}: U_{\alpha} \rightarrow\left(R^{n}\right)^{+}=\left\{\left(x_{1}, \ldots, x_{n}\right) \in R^{n}: x_{n} \geq 0\right\}
$$

such that
i. $M=\mathrm{U}_{\alpha} U_{\alpha}$
ii. $\varphi_{\alpha}: U_{\alpha} \rightarrow \varphi_{\alpha}\left(U_{\alpha}\right)$ is a bijection onto an open set of $\left(R^{n}\right)^{+}$and $\varphi_{\alpha}\left(U_{\alpha} \cap U_{\beta}\right)$ is open for all $\alpha, \beta$,
iii. $\varphi_{\beta} \varphi_{\alpha}^{-1}: \varphi_{\alpha}\left(U_{\alpha} \cap U_{\beta}\right) \rightarrow \varphi_{\beta}\left(U_{\alpha} \cap U_{\beta}\right)$ is the restriction of a $C^{\infty}$ map from a neighborhood of $\varphi_{\alpha}\left(U_{\alpha} \cap U_{\beta}\right) \subseteq\left(R^{n}\right)^{+} \subset R^{n}$ to $R^{n}$.

The boundary $\partial \mathrm{M}$ of M is defined as

$$
\partial M=\left\{x \in M: \varphi_{\alpha}(x) \in\left\{\left(x_{1}, \ldots, x_{n-1}, 0\right) \in R^{n}\right\}\right.
$$

and these charts define the structure of an $(n-1)$-manifold on $\partial M$.

## Example (3.2.10):

1. The model space $\left(R^{n}\right)^{+}$is a manifold with boundary $x_{n}=0$.
2. The unit ball $\left\{x \in R^{n}:\|x\| \leq 1\right\}$ is a manifold with boundary $S^{n-1}$.
3. The Möbius band is a 2-dimensional manifold with boundary the circle:

4. The cylinder $I \times S^{\prime}$ is a 2-dimensional manifold with boundary the union of two circles - a manifold with two components.


We can define differential forms etc. On manifolds with boundary in a straightforward way. Locally, they are just the restrictions of smooth forms on some open set in $R^{n}$ to $\left(R^{n}\right)^{+}$. A form on $M$ restricts to a form on its boundary.

## Proposition (3.2.11):

If a manifold M with boundary is oriented, there is an induced orientation on its boundary.

## Proof:

We choose local coordinate systems such that $\partial M$ is defined by $x_{n}=0$ and $\operatorname{det}\left(\partial y_{i} / \partial x_{j}\right)>0$. So, on overlapping neighbourhoods,

$$
y_{i}=y_{i}\left(x_{1}, \ldots, x_{n}\right), y_{n}\left(x_{1}, \ldots, x_{n-1}, 0\right)=0
$$

Then the Jacobian matrix has the form

$$
\left(\begin{array}{cccc}
\partial y_{1} / \partial x_{1} & \partial y_{1} / \partial x_{2} & \cdots & \partial y_{1} / \partial x_{n}  \tag{3.9}\\
& & \cdots & \cdots: \\
::: & \cdots & \cdots & \ddot{:}_{3} \\
0 & 0 & 0 & \partial y_{n} / \partial x_{n}
\end{array}\right)
$$

From the definition of manifold with boundary, $\varphi_{\beta} \varphi_{\alpha}^{-1}$ maps $x_{n}>0$ to $y_{n}>0$, so yn has the property that if $x_{n}=0, y_{n}=0$ and if $x_{n}>0, y_{n}>0$ It follows that

$$
\left.\frac{\partial y_{n}}{\partial x_{n}}\right|_{x_{n=0}}>0
$$

From (3.9) the determinant of the Jacobian for $\partial M$ is given by

$$
\left.\operatorname{det}\left(J_{\partial M}\right) \frac{\partial y_{n}}{\partial x_{n}}\right|_{x_{n=0}}=\operatorname{det}\left(J_{M}\right)
$$

so if $\operatorname{det}\left(J_{M}\right)>0$ so is $\operatorname{det}\left(J_{\partial M}\right)$.

## Remark (3.2.12):

The boundary of an oriented manifold has an induced orientation, but there is a convention about which one to choose: for a surface in $\boldsymbol{R}^{\mathbf{3}}$ this is the choice of an "inward" or "outward" normal. Our choice will be that if $d x_{1} \wedge \ldots \wedge d x_{n}$ defines the orientation on $M$ with $x_{n} \geq 0$ defining $M$ locally, then $(-1)^{n} d x_{1} \wedge \ldots \wedge d x_{n-1}$ (the "outward" normal) is the induced orientation on $\partial M$. The boundary of the cylinder gives opposite orientations on the two circles. The Möbius band is not orientable, though its boundary the circle of course is.

We can now state the full version of Stokes' theorem:

## Theorem (3.2.13): (Stokes' theorem)

Let $M$ be an $n$-dimensional oriented manifold with boundary $\partial M$ and let $\omega \in \Omega^{n-1}(M)$ be a form of compact support. Then, using the induced orientation

$$
\int_{M} d \omega=\int_{\partial M} \omega .
$$

## Proof:

We write again

$$
\omega=\sum \varphi_{i} \omega
$$

and then

$$
\int_{M} d \omega=\sum \int_{M} d\left(\varphi_{i} \omega\right)
$$

We work as in the previous version of the theorem, with

$$
\begin{gathered}
\varphi_{i} \omega=a_{1} d x_{2} \wedge \ldots \wedge d x_{n}-a_{2} d x_{1} \wedge d x_{3} \wedge \ldots \wedge d x_{n}+\cdots+(-1)^{n-1} a_{n} d x_{1} \wedge \\
d x_{2} \wedge \ldots \wedge d x_{n-1}
\end{gathered}
$$

(3.2.6), but now there are two types of open sets. For those which do not intersect $\partial M$ the integral is zero by Theorem (3.2.6). For those which do, we have

$$
\begin{gathered}
\int_{M} d\left(\varphi_{i} \omega\right)=\int_{x_{n} \geq 0}\left(\frac{\partial a_{1}}{\partial x_{1}}+\ldots+\frac{\partial a_{n}}{\partial x_{n}}\right) d x_{1} d x_{2} \ldots d x_{n}= \\
\int_{R^{n-1}}\left[a_{n}\right]_{0}^{\infty} d x_{1} \ldots d x_{n-1}=-\int_{R^{n-1}} a_{n}\left(x_{1}, x_{2}, \ldots, x_{n-1}, 0\right) d x_{1} \ldots d x_{n-1}= \\
\int_{\partial M} \varphi_{i} \omega
\end{gathered}
$$

where the last line follows since

$$
\left.\varphi_{i} \omega\right|_{\partial M}=(-1)^{n-1} a_{n} d x_{1} \wedge d x_{2} \wedge \ldots \wedge d x_{n-1}
$$

and we use the induced orientation $(-1)^{n} d x_{1} \wedge \ldots \wedge d x_{n-1}$.
An immediate corollary is the following classical result, called the Brouwer fixed point theorem.

## Theorem (3.2.14):

Let B be the unit ball $\left\{x \in \boldsymbol{R}^{n}:\|x\| \leq 1\right\}$ and let $F: B \rightarrow B$ be a smooth map from $B$ to itself. Then $F$ has a fixed point.

## Proof:

Suppose there is no fixed point, so that $F(x) \neq x$ for all $x \in B$. For each $x \in B$, extend the straight line segment $\overline{F(x) x}$ until it meets the boundary sphere of $B$ in the point $f(x)$. Then we have a smooth function

$$
f: B \rightarrow \partial B
$$

such that if $x \in \partial B, f(x)=x$.
Let $\omega$ be the standard non-vanishing $(n-1)$-form on $S^{n-1}=\partial B$, with

$$
\int_{\partial B} \omega=1 .
$$

Then

$$
1=\int_{\partial B} \omega=\int_{\partial B} f^{*} \omega
$$

since $f$ is the identity on $S^{n-1}$. But by Stokes' theorem,

$$
\int_{\partial B} f^{*} \omega=\int_{B} d\left(f^{*} \omega\right)=\int_{B} f^{*}(d \omega)=0
$$

since $d \omega=0$ as $\omega$ is in the top dimension on $S^{n-1}$.

The contradiction $l=0$ means that there must be a fixed point.

## Chapter(4)

## Smooth Map and Riemannian Metric

## Section (4.1): The Degree of Smooth Map

We begin this section by studying the degree of a smooth map.By using integration of forms we have seen that for a compact orientable manifold of dimension $n$ the de Rham cohomology group $H^{n}(M)$ is non-zero, and that this fact enable us to prove some global topological results about such manifolds. We shall now refine this result, and show that the group is (for a compact, connected, orientable manifold ) just one-dimensional. This gives us a concrete method of determining the comology class of an $n$-form: it is exact if and only if its integral is zero.

Now we will study de Rham cohomology in the top dimension. First a lemma:

## Lemma(4.1.1):

Let $U^{n}=\left\{x \in R^{n}:\left|x_{i}\right|<1\right\}$ and let $\omega \in \Omega^{n}\left(R^{n}\right)$ be a form with support in $U^{n}$ such that

$$
\int_{U^{n}} \omega=0
$$

Then there exists $\beta \in \Omega^{n-1}\left(R^{n}\right)$ with support in $U^{n}$ such that $\omega=d \beta$

## Proof:

We prove the result by induction on the dimension $n$, but we make the inductive assumption that $\omega$ and $\beta$ depend smoothly on a parameter $\lambda \in R^{m}$, and also that if $\omega$ vanishes identically for some $\lambda$, so dose $\beta$.

Consider the case $n=1$, so $=f(x, \lambda) d x$. Clearly taking

$$
\begin{equation*}
\beta(x, \lambda)=\int_{-1}^{x} f(u, \lambda) d u \tag{4.1}
\end{equation*}
$$

Gives us a function with $d \beta=\omega$. But also, since f has support in $U$,there is a $\delta>0$ such that $f$ vanishes for $x>1-\delta$ or $x<-1+\delta$. Thus

$$
\int_{-1}^{x} f(u, \lambda) d u=\int_{-1}^{1} f(u, \lambda) d u=0
$$

For $x>1-\delta$ and similarly for $x<-1+\delta$ which means that $\beta$ itself has support in. If $f(x, \lambda)=0$ for all, then from the integration (4.1) so dose $\beta(x, \lambda)$. Now assume the result for dimensions less than $n$ and let

$$
\omega=f\left(x_{1}, \ldots, x_{n}, \lambda\right) d x_{1} \wedge \ldots \wedge d x_{n-1}
$$

Be the given form. Fix $x_{n}=t$ and consider

$$
f\left(x_{1}, \ldots, x_{n-1}, t, \lambda\right) d x_{1} \wedge \ldots \wedge d x_{n-1}
$$

As a form on $R^{n-1}$, depending smoothly on $t$ and $\lambda$. Its integral is no longer zero, but if $\sigma$ is a bump funtion on $U^{n-1}$ such that the integral of $\sigma d x_{1} \wedge \ldots \wedge d x_{n-1}$ is 1 , then putting

$$
g(t, \lambda)=\int_{U^{n-1}} f\left(x_{1}, \ldots, x_{n-1}, t, \lambda\right) d x_{1} \wedge \ldots \wedge d x_{n-1}
$$

We have a form

$$
f\left(x_{1}, \ldots, x_{n-1}, t, \lambda\right) d x_{1} \wedge \ldots \wedge d x_{n-1}-g(t, \lambda) \sigma d x_{1} \wedge \ldots \wedge d x_{n-1}
$$

With support in $U^{n-1}$ and zero integral. Apply induction to this and we can write it as $d \gamma$ where $\gamma$ has support in $U^{n-1}$.

Now put $=x_{n}$, and consider $d\left(\gamma \wedge d x_{n}\right)$. The $x_{n}$-derivative of $\gamma$ doesn't contribute because of the $d x_{n}$ factor, and $\sigma$ is independent of $x_{n}$, so we get

$$
d\left(\gamma \wedge d x_{n}\right)=f\left(x_{1}, \ldots, x_{n-1}, t, \lambda\right) d x_{1} \wedge \ldots \wedge d x_{n-1}-g(t, \lambda) \sigma d x_{1} \wedge \ldots \wedge d x_{n}
$$

Putting

$$
\xi\left(x_{1}, \ldots, x_{n}, \lambda\right)=(-1)^{n-1}\left(\int_{-1}^{x_{n}} g(t, \lambda) d t\right) \sigma d x_{1} \wedge \ldots \wedge d x_{n-1}
$$

Also gives

$$
d \xi=g\left(x_{n}, \lambda\right) \sigma d x_{1} \wedge \ldots \wedge d x_{n}
$$

We can therefore write

$$
f\left(x_{1}, \ldots, x_{n-1}, t, \lambda\right) d x_{1} \wedge \ldots \wedge d x_{n}=d\left(\gamma \wedge d x_{n}+\xi\right)=d \beta
$$

Now by construction $\beta$ has support in $\left|x_{i}\right|<1$ for $1 \leq i \leq n-1$, but what a but the $x_{n}$ direction? Since $f\left(x_{1}, \ldots, x_{n-1}, t, \lambda\right)$ vanishes for $t>1-\delta$ or $<-1+\delta$, the inductive assumption tells us that $\gamma$ does also for $x_{n}>1-\delta$.as for $\xi$, if $t>1-\delta$

$$
\begin{aligned}
\int_{-1}^{t} g(s, \lambda) d s & =\int_{-1}^{t}\left(\int_{U^{n-1}} f\left(x_{1}, \ldots, x_{n-1}, t, \lambda\right) d x_{1} \wedge \ldots \wedge d x_{n-1}\right) d t \\
& =\int_{-1}^{1}\left(\int_{U^{n-1}} f\left(x_{1}, \ldots, x_{n-1}, t, \lambda\right) d x_{1} \wedge \ldots \wedge d x_{n-1}\right) d t \\
& =\int_{U^{n}} f\left(x_{1}, \ldots, x_{n}, \lambda\right) d x_{1} \wedge \ldots \wedge d x_{n}=0
\end{aligned}
$$

By assumption. Thus the support of $\xi$ is in $U^{n}$ again, examining the integrals, if $f(x, \lambda)$ is identically zero for some $\lambda$ so is $\beta$.

Using the lemma, we prove :

## Theorem (4.1.2):

If $M$ is a compact , connected orientable $n$-dimensional manifold, then $H^{n}(M) \cong R$

## Proof:

Take a covering by coordinate neighbourhoods which map to $U^{n}=\{x \in$ $\left.R^{n}:\left|x_{i}\right|<1\right\}$ and a corresponding partition of unity $\left\{\varphi_{i}\right\}$. by compactness, we can assume we have a finite number $U_{1}, \ldots, U_{N}$ of open sets. Using a bump function , fix an $n$-form $\alpha_{0}$ with support in $U_{1}$ and

$$
\int_{M} \alpha_{0}=1
$$

Thus, by theorem (3.2.7) the cohomology class $\left[\alpha_{0}\right]$ is non-zero. To prove the theorem we want to show that for any $n$-form $\alpha$,

$$
[\alpha]=c\left[\alpha_{0}\right]
$$

i.e that $\alpha=c \alpha_{0}+d \gamma$.
given $\alpha$ use the partition of unity to write

$$
\alpha=\sum \varphi_{i} \alpha
$$

then by linearity it is sufficient to prove the result for each $\varphi_{i} \alpha$, so we may assume that the support of $\alpha$ lies in a coordinate neigbourhood $U_{m}$. Because $M$ is connected we can connect $p \in U_{1}$ and $q \in U_{m}$ by a path and by the connectedness of open
intervals we can assume that the path is covered by a sequence of $U_{i}$ 's, each of which intersects the next : i.e. renumbering we have.

$$
p \in U_{1}, \quad U_{i} \cap U_{i+1} \neq \emptyset, \quad q \in U_{m}
$$

now for $\quad 1 \leq i \leq m-1$ choose an $n$-form $\alpha_{i}$ with support in $U_{i} \cap U_{i+1}$ and integral 1. On $U_{1}$ we have

$$
\int\left(\alpha_{0}-\alpha_{1}\right)=0
$$

and so applying lemma (4.1.1), there is a form $\beta_{0}$ with support in $U_{1}$ such that

$$
\alpha_{0}-\alpha_{1}=d \beta_{1}
$$

continuing, we get altogether

$$
\begin{gathered}
\alpha_{0}-\alpha_{1}=d \beta_{1} \\
\alpha_{1}-\alpha_{2}=d \beta_{2} \\
\ldots=\cdots \\
\alpha_{m-2}-\alpha_{m-1}=d \beta_{m-1}
\end{gathered}
$$

and adding, we find

$$
\alpha_{0}-\alpha_{m-1}=d\left(\sum_{i} \beta_{i}\right)
$$

On $U_{m}$, we have

$$
\int \alpha=c=\int c \alpha_{m-1}
$$

and applying the limma again, we get $\alpha-c \alpha_{m-1}=d \beta$ and so from (4.2)

$$
\alpha=c \alpha_{m-1}+d \beta=c \alpha_{0}+d\left(\beta-c \sum_{i} \beta_{i}\right)
$$

as required.

Theorem (4.2.1) tells us that for a compact connected oriented $n$-dimensional manifold, $H^{n}(M)$ is one-dimensional. Take a form $\omega_{M}$ whose integral over $M$ is 1 , then $\left[\omega_{M}\right]$ is a natural basis element for $H^{n}(M)$. suppose

$$
F: M \rightarrow N
$$

Is a smooth map of compact connected oriented manifolds of the same dimension $n$. Then we have the induced map

$$
F^{*}: H^{n}(N) \rightarrow H^{n}(M)
$$

And relative to our bases

$$
\begin{equation*}
F^{*}\left[\omega_{N}\right]=k\left[\omega_{M}\right] \tag{4.3}
\end{equation*}
$$

For some real number $k$. we now show that $k$ is an integer.

## Theorem(4.1.3):

Let $M, N$ be oriented, compact, connected manifolds of the same dimension $n$, and $F: M \rightarrow N$ a smooth map. There exists an integer, called degree of $F$ such that
(i) If $\omega \in \Omega^{n}(N)$ then

$$
\int_{M} F^{*} \omega=\operatorname{deg} \int_{N} \omega
$$

(ii) If $a$ is a regular value of $F$ then

$$
\operatorname{deg} F=\sum_{x \in F^{-1}(a)} \operatorname{sgn}\left(\operatorname{deg} D F_{x}\right)
$$

## Remark(4.1.4):

1. A regular value for a smooth map $F: M \rightarrow N$ is a point $a \in N$ such that for each $x \in F^{-1}(a)$. The derivative $D F_{x}$ is surjective. When $\operatorname{dim} M=\operatorname{dim} N$ this means that $D F_{x}$ is invertible. Sard's theorem shows that for any smooth map most points in $N$ are regular values.
2. The expression $\operatorname{sgn}\left(\operatorname{det} D F_{x}\right)$ in the theorem can be interpreted in two ways, but depends crucially on the notion of orientation - consistently associating the right sign for all the point $x \in F^{-1}(a)$. The straightforword approach uses proposition (3.2.5) to associate to an orientation a class of coordinates whose jacobians have positive determinant. If $\operatorname{det} D F_{x}$ is written as a Jacobian matrix in such a set of coordinates for M and N , then $\operatorname{sgn}\left(\operatorname{det} D F_{x}\right)$ is just the sign of the determinant. More invariantly, $D F_{x}: T_{x} M \rightarrow T_{a} N$ defines a linear map

$$
\Lambda^{n}\left(D F_{x}\right): \wedge T^{*} N_{a} \rightarrow \Lambda T_{x}^{*} M
$$

Orientations on M and N are define by non-vanishing forms $\omega_{M}, \omega_{N}$ and

$$
\Lambda^{n}\left(D F_{x}\right)\left(\omega_{N}\right)=\lambda_{\omega} M
$$

Then $\operatorname{sgn}\left(\operatorname{det} D F_{x}\right)$ is the sign of $\lambda$.
3. Note the immediate corollary of the theorem: if $F$ is not surjective, then deg $F=0$.

## Proof:

For the first part of the theorem, the cohomology class of $\omega$ is $[\omega]=c\left[\omega_{N}\right]$ and so integrating ( and using proposition(3.2.6)),

$$
\int_{N} \omega=c \int_{N} \omega_{N}=c
$$

Using the number k in (4.3)

$$
F^{*}[\omega]=c F^{*}\left[\omega_{N}\right]=\operatorname{ck}\left[\omega_{M}\right]
$$

And integration,

$$
\int_{M} F^{*} \omega=c k \int_{M} \omega_{M}=c k=k \int_{N} \omega
$$

For the second part, since $D F_{x}$ is an isomorphism at all points in $F^{-1}(a)$, from theorem (1.2.8) , $F^{-1}(a)$ is a zero -dimensional manifold. Since it is compact (closed inside a compact space $M$ ) it is a finite set of points. The inverse function theorem applied to these $m$ points shows that there is a coordinate neighbouhood Uof $a \in N$ such that $F^{-1}(U)$ is a disjoint union of $m$ open sets $U_{i}$ such that

$$
F: U_{i} \rightarrow U
$$

Is a diffeomorphism .

Let $\sigma$ be an $n$-form supported in U with $\int_{N} \sigma=1$ and consider the diffeomorphism $F: U_{i} \rightarrow U$. Then by the coordinate invariance of integration of forms, and using the orientations on $M$ and $N$,

$$
\int_{U_{i}} F^{*} \sigma=\operatorname{sgn} D F_{x_{i}} \int_{U} \sigma=\operatorname{sgn} D F_{x_{i}}
$$

Hence, summing

$$
\int_{M} F^{*} \sigma=\sum_{i} \operatorname{sgn} D F_{x_{i}}
$$

And this is from the first part

$$
k=k \int_{N} \sigma=\int_{M} F^{*} \sigma
$$

Which gives

$$
k=\sum_{i} \operatorname{sgn} D F_{x_{i}}
$$

## Example(4.1.5):

Let $M$ be the extended complex plane : $M=\boldsymbol{C} \cup\{\infty\}$.this is a compact, connected , orientable 2-manifold. In fact it is the 2 -sphere . Consider the map $F: M \rightarrow$ $M$ defined by

$$
\begin{gathered}
F(z)=z^{k}+a_{1} z^{k-1}+\cdots+a_{k,} z \neq \infty \\
F(\infty)=\infty
\end{gathered}
$$

This is smooth because in coordinate near $z=\infty, F$ is defined (for $\omega=1 / z$ ) by

$$
\omega \mapsto \frac{\omega^{k}}{1+a_{1} \omega+\cdots+a_{k} \omega^{k}}
$$

To find the degree of , consider

$$
F_{t}(z)=z^{k}+t\left(a_{1} z^{k-1}+\cdots+a_{k}\right)
$$

This is a smooth map for all $t$ and by theorem (3.1.19) the action on cohomology is independent of , so

$$
\operatorname{deg} F=\operatorname{deg} F_{0}
$$

Where $F_{0}(z)=z^{k}$

We can calculate this degree by taking a 2-form, with $|z|=r$ and $z=x+i y$

$$
f(r) d x \wedge d y=f(r) r d r \wedge d \theta
$$

With $f(r)$ of compact support. Then the degree is given by

$$
\operatorname{deg} F_{0} \int_{R^{2}} f(r) r d r \wedge d \theta=\int_{R^{2}} f\left(r^{k}\right) r^{k} d\left(r^{k}\right) k d \theta=k \int_{R^{2}} f(r) r d r \wedge d \theta
$$

Thus $\operatorname{deg} F=k$. if $k>0$ this means in particular that $F$ is surjective and therefore takes the value 0 somewhere, so that

$$
z^{k}+a_{1} z^{k-1}+\cdots+a_{k}=0
$$

Has a solution. This is the fundamental theorem of algebra.

## Example(4.1.6):

Take two smooth maps $f_{1}, f_{2}: S^{1} \rightarrow R^{3}$. These give two circle in $R^{3}$ - suppose they are disjoint. Define

$$
F: S^{1} \times S^{1} \rightarrow S^{2}
$$

By

$$
F(s, t)=\frac{f_{1}(s)-f_{2}(t)}{\left\|f_{1}(s)-f_{2}(t)\right\|}
$$

The degree of this map is called the linking number.

## Example(4.1.7):

Let $M \subset R^{3}$ be a compact surface and $\boldsymbol{n}$ its unit normal. The Gauss map is the map

$$
F: M \rightarrow S^{2}
$$

Defined by $(x)=n(x)$. If we work out the degree by integration, we take the standard 2-form $\omega$ on $S^{2}$. Then one finds that

$$
\int_{M} F^{*} \omega=\int_{M} K \sqrt{E G-F^{2}} d u d v
$$

Where $K$ is the Gaussian curvature. The Gauss-Bonnet theorem tells us that the degree is half the Euler characteristic of $M$.

## Section (4.2): The Metric Tensor

We begin this section by study Riemannian metrics . Differential forms and the exterior derivative provide one piece of analysis on manifolds which, as we have seen, links in with global topological questions. There is much more one can do when one introduces a Riemannian metric. Since the whole subject of Riemannian geometry is a huge one, we shall here look at only two aspects which relate to the use of differential forms: the study of harmonic forms and of geodesics.

In particular, we ignore completely here questions related to curvature.
Now we will discuss the metric tensor . In informal terms, Riemannian metric on a manifold $M$ is a smoothly varying positive definite inner product on the tangent spaces $T_{x}$. To make global sense of this, note that an inner product is a bilinear form, so at each point $x$ we want a vector in the tensor product

$$
T_{x}^{*} \otimes T_{x}^{*}
$$

We can put, just as we did for the exterior forms , a vector bundle structure on

$$
T^{*} M \otimes T^{*} M=\mathrm{U}_{x \in M} T_{x}^{*} \otimes T_{x}^{*}
$$

The conditions we need to satisfy for a vector bundle are provided by two facts we used for the bundle of $p$ - forms:
(i) Each coordinate system $x_{1}, \ldots, x_{n}$ defines a basic $d x_{1}, \ldots, d x_{n}$ for each $T_{x}^{*}$ in the coordinate neighbourhood and the $n^{2}$ elements

$$
d x_{i} \otimes d x_{j}, \quad 1 \leq i, j \leq n
$$

Give a corresponding basis for $T_{x}^{*} \otimes T_{x}^{*}$
(ii) The jacobian of a change of coordinates defines an invertible linear transformation $j: T_{x}^{*} \rightarrow T_{x}^{*}$ and we have a corresponding invertible linear transformation $J \otimes J: T_{x}^{*} \otimes T_{x}^{*} \rightarrow T_{x}^{*} \otimes T_{x}^{*}$

Given this, we define :

## Definition (4.2.1):

A Riemannain metric on a manifold $M$ is a section $g$ of $T^{*} \otimes T^{*}$ which at each point is symmetric and positive definite.

In a local coordinate system we can write

$$
g=\sum_{i, j} g_{i j}(x) d x_{i} \otimes d x_{j}
$$

Where $g_{i j}=g_{j i}(x)$ and is a smooth function, with $g_{i j}(x)$ positive definite. Often the tensor product symbol is omitted and one simply writes

$$
g=\sum_{i, j} g_{i j}(x) d x_{i} d x_{j}
$$

## Examples (4.2.2):

1. The Euclidean metric on $R^{n}$ is defined by

$$
g=\sum d x_{i} \otimes d x_{i}
$$

So

$$
g\left(\frac{\partial}{\partial x_{i}}, \frac{\partial}{\partial x_{j}}\right)=\delta_{i j}
$$

2. A submanifold of $R^{n}$ has an induced Riemannain metric: The tangent space at $x$ can be thought of as a subspace of $R^{n}$ and we take the Euclidean inner product on $R^{n}$.

Given a smooth map $F: M \rightarrow N$ and a metric $g$ on, we can pull back $g$ to a section $F^{*} g$ of $T^{*} M \otimes T^{*} M$ :

$$
\left(F^{*} g\right)_{x}(X, Y)=g_{F}(x)\left(D F_{x}(X), D F_{x}(Y)\right)
$$

If $D F_{x}$ is invertible ,this will again be positive definite, so in particular if $F$ is a diffeomorphism.

## Definition (4.2.3):

A diffeomorphism $F: M \rightarrow N$ between two Riemannian manifold is an isometry if $F^{*} g_{N}=g_{M}$.

## Example (4.2.4):

Let $M=\left\{(x, y) \in R^{2}: y>0\right\}$ and

$$
g=\frac{d x^{2}+d y^{2}}{y^{2}}
$$

If $z=x+i y$ and

$$
F(z)=\frac{a z+b}{c z+d}
$$

With $a, b, c$ real and $-b c>0$, then

$$
F^{*} d z=(a d-b c) \frac{d z}{(c z-d)^{2}}
$$

And

$$
F^{*} y=y \circ F=\frac{1}{i}\left(\frac{a z+b}{c z+d}-\frac{a \bar{z}+b}{c \bar{z}+d}\right)=\frac{a d-b c}{|c z+d|^{2}} y
$$

Then

$$
F^{*} g=(a d-b c)^{2} \frac{d x^{2}+d y^{2}}{\left|(c z+d)^{2}\right|^{2}} \frac{|c z+d|^{4}}{(a d-b c)^{2} y^{2}}=\frac{d x^{2}+d y^{2}}{y^{2}}=g
$$

So these Möbius transformations are isometries of a Riemannian metric on the upper half-plane.

We have learned in section ((one- parameter group of diffeomorphisms)) that a one- parameter group $\varphi_{t}$ of diffeomorphism defines a vector field $X$. Then we can define the lie derivative of a Riemannian metric by

$$
\mathcal{L}_{X g}=\left.\frac{d}{d t} \varphi_{t}^{*} g\right|_{t=0}
$$

If this is a group of isometries then since $\varphi_{t}^{*} g=g$, we have $\mathcal{L}_{X g}=0$. Such a vector field is a called a Killing vector field or an infinitesimal isometry.

The lie derivative obeys the usual derivation rules, and commutes with . Since $\mathcal{L}_{X} f=X f$ we have

$$
\begin{aligned}
\mathcal{L}_{X} \sum_{i} g_{i j} d x_{i} \otimes d x_{j}= & \sum_{i}\left(X g_{i j}\right) d x_{i} \otimes d x_{j}+\sum_{i} g_{i j} d\left(X x_{i}\right) \otimes d x_{j} \\
& +\sum_{i} g_{i j} d x_{i} \otimes d\left(X x_{j}\right)
\end{aligned}
$$

## Example (4.2.5):

Take the Euclidean metric $g=\sum_{i} d x_{i} \otimes d x_{i}$, and a vector field of the form

$$
X=\sum_{i, j} A_{k i} x_{k} \frac{\partial}{\partial x_{j}}
$$

Where $A_{i j}$ is a constant matrix.
This is a Killing vector field if and only if

$$
0=\sum_{k, i} d\left(A_{k i} x_{k}\right) \otimes d x_{i}+d x_{i} \otimes d\left(A_{k i} x_{k}\right)=\sum_{k, i}\left(A_{k i}+A_{i k}\right) d x_{i} \otimes d x_{k}
$$

In other words if $A$ is skew0symmetric.

With a Riemannian metric one can define the length of the a curve:

## Definition(4.2.6):

Let $M$ be a Riemannian manifold and $\gamma:[0,1] \rightarrow M$ a smooth map (i.e. a smooth curve in $M$ ). The length of the curve is

$$
\ell(\gamma)=\int_{0}^{1} \sqrt{g\left(\gamma^{\prime}, \gamma^{\prime}\right)} d t
$$

Where $\gamma^{\prime}(t)=D \gamma_{t}(d / d t)$
With this definition, any Riemannian manifold is a metric space : Define

$$
d(x, y)=\inf \{\ell(\gamma) \in R: \gamma(0)=x, \gamma(1)=y\}
$$

In fact a metric defines an inner product on $T^{*}$ as well as on , for the map

$$
X \mapsto g(X,-)
$$

Defines an isomorphism from $T$ to $T^{*}$. In concrete terms, if $g^{*}$ is the inner product on $T^{*}$, then

$$
g^{*}\left(\sum_{j} g_{i j} d x_{j}, \sum_{k} g_{k l} d x_{l}\right)=g_{i k}
$$

Which means that

$$
g^{*}\left(d x_{j}, d x_{k}\right)=g^{j k}
$$

Where $g^{j k}$ denotes the inverse matrix to $g_{j k}$.
We can also define an inner product on the exterior product spaces .

$$
\begin{equation*}
\left(\alpha_{1} \wedge \alpha_{2} \wedge \ldots \wedge \alpha_{p}, \beta_{1} \wedge \beta_{2} \wedge \ldots \wedge \beta_{p}\right)=\operatorname{deg}^{*}\left(\alpha_{i}, \beta_{j}\right) \tag{4.4}
\end{equation*}
$$

In particular, on an $n$-manifold there is an inner product in each fibre of the bundle $\Lambda^{n} T^{*}$. S ince each fibre is one-dimensional there are only two unit vectors $\pm u$.

## Definition(4.2.7):

Let $M$ be an oriented Riemannian manifold, then the volume form is the unique $n$ form $\omega$ of unit length in the equivalence class define by the orientation.
In local coordinates, the definition of the inner product (4.4)gives
$\left(d x_{1} \wedge \ldots \wedge d x_{n}, d x_{1} \wedge \ldots \wedge d x_{n}\right)=\operatorname{det} \mathrm{g}_{\mathrm{ij}}^{*}=\left(\operatorname{detg}_{i j}\right)^{-1}$
Thus if $d x_{1} \wedge \ldots \wedge d x_{n}$ defines the orientation,

$$
\omega=\sqrt{\operatorname{deg} g_{i j}} d x_{1} \wedge \ldots \wedge d x_{n}
$$

On a compact manifold we can integrate this to obtain the total volume - so a metric defines not only lengths but also volumes.

Are Riemannian manifolds special? No , because:

## Proposition(4.2.8):

Any manifold admits a Riemannian metric.

## Proof:

Take a covering by coordinate neighbourhood and a partition of unity subordinate to the covering. On each open set $U_{\alpha}$ we have a metric

$$
g_{\alpha}=\sum_{i} d x_{i}^{2}
$$

In the local coordinate . define

$$
g=\sum \varphi_{i} g_{\alpha(i)}
$$

This sum is well - defined because the supports of $\varphi_{i}$ are locally finite. Since $\varphi_{i} \geq 0$ at each point every term in the sum is positive definite or zero, but at least one is positive definite so the sum is positive definite.

## Section (4.3): The Geodesic Flow.

We begin this section by studying the the geodesic flow. Consider any manifold $M$ and its cotangent bundle $T^{*} M$, with projection to the base $p: T^{*} M \rightarrow M$. Let X be tangent vector to $T^{*} M$ at the point $\xi_{a} \in T_{a}^{*}$. Then

$$
D p_{\xi_{a}}(X) \in T_{a} M
$$

So

$$
\theta(X)=\xi_{a}\left(D p_{\xi_{a}}(X)\right)
$$

defines a canonical 1-form $\theta$ on $T^{*} M$. In coordinates $(x, y) \mapsto \sum_{i} y_{i} d x_{i}$, the projection $p$ is

$$
p(x, y)=x
$$

so if

$$
X=\sum a_{i} \frac{\partial}{\partial x_{i}}+\sum \mathrm{b}_{\mathrm{i}} \frac{\partial}{\partial y_{i}}
$$

then

$$
\theta(X)=\sum_{i} y_{i} d x_{i}\left(D p_{\xi_{a}} X\right)=\sum_{i} y_{i} a_{i}
$$

which gives

$$
\theta=\sum_{i} y_{i} d x_{i}
$$

We now take the exterior derivative

$$
\omega=-d \theta=\sum d x_{i} \wedge d y_{i}
$$

which is the canonical 2 -form on the cotangent bundle. It is non-degenerate, so that the map

$$
X \mapsto i_{X} \omega
$$

from the tangent bundle of $T^{*} M$ to its cotangent bundle is an isomorphism.
Now suppose $f$ is a smooth function on $T^{*} M$. Its derivative is a 1 -form $d f$. Because of the isomorphism above, there is a unique vector field $X$ on $T^{*} M$ such that

$$
i_{X} \omega=d f .
$$

If $g$ is another function with vector field $Y$, then

$$
\begin{equation*}
Y(f)=d f(Y)=i_{Y} i_{X} \omega=-i_{X} i_{Y} \omega=-X(g) \tag{4.5}
\end{equation*}
$$

On a Riemannian manifold there is a natural function on $T^{*} M$ given by the induced inner product: we consider the function on $T^{*} M$ defined by

$$
H\left(\xi_{a}\right)=g^{*}\left(\xi_{a}, \xi_{a}\right)
$$

In local coordinates this is

$$
H(x, y)=\sum_{i j} g^{i j}(x) y_{i} y_{j}
$$

## Deffinition (4.3.1):

The vector field $X$ on $T^{*} M$ given by $i_{X} \omega=d H$ is called the geodesic flow of the metric $g$.

## Definition (4.3.2):

If $\gamma:(a, b) \rightarrow T^{*} M$ is an integral curve of the geodesic flow, then the curve $p(\gamma)$ in $M$ is called a geodesic.

In local coordinates, if the geodesic flow is

$$
X=\sum a_{i} \frac{\partial}{\partial x_{i}}+\sum b_{i} \frac{\partial}{\partial y_{i}}
$$

Then

$$
i_{X} \omega=\sum_{k}\left(a_{k} d y_{k}-b_{k} d x_{k}\right)=d H=\sum_{i j} \frac{\partial g^{i j}}{\partial x_{k}} d x_{k} y_{i} y_{j}+2 \sum_{i j} g^{i j} y_{i} d y_{j}
$$

Thus the integral curves are solutions of

$$
\begin{equation*}
\frac{d x_{k}}{d t}=2 \sum_{j} g^{k j} y_{j} \tag{4.6}
\end{equation*}
$$

$$
\begin{equation*}
\frac{d y_{k}}{d t}=-\sum_{i j} \frac{\partial g^{i j}}{\partial x_{k}} y_{i} y_{j} \tag{4.7}
\end{equation*}
$$

Before we explain why this is a geodesic, just note the qualitative behaviour of these curves. For each point $a \in M$, choose a point $\xi_{a} \in T_{\alpha}^{*}$ a and consider the unique integral curve starting at $\xi_{a}$. Equation (4.6) tells us that the projection of the integral curve is parallel at $a$ to the tangent vector $X_{a}$ such that $g\left(X_{a},-\right)=\xi_{a}$. Thus these curves have the property that through each point and in each direction there passes one geodesic.

Geodesics are normally thought of as curves of shortest length, so next we shall link up this idea with the definition above. Consider the variational problem of looking for critical points of the length functional

$$
\ell(\gamma)=\int_{0}^{1} \sqrt{g\left(\gamma^{\prime}, \gamma^{\prime}\right) d t}
$$

for curves with fixed end-points $\gamma(0)=a, \gamma(1)=b$. For simplicity assume $a, b$ are in the same coordinate neighbourhood. If

$$
F(x, z)=\sum_{i j} g_{i j}(x) z_{i} z_{j}
$$

then the first variation of the length is

$$
\delta \ell=\int_{0}^{1} \frac{1}{2} F^{-1 / 2}\left(\frac{\partial F}{\partial x_{i}} \dot{x}_{l}+\frac{\partial F}{\partial z_{i}} \frac{d \dot{x}_{l}}{d t}\right) d t=\int_{0}^{1} \frac{1}{2} F^{-1 / 2} \frac{\partial F}{\partial x_{i}} \dot{x}_{l}-\frac{d}{d t}\left(\frac{1}{2} F^{-1 / 2} \frac{\partial F}{\partial z_{i}}\right) \dot{x}_{l} d t
$$

on integrating by parts with $\dot{x}_{l}(0)=\dot{x}_{l}(1)=0$. Thus a critical point of the functional is given by

$$
\frac{1}{2} F^{-1 / 2} \frac{\partial F}{\partial x_{i}}-\frac{d}{d t}\left(\frac{1}{2} F^{-1 / 2} \frac{\partial F}{\partial z_{i}}\right)=0
$$

If we parametrize this critical curve by arc length:

$$
s=\int_{0}^{t} \sqrt{g\left(\gamma^{\prime}, \gamma^{\prime}\right) d t}
$$

then $F=1$, and the equation simplifes to

$$
\frac{\partial F}{\partial x_{i}}-\frac{d}{d s}\left(\frac{\partial F}{\partial z_{i}}\right)=0
$$

But this is

$$
\begin{equation*}
\sum \frac{\partial g_{j k}}{\partial x_{i}} \frac{d x_{j}}{d s} \frac{d x_{k}}{d s}-\frac{d}{d s}\left(2 g_{j k} \frac{d x_{k}}{d s}\right)=0 \tag{4.8}
\end{equation*}
$$

But now define $y_{i}$ by

$$
\frac{d x_{k}}{d t}=2 \sum_{j} g^{k j} y_{j}
$$

as in the first equation for the geodesic flow (4.6) and substitute in (4.8) and we get

$$
4 \sum \frac{\partial g_{j k}}{\partial x_{i}} g^{j a} y_{a} g^{k b} y_{b}-\frac{d}{d t}\left(4 g_{i k} g^{k a} y_{a}\right)=0
$$

and using

$$
\sum_{j} g^{i j} g_{j k}=\delta_{k}^{i}
$$

this yields

$$
-\frac{\partial g^{j k}}{\partial x_{i}} y_{j} y_{k}=\frac{d y_{i}}{d t}
$$

which is the second equation for the geodesic flow. (Here we have used the formula for the derivative of the inverse of a matrix $\left.G: D\left(G^{-1}\right)=-G^{-1} D G G^{-1}\right)$.

The formalism above helps to solve the geodesic equations when there are isometries of the metric. If $F: M \rightarrow M$ is a diffeomorphism of $M$ then its natural
action on 1-forms induces a diffeomorphism of $T^{*} M$. Similarly with a oneparameter group $\varphi_{t}$.

Differentiating at $t=0$ this means that a vector field $X$ on $M$ induces a vector field $\tilde{X}$ on $T^{*} M$. Moreover, the 1 -form $\theta$ on $T^{*} M$ is canonically defined and hence invariant under the induced action of any diffeomorphism. This means that

$$
\mathcal{L}_{\tilde{X}} \theta=0
$$

and therefore, using (3.1.13) that

$$
i_{\tilde{X}} d \theta+d\left(i_{\tilde{X}} \theta\right)=0
$$

so since $\omega=-d \theta$

$$
i_{\tilde{X}} \omega=d f
$$

where $f=i_{\tilde{X}} \theta$.

## Proposition (4.3.3):

The function $f$ above is $f\left(\xi_{x}\right)=\xi_{x}\left(X_{x}\right)$.
Proof: Write in coordinates

$$
\tilde{X}=\sum a_{i} \frac{\partial}{\partial x_{i}}+\sum b_{i} \frac{\partial}{\partial y_{i}}
$$

where $\theta=\sum_{i} y_{i} d x_{i}$. Since $\tilde{X}$ projects to the vector field $X$ on $M$, then

$$
X=\sum a_{i} \frac{\partial}{\partial x_{i}}
$$

and
$i_{\tilde{X}} \theta=\sum_{i} a_{i} y_{i}=\xi_{x}\left(X_{x}\right)$
by the definition of $\theta$.

Now let $M$ be a Riemannian manifold and $H$ the function on $T^{*} M$ defined by the metric as above. If $\varphi_{t}$ is a one-parameter group of isometries, then the induced diffeomorphisms of $T^{*} M$ will preserve the function $H$ and so the vector field $\tilde{Y}$ will satisfy

$$
\tilde{Y}(H)=0
$$

But from (4.5) this means that $X(f)=0$ where $X$ is the geodesic flow and $f$ the function $i_{\tilde{Y}} \theta$.This function is constant along the geodesic flow, and is therefore a constant of integration of the geodesic equations.

To see what this constant is, we note that $\tilde{Y}$ is the natural lift of a Killing vector field

$$
Y=\sum_{i} a_{i} \frac{\partial}{\partial x_{i}}
$$

so the function $f$ is $f=\sum_{i} a_{i} y_{i}$.
The first geodesic equation is

$$
\frac{d x_{k}}{d t}=2 \sum_{j} g^{k j} y_{j}
$$

so

$$
\sum_{k} g_{j k} \frac{d x_{k}}{d t}=2 y_{j}
$$

and

$$
f=\frac{1}{2} \sum_{k} g_{j k} a_{j} \frac{d x_{k}}{d t}=\frac{1}{2} g\left(\gamma^{\prime}, X\right)
$$

Sometimes this observation enables us to avoid solving any differential equations as in this example.

## Example (4.3.4):

Consider the metric

$$
g=\frac{d x_{1}^{2}+d x_{2}^{2}}{x_{2}^{2}}
$$

on the upper half plane and its geodesic flow X .

The map $\left(x_{1}, x_{2}\right) \mapsto\left(x_{1}+t, x_{2}\right)$ is clearly a one-parameter group of isometries (the Möbius transformations $z \mapsto z+t$ ) and defines the vector field

$$
Y=\frac{\partial}{\partial x_{1}} .
$$

On the cotangent bundle this gives the function

$$
f(x, y)=y_{1}
$$

which is constant on the integral curve.

The map $z \mapsto e^{t} z$ is also an isometry with vector field

$$
Z=x_{1} \frac{\partial}{\partial x_{1}}+x_{2} \frac{\partial}{\partial x_{2}}
$$

so that

$$
g(x, y)=x_{1} y_{1}+x_{2} y_{2}
$$

is constant.

We also have automatically that $H=x_{2}^{2}\left(y_{1}^{2}+y_{2}^{2}\right)$ is constant since

$$
X(H)=i_{X} i_{X} \omega=0
$$

We therefore have three equations for the integral curves of the geodesic flow:

$$
\begin{aligned}
y_{1} & =c_{1} \\
x_{1} y_{1}+x_{2} y_{2} & =c_{2} \\
x_{2}^{2}\left(y_{1}^{2}+y_{2}^{2}\right) & =c 3
\end{aligned}
$$

Eliminating $y_{1}, y_{2}$ gives the geodesics:

$$
\left(c_{1} x_{1}-c_{2}\right)^{2}+c_{1}^{2} x_{2}^{2}=c_{3}
$$

If $c_{1} \neq 0$, this is a semicircle with centre at $\left(x_{1}, x_{2}\right)=\left(c_{2} / c_{1}, 0\right)$. If $c_{1}=0$ then $y_{1}=0$ and the geodesic equation gives $x_{1}=$ const. Together, these are the straight lines of non-Euclidean geometry.

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